



FWF

Der Wissenschaftsfonds.

***Dark matter as a complex scalar field:
new cosmological constraints
and detectability by LIGO***

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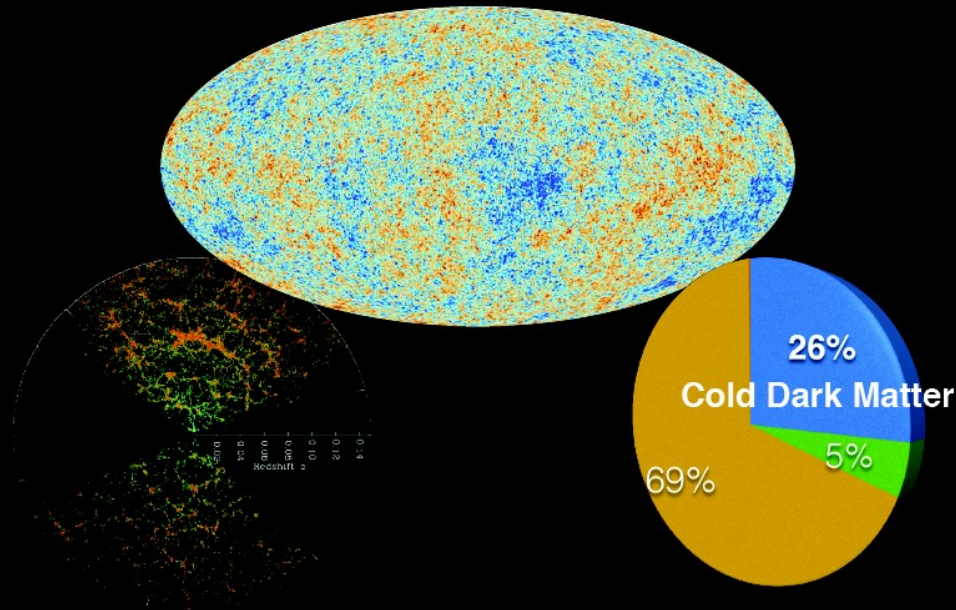
& Michigan Center f. Theoretical Physics, U Michigan

in collaboration with Bohua Li and Paul R Shapiro (UT Austin)

HEP Seminar, Dep.of Physics, IIT Hyderabad

2017/3/15

Scalar Field Dark Matter: Cold Dark Matter Variant



Cold Dark Matter (CDM) candidates

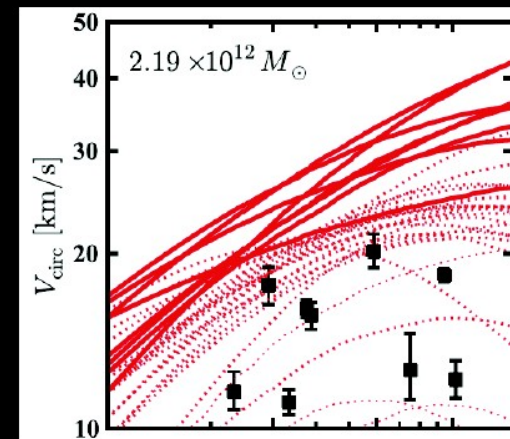
- *Standard* CDM: WIMP, QCD Axion, etc.

- **Scalar Field Dark Matter**

Small-Scale Problems:

Discrepancies between dark matter simulations and observations on sub-halo scales

- **Missing satellites**
- **Cuspy Core**
- **Too big to fail**



Boylan-Kolchin et. al. 2012

Motivation for **Scalar Field Dark Matter (SFDM)**: *alternative DM candidate to explain discrepancies*

SFDM: Ultra-light ($m \ll 10^{-5}$ eV), “cold” bosons ($T \approx 0$) in an (effective) classical field description

Astrophysics motiv':

Suppression of DM clustering below a scale compatible with observations

Suppression of central DM densities in galaxies to explain cored density profiles in DM-dominated (dwarf) galaxies

Accomplish all this, even if we never detect DM

Particle physics motiv':

Ultra-light bosons in extensions of the SM, e.g.

(a multitude of) axion-like particles (“string axiverse”),

other extra-dimensional cosmologies (akin to KK modes),

pseudo-Nambu-Goldstone bosons (upon symmetry breaking in early Universe)

Scalar Field Dark Matter (SFDM), aka Bose-Einstein Condensed Cold Dark Matter (BEC-CDM)

- **Particles created with low entropy per particle → BEC**
→ **classical field Lagrangian for the DM condensate („order parameter“)**
- **complex scalar field** $\psi = |\psi| e^{i\theta}$ → **U(1) symmetry: charge conservation**

$$\mathcal{L} = \frac{\hbar^2}{2m} g^{\mu\nu} \partial_\mu \psi^* \partial_\nu \psi - V(\psi)$$

units: $[\mathcal{L}] = [\text{eV}/\text{cm}^3]$, $[\psi] = \text{cm}^{-3/2}$, (+, -, -, -)

- **assume no coupling to the SM within this EFT description**
- **choice of $V(\psi)$ and initial conditions determine**
 - i) the evolution of SFDM, hence the evolution of the background Universe**
 - ii) suppression of small-scale structure for $L < L_{\text{SFDM}}$ with $L_{\text{SFDM}} = \max\{\lambda_{\text{deB}}, l_{\text{SI}}\}$**

Scalar Field Dark Matter (SFDM), aka Bose-Einstein Condensed Cold Dark Matter (BEC-CDM)

- **choice of potential V:**

(rest-mass) quadratic term, $(mc^2/2) |\psi|^2$

(→ CDM-like in late Universe),

plus a possible repulsive self-interaction, $(\lambda/2) |\psi|^4$

(→ radiation-like in early Universe)

→ **fundamental SFDM parameters:** m and λ

$$\lambda = \hat{\lambda} \frac{\hbar^3}{m^2 c}$$

- **Initial condition:**

conserved charge Q of SFDM determines energy density of DM today

large- Q limit „spintessence“: $\rho_{SFDM,0} = Qmc^2$

(a version of “asymmetric DM”)

$$\rho_{SFDM,0} = n_{SFDM,0} mc^2 = \Omega_{DM} \rho_{crit,0}$$

Equations of motion

Klein-Gordon equation for the SFDM field ψ

$$g^{\mu\nu} \partial_\mu \partial_\nu \psi - g^{\mu\nu} \Gamma^\sigma_{\mu\nu} \partial_\sigma \psi + \frac{m^2 c^2}{\hbar^2} \psi + \frac{2\lambda m}{\hbar^2} |\psi|^2 \psi = 0$$

...which evolves in a **classical GR** background

$$R_{\mu\nu} - \frac{1}{2} g_{\mu\nu} R = \frac{8\pi G}{c^4} T_{\mu\nu}$$

**Well into the matter-dominated era,
they simplify to ...**

nonlinear Schrödinger-Poisson system

$$i\hbar \frac{\partial \psi}{\partial t} = -\frac{\hbar^2}{2m} \Delta \psi + \lambda |\psi|^2 \psi + m\Phi \psi$$

$$\Delta \Phi = 4\pi G m |\psi|^2$$

Numbers for SFDM parameters

„Typical“ numbers of models:

QCD axion: $m \simeq 6 \times 10^{-6} \text{ eV} \frac{10^{12} \text{ GeV}}{f}$ $10^9 \lesssim f \lesssim 10^{12} \text{ GeV}$

(attractive) SI usually neglected: $\lambda \sim 10^{-57}$

$$\lambda = \hat{\lambda} \frac{\hbar^3}{m^2 c}$$

ultra-light axion-like particles: $m \sim (10^{-33} - 1) \text{ eV}$,
f unknown, but reasons for $f \sim 10^{16} \text{ GeV}$

