

Splashback radius as probes of cosmology, dark matter and galaxy evolution

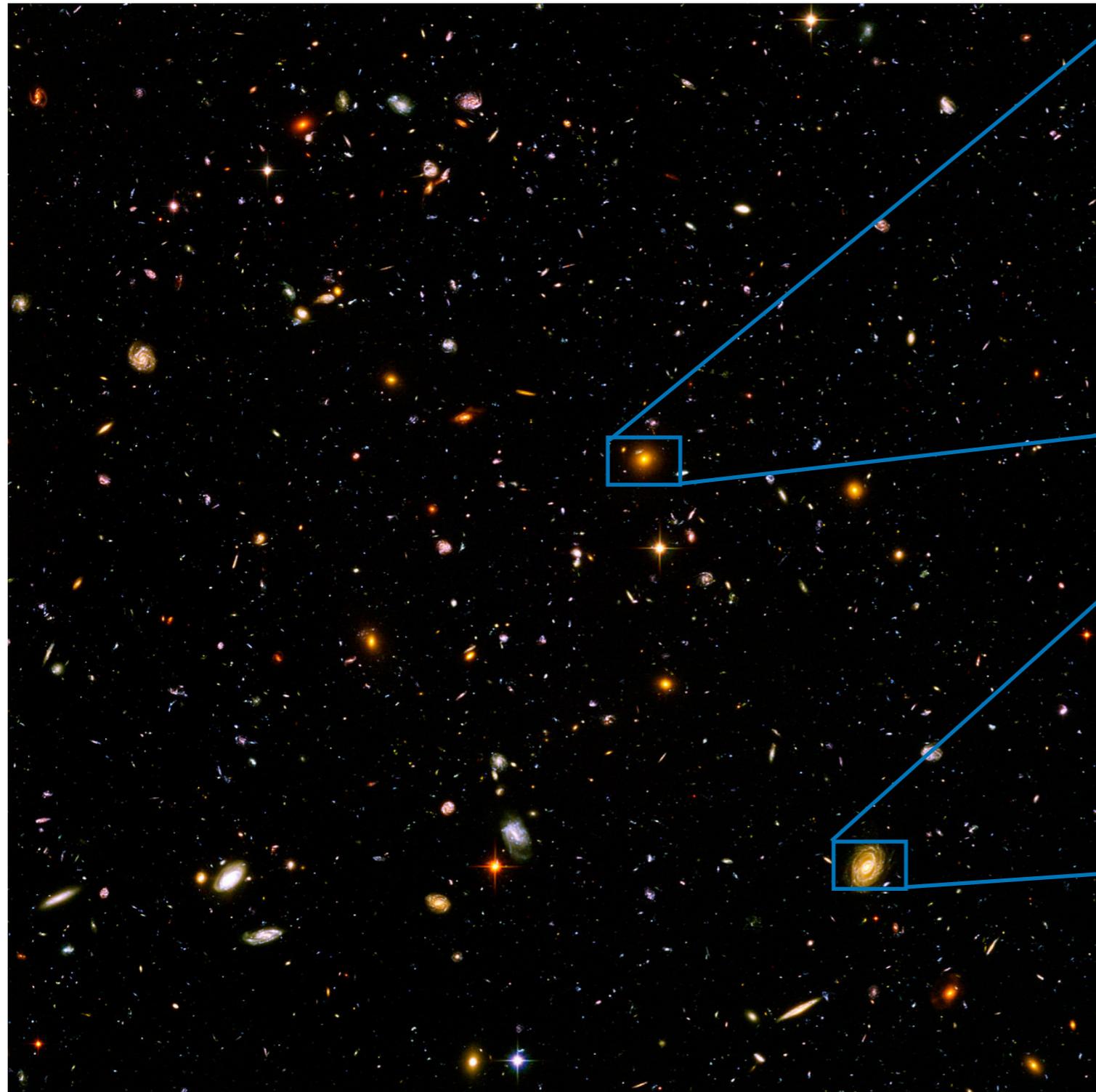
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IIT Hyderabad (15th July 2020)

Collaborators- Tae-hyeon Shin, Ethan Nadler, Arka Banerjee, Eric Baxter, Chihway Chang, Neal Dalal, Bhuvnesh Jain, Andrey Kravtsov, Jeremy Sakstein, Risa Wechsler

Image of the night sky taken with the Hubble Space telescope



clusters of bound galaxies



**Galaxies are formed by baryonic matter
and are held together by gravity and
hydrodynamic forces**

Evidence for a dark component to gravity