SI usually neglected

(other) bosons for DM: $m \sim (10^{-27} - 1) \text{ eV}$,

SI usually neglected, but we don't: (positive) $\lambda \sim 10^{-93} - 10^{-83}$

For the latter two: Jeans/virial scale can extend to galactic scales $\sim \text{kpc}$!

| | Halo mass [M_\odot] | Size [kpc] | Boson mass [eV] |
|-------------------------|-------------------------|------------|------------------------|
| Milky Way (MW) | 10^{12} | 100 | $1.066 \cdot 10^{-25}$ |
| Dwarf galaxy (DG) | 10^{10} | 10 | $3.371 \cdot 10^{-24}$ |
| Dwarf spheroidal (dSph) | 10^8 | 1 | $1.066 \cdot 10^{-22}$ |
| Minihalo (MH) | 10^6 | 0.1 | $3.371 \cdot 10^{-21}$ |

TRD, Shapiro (1209.1835)

Jeans / virial scales bounded from below by

- no self-interaction: $L_{\text{deB}} \lesssim R \rightarrow m \gtrsim m_H$ (Heisenberg uncertainty “pressure”)

$$m_H := \frac{\hbar}{R^2(\pi G \bar{\rho})^{1/2}} = 1.066 \cdot 10^{-22} \left(\frac{R}{1 \text{ kpc}} \right)^{-1/2} \left(\frac{M}{10^8 M_\odot} \right)^{-1/2} \text{ eV}$$

- or else gravity is balanced by a *positive* self-interaction coupling:

$$L_{\text{deB}} \ll R \rightarrow m \gg m_H \text{ and } \lambda \gg \lambda_H \text{ (self-interaction „pressure“)}$$

$$\lambda_H := \frac{\hbar^2}{2\bar{\rho}R^2} = 2.252 \cdot 10^{-62} \left(\frac{R}{1 \text{ kpc}} \right) \left(\frac{M}{10^8 M_\odot} \right)^{-1} \text{ eV cm}^3$$

ΛSFDM Model (2014) + GW (2016)

2014: take the same cosmic inventory as the basic ΛCDM model, except that

CDM is replaced by SFDM → ΛSFDM (1310.6061, PRD 89, 083536 (2014))

2016: add stochastic GW background (SGWB) from inflation self-consistently to it

(1611.07961)

$$\Omega_m = \Omega_b + \Omega_c$$

Cosmological parameters from **Planck 2013/2015**

(assume SM neutrinos massless)

$$\Omega_\Lambda = 1 - \Omega_m - \Omega_r \quad (2014)$$

$$\Omega_\Lambda = 1 - \Omega_m - \Omega_r - \Omega_{\text{GW}} \quad (2016)$$

- SFDM particle parameters: $m, \lambda/(mc^2)^2$

$$\lambda/(mc^2)^2 = 1 \times 10^{-18} \text{ eV}^{-1} \text{ cm}^3 \Rightarrow l_{SI} \approx 0.8 \text{ kpc}$$

$$\mathcal{L} = \frac{\hbar^2}{2m} g^{\mu\nu} \partial_\mu \psi^* \partial_\nu \psi - \frac{1}{2} mc^2 |\psi|^2 - \frac{\lambda}{2} |\psi|^4,$$

- Global U(1) symmetry ⇒ Charge (particle number density) conservation

$$Q \equiv n - \bar{n} = \rho_{\text{SFDM},0} / (mc^2)$$

- Tensor-to-scalar ratio: r
- Reheating temperature: T_{reheat}

Holistic Evolution of the Λ SFDM Universe

- Friedmann equation

$$H^2(t) \equiv \left(\frac{da/dt}{a}\right)^2 = \begin{cases} H_{\text{inf}}^2, & a < a_{\text{inf}}, \\ H_{\text{inf}}^2 \left(\frac{a_{\text{inf}}}{a(t)}\right)^3, & a_{\text{inf}} < a < a_{\text{reheat}}, \\ \frac{8\pi G}{3c^2} [\rho_r(t) + \rho_b(t) + \rho_\Lambda(t) + \rho_{\text{SFDM}}(t) + \rho_{\text{GW}}(t)], & a > a_{\text{reheat}}, \end{cases}$$

SGWB contribution to the expansion history *self-consistently* included

$$\begin{aligned} \Omega_{\text{GW}}(k, a) &\equiv \frac{d\Omega_{\text{GW}}(a)}{d \ln k} = \frac{1}{\rho_{\text{crit}}(a)} \frac{d\rho_{\text{GW}}(a)}{d \ln k} \\ &= \frac{\Delta_h^2(k, a)c^2}{24a^2 H^2(a)} \left(\left| \frac{h'_k(a(\tau))}{h_k(a(\tau))} \right|^2 + k^2 \right) \end{aligned}$$

- Klein-Gordon Equation

$$\frac{\hbar^2}{2mc^2} \ddot{\psi} + 3 \frac{\hbar^2}{2mc^2} \frac{\dot{a}}{a} \dot{\psi} + \frac{1}{2} mc^2 \psi + \lambda |\psi|^2 \psi = 0,$$

Cosmological evolution of SFDM in an FLRW Universe

Compare size of SF oscillation freq ω to Hubble expansion rate H

- Fast oscillation regime („oscillation“):

$$\omega / H \gg 1$$

disp.relation: $\omega = \omega(V)$,

e.g.

$$\omega = \frac{mc^2}{\hbar} \sqrt{1 + \frac{2\lambda}{mc^2} |\psi|^2}$$

“easier”

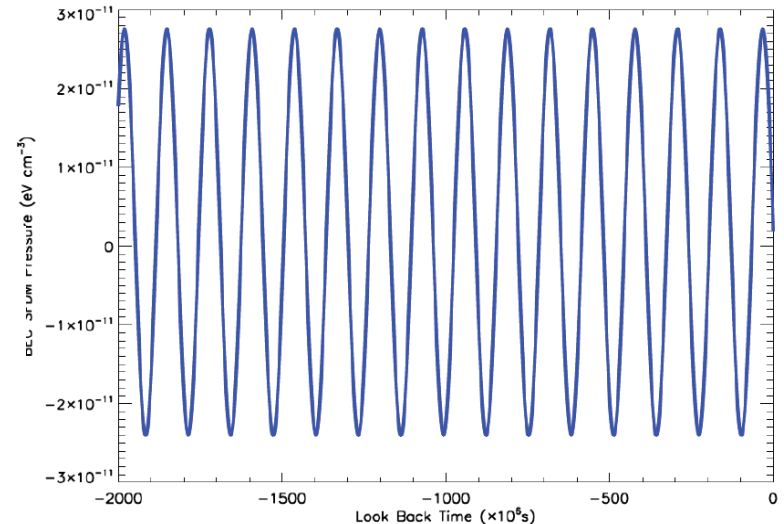
- Slow oscillation regime:

$$\omega / H \ll 1$$

“harder”

kinetic energy $\equiv 0$: $w = -1$ CC EOS

kinetic energy $\neq 0$: $w = 1$ stiff EOS (“kination”)



Cosmological evolution of SFDM in an FLRW Universe

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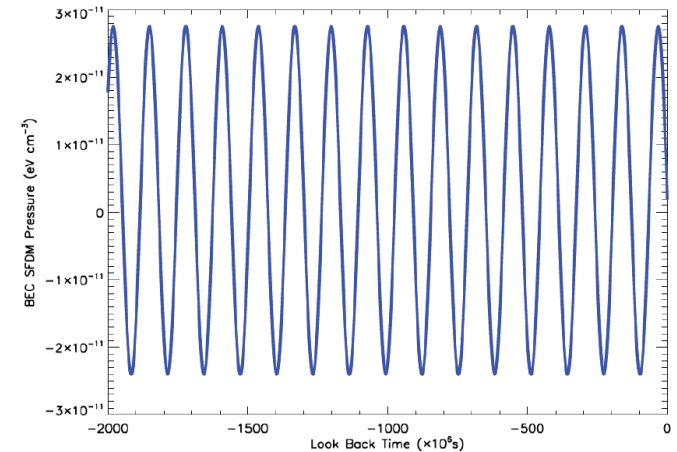
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*affects
the early
Universe
differently!*

Behavior is determined by choice of initial conditions !

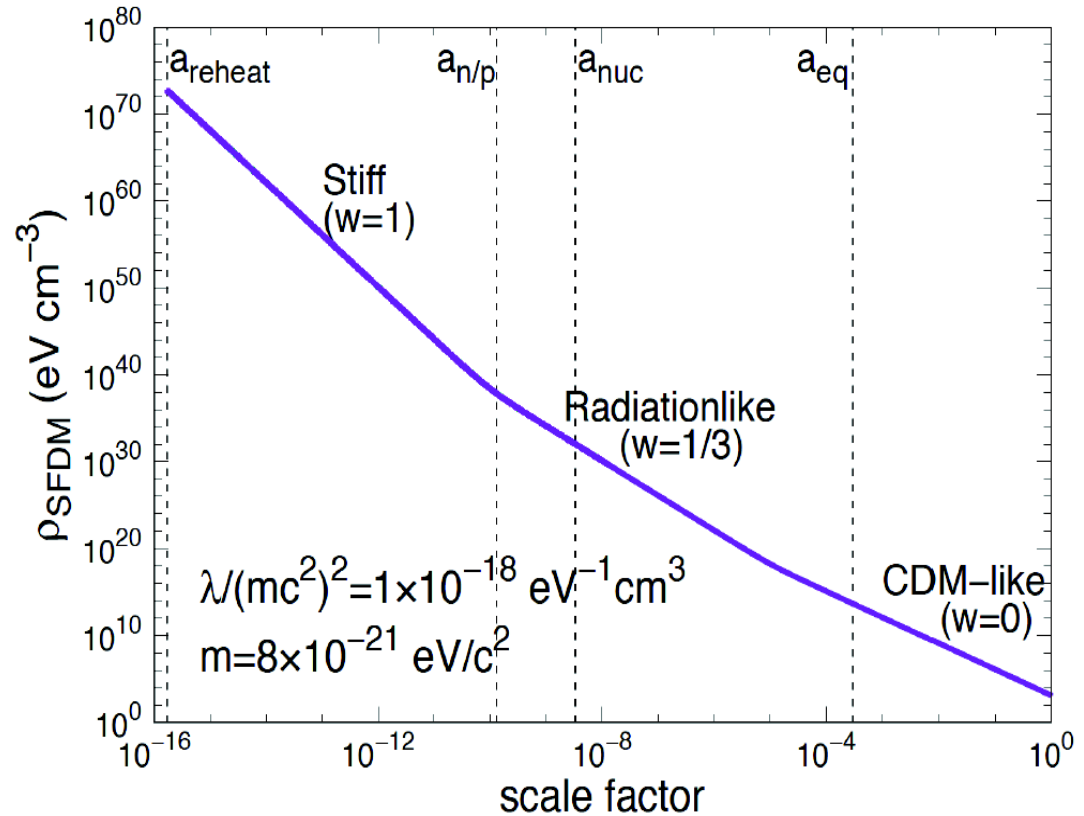
„spintessence“ SFDM with repulsive SI has 3 phases

EOS: $(p/\rho)_{SFDM} = w(t)$

- (1) Early: $w = 1$**
(stiff)
- (2) Intermediate: $w = 1/3$**
(radiationlike)
- (3) Late: $w = 0$**
(non-relativistic matter)

→ change of expansion history !

$\Omega_{SFDM} \rightarrow 1$ at early times



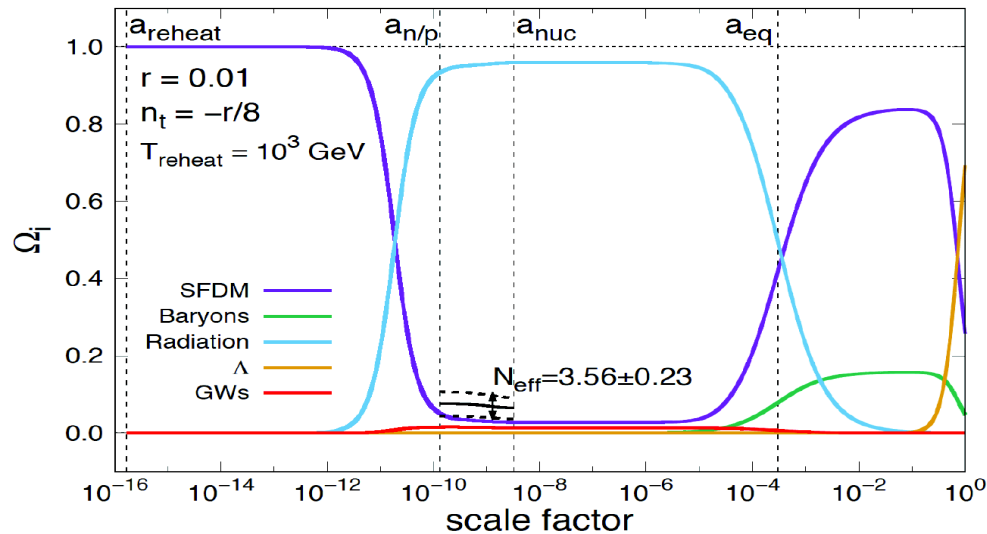
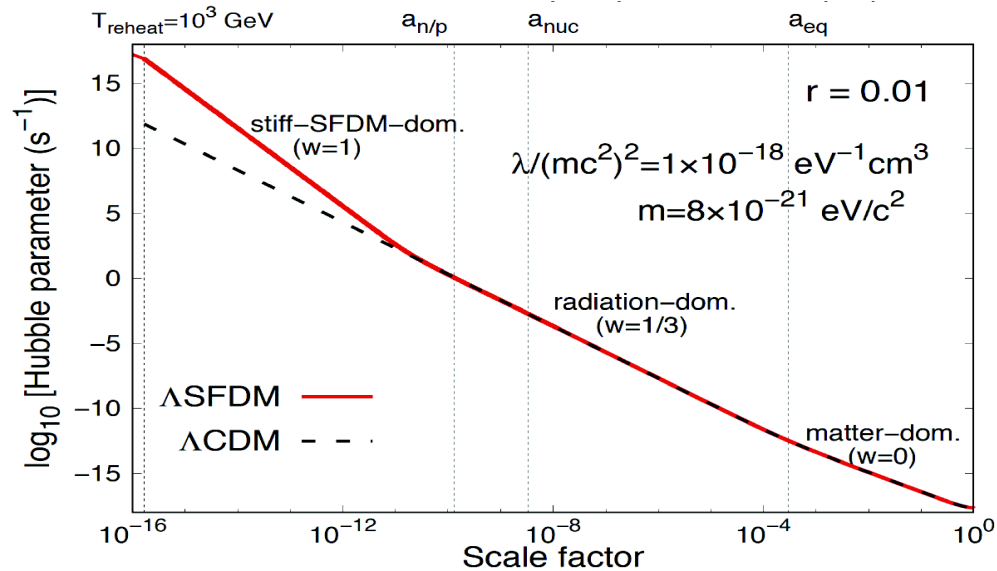
Stiff-SFDM-dominated early Universe

→ additional N_{eff} during (1) and (2)

→ amplifies primordial GWs from inflation during (1)

Λ SFDM+SGWB: the Universe has 6 eras

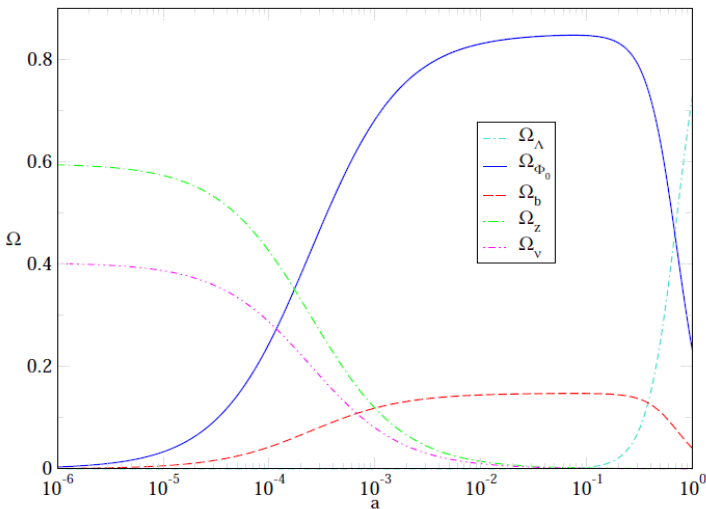
Inflation \rightarrow Reheating \rightarrow Stiff-SFDM-dom. \rightarrow Radiation-dom. \rightarrow Matter-dom. \rightarrow Λ -dom.
 ($w=-1$) ($w=0$) ($w=1$) ($w=1/3$) ($w=0$) ($w=-1$)



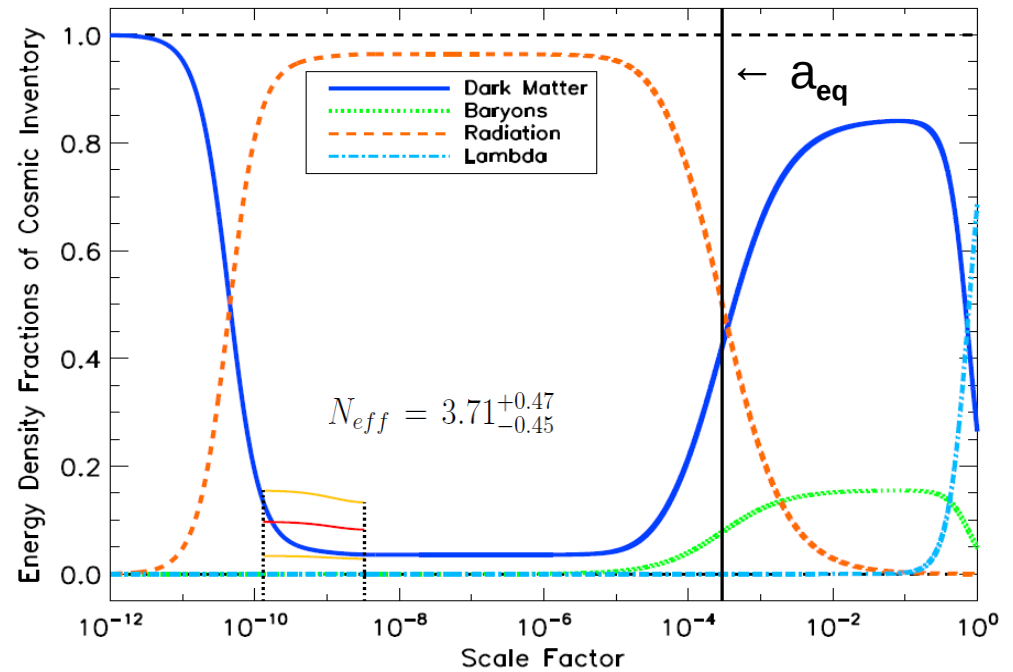
Cosmological evolution in an FLRW Universe of real vs. complex SFDM

Li, TRD, Shapiro (1310.6061)

Magaña, Matos (2012)



$m = 10^{-22}$ eV



$$(m, \lambda)_{\text{fiducial}} = (3 \times 10^{-21} \text{ eV}/c^2, 1.8 \times 10^{-59} \text{ eV cm}^3)$$

Λ SFDM + SGWB

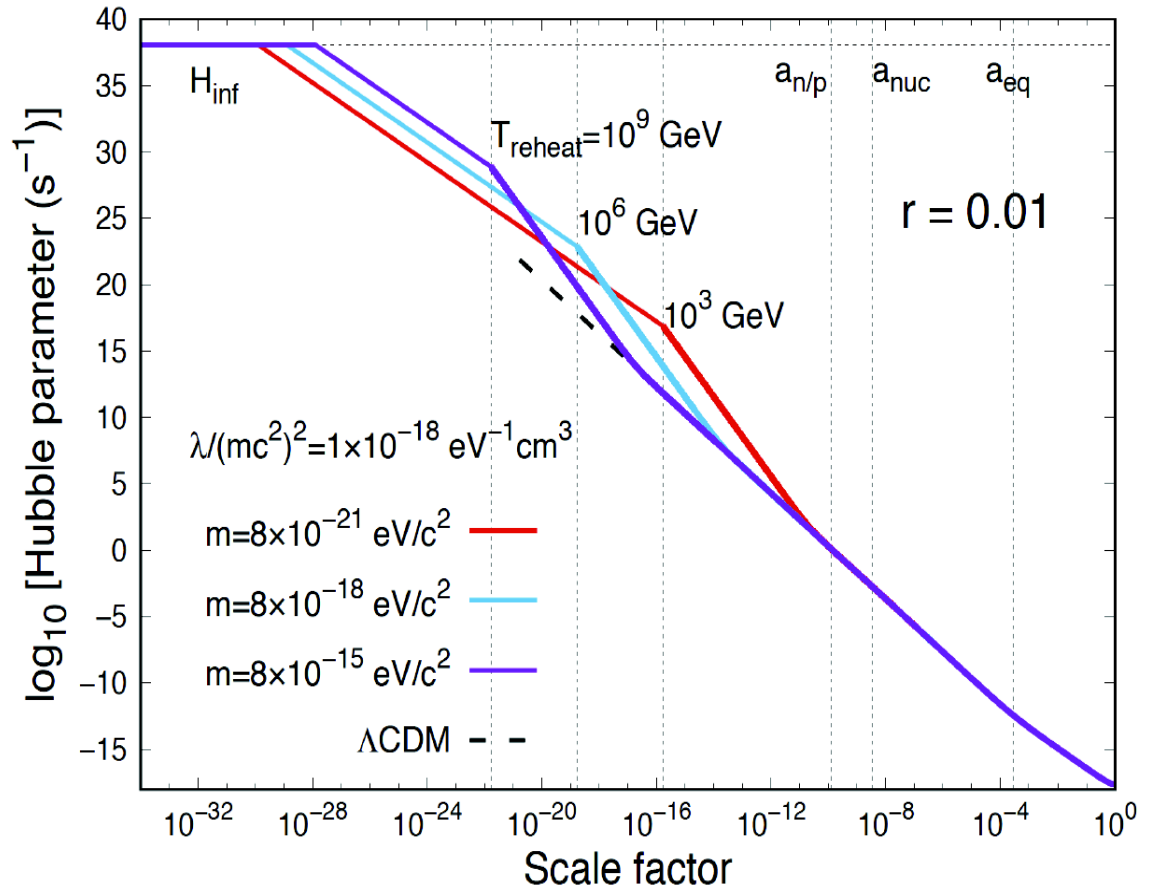
Limiting the duration of the stiff phase after reheating and before BBN constrains SFDM parameters via their contribution to N_{eff}

- for given r :
the smaller the DM mass,
the later must
reheating occur
- Matter-radiation equality:

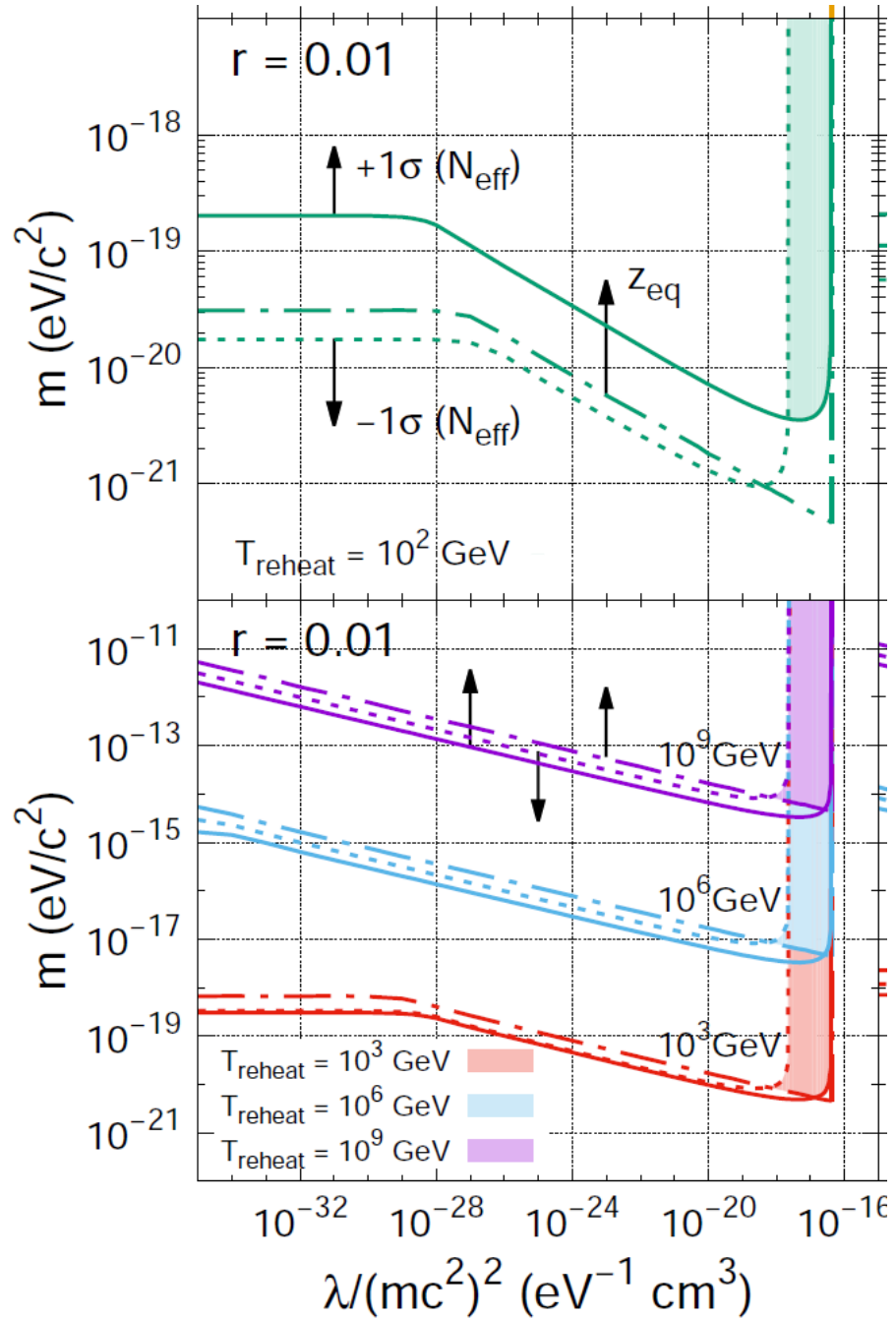
$$1 + z_{\text{eq}} \equiv \frac{1}{a_{\text{eq}}} = \frac{\Omega_b h^2 + \Omega_c h^2}{\Omega_r h^2 + \Omega_{\text{GW}} h^2}$$

- N_{eff} during BBN:

$$\frac{\Delta N_{\text{eff, BBN}}(a)}{N_{\text{eff, standard}}} = \frac{\Omega_{\text{SFDM}}(a) + \Omega_{\text{GW}}(a)}{\Omega_\nu(a)}$$



Cosmological Constraints on the SFDM parameters



ASFDM + SGWB:

enhanced signal of inflationary SGWB due to DM !

Any grav.waves which enter horizon while the background

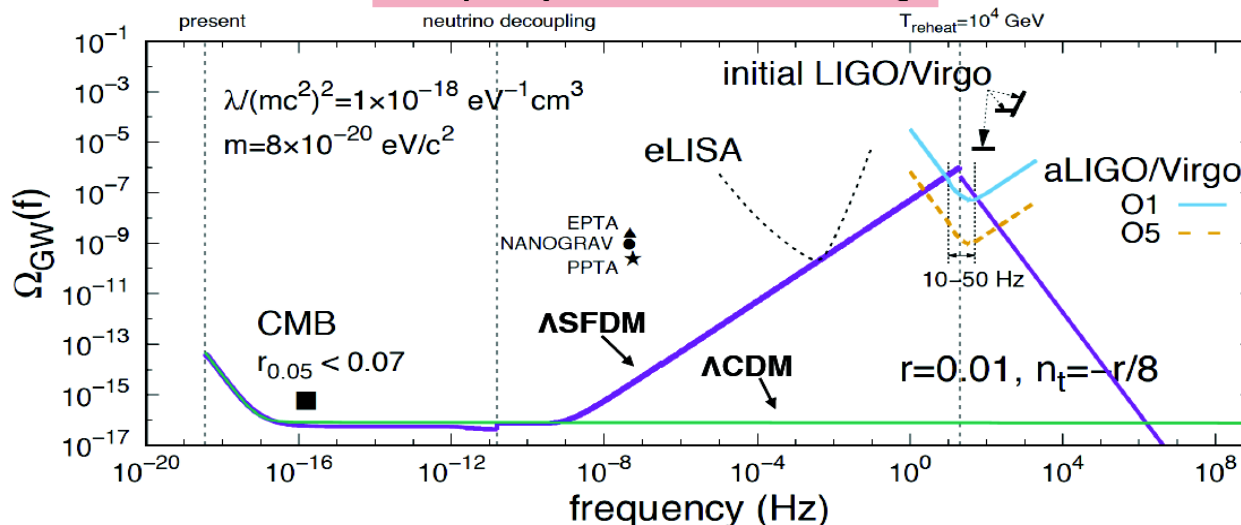
EOS obeys $w > 1/3$ amplifies Ω_{GW} ! (Grishtshuk, Giovannini, Boyle,))

Stiff-SFDM-dominated era amplifies SGWB from (standard) inflation:
can be measured/constrained by GW laser interferometers !

ASFDM predicts 2-parameter broken power-law spectrum at high frequencies:

$$\Omega_{GW}(f) = \Omega_{GW,peak} \times \begin{cases} f / f_{peak}, & f \leq f_{peak} \\ \frac{9\pi}{64} (f / f_{peak})^{-2}, & f > f_{peak} \end{cases}$$

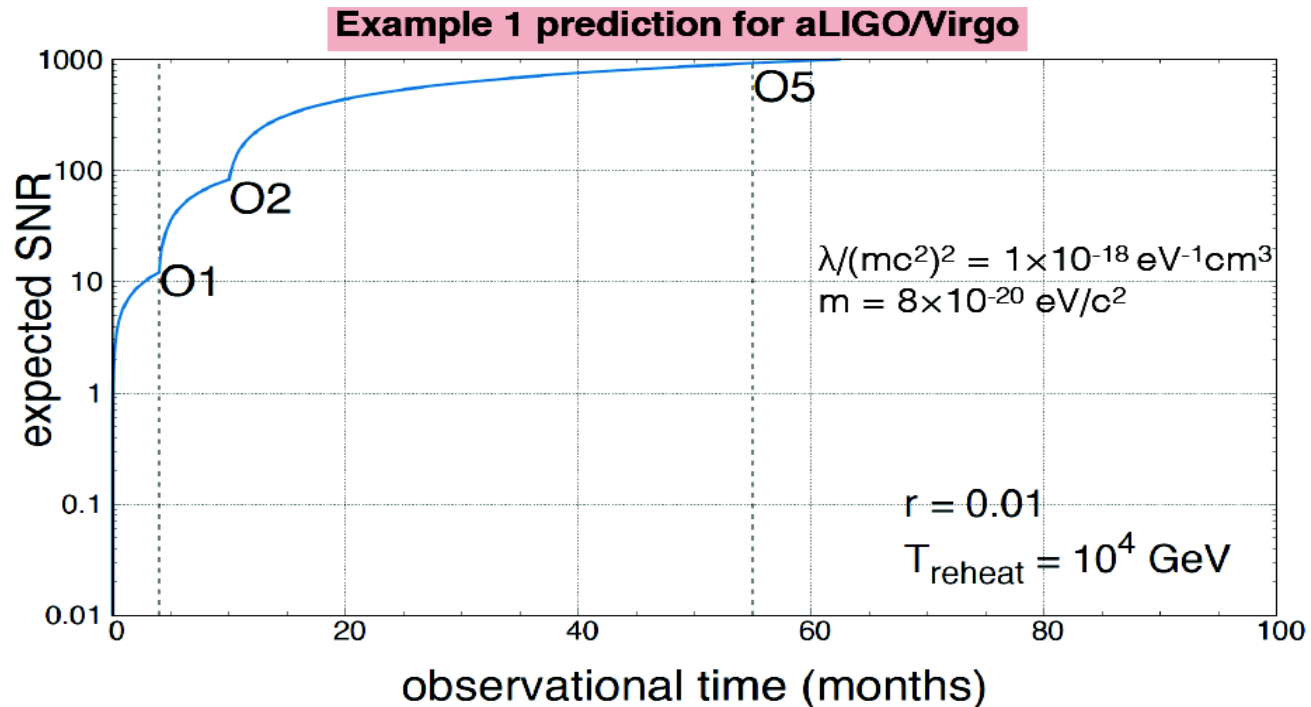
Example 1 prediction for aLIGO/Virgo



ΛSFDM + SGWB:

enhanced signal of inflationary SGWB due to DM !

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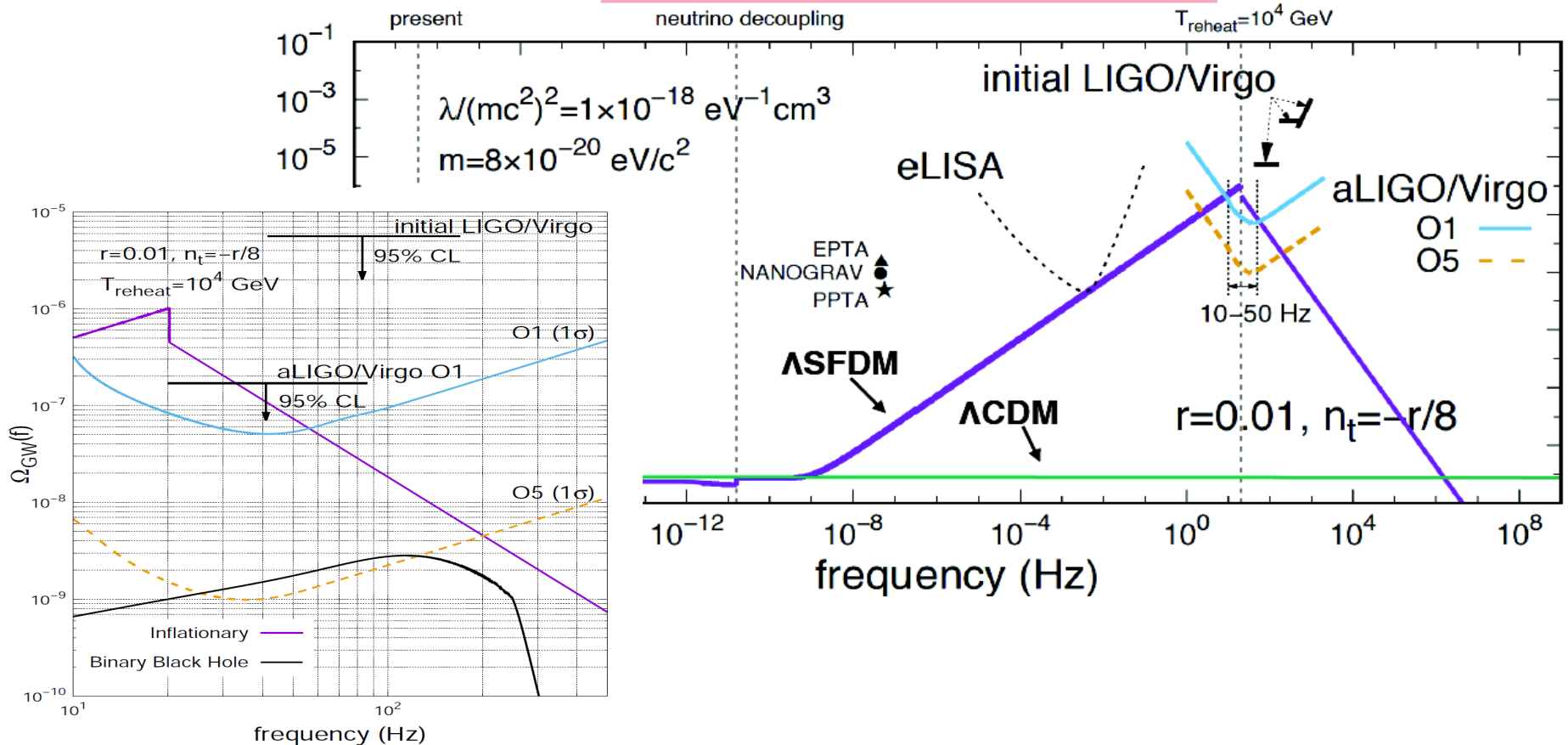
NEW upper limit from O1 excludes this example case at 95% CL
→ **The Age of DM Search/Constraints by GW Detection has begun !**

ΛSFDM + SGWB:

enhanced signal of inflationary SGWB due to DM !

**Stiff-SFDM-dominated era amplifies SGWB from (standard) inflation:
can be measured/constrained by GW laser interferometers !**

Example 1 prediction for aLIGO/Virgo

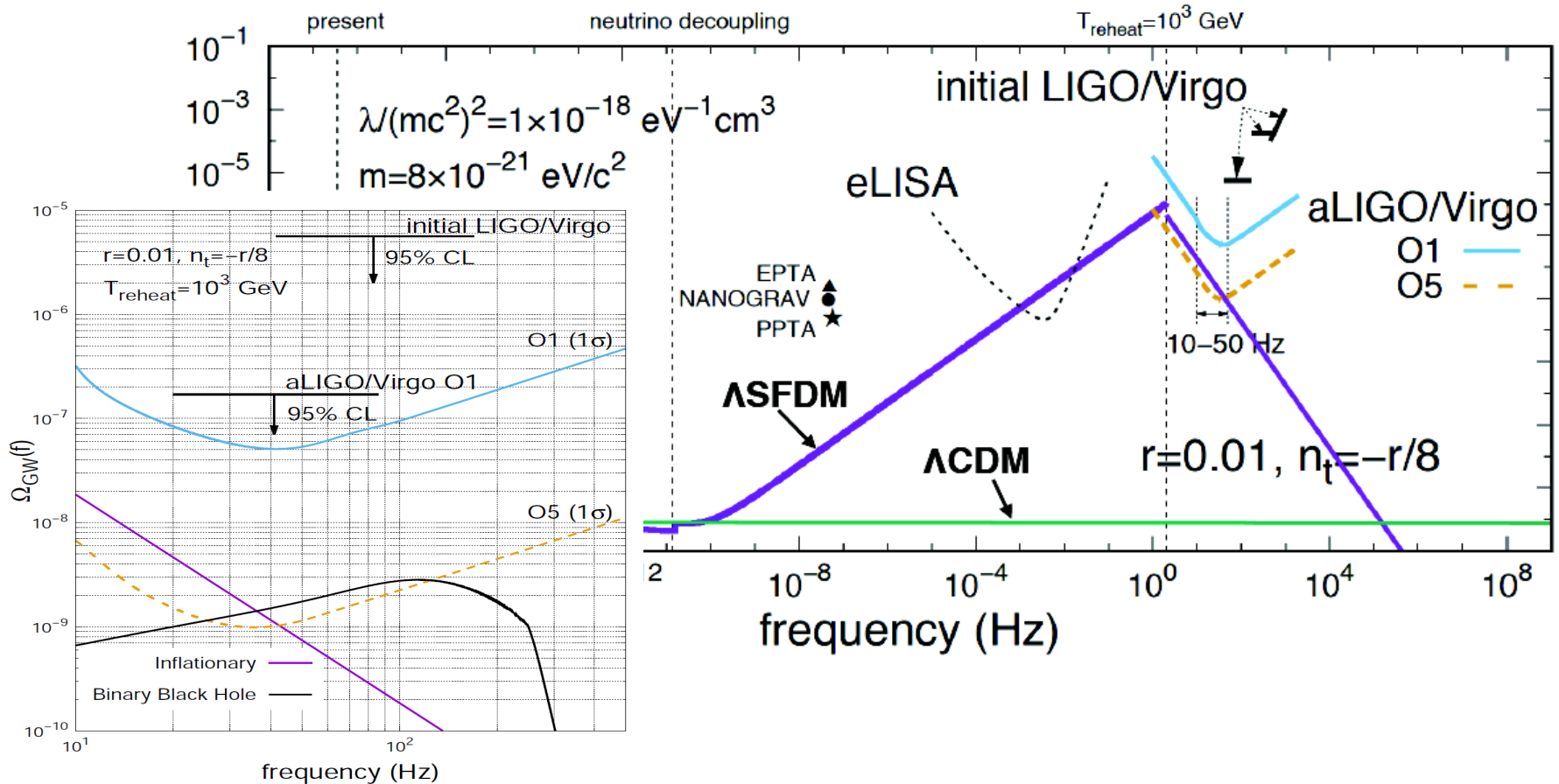


ΛSFDM + SGWB:

enhanced signal of inflationary SGWB due to DM !

Stiff-SFDM-dominated era amplifies SGWB from (standard) inflation:
can be measured/constrained by GW laser interferometers !

Example 2 prediction for aLIGO/Virgo

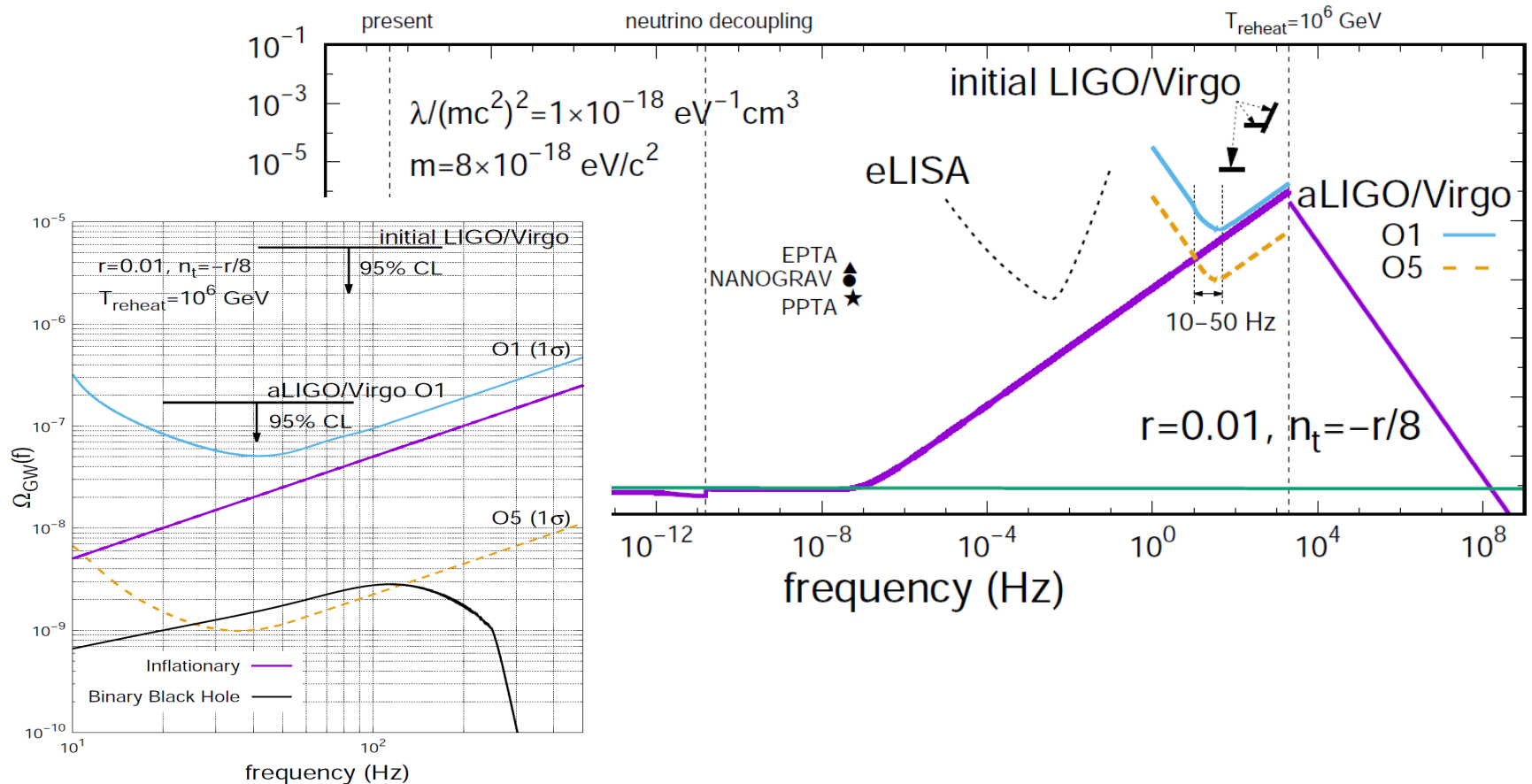


ASFDM + SGWB:

enhanced signal of inflationary SGWB due to DM !

**Stiff-SFDM-dominated era amplifies SGWB from (standard) inflation:
can be measured/constrained by GW laser interferometers !**

Example 3 prediction for aLIGO/Virgo

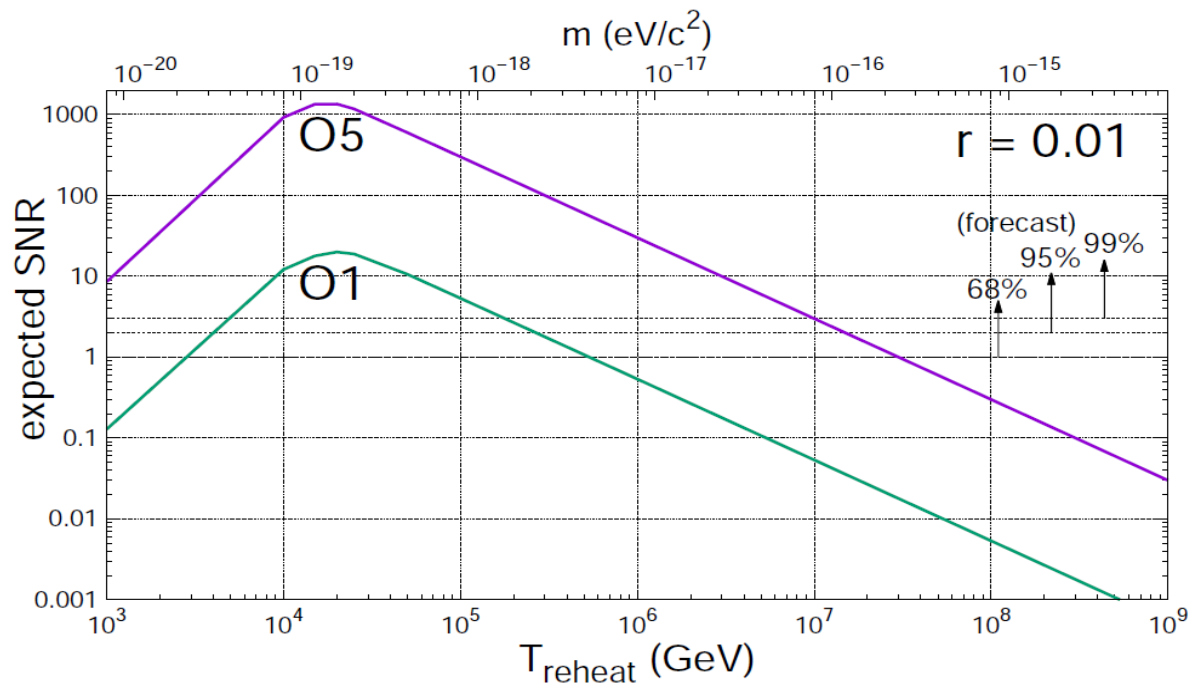


Λ SFDM + SGWB:

enhanced signal of inflationary SGWB due to DM !

Stiff-SFDM-dominated era amplifies SGWB from (standard) inflation:

can be measured/constrained by GW laser interferometers ! \rightarrow FORECASTS



Marginally allowed Λ SFDM models for $\lambda/(mc^2)^2 = 1 \times 10^{-18} \text{ eV}^{-1} \text{ cm}^3$

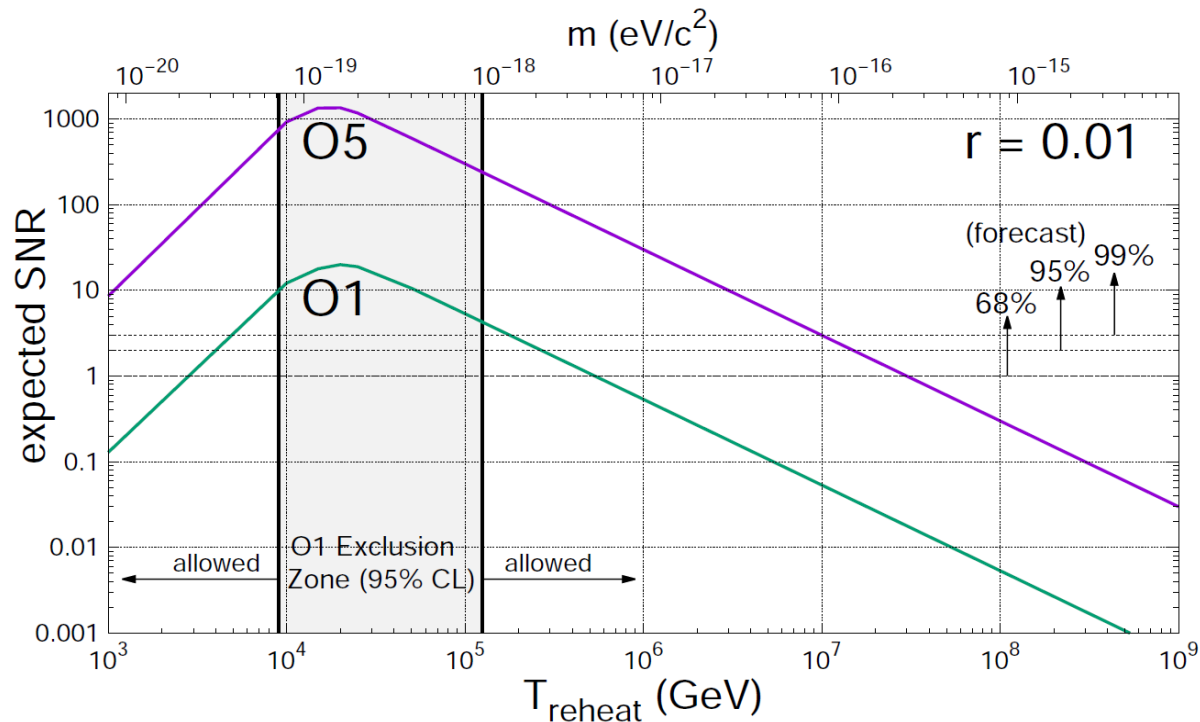
Λ SFDM + SGWB:

enhanced signal of inflationary SGWB due to DM !

Stiff-SFDM-dominated era amplifies SGWB from (standard) inflation:

can be measured/constrained by GW laser interferometers !

Current limits: O1 from LIGO (1612.02029)



Marginally allowed Λ SFDM models for $\lambda/(mc^2)^2 = 1 \times 10^{-18} \text{ eV}^{-1} \text{ cm}^3$

Conclusions: SFDM is a good DM candidate

- may resolve small-scale problems of CDM *structure formation*
 - rich variety of models allows non-standard *expansion histories* in the early Universe:
 - particle parameters are constrained by the CMB, BBN, other PTs, primordial GWs from inflation (SGWB)
 - puts into reach the *possible detection of the inflationary SGWB*:
 - a wide range of DM particle parameters and reheat temperatures can be already tested by aLIGO/VIRGO O1 run, and more with O5 !
 - (some examples are already ruled out from O1)
- **ongoing/upcoming GW laser interferometer experiments can detect/constrain DM !**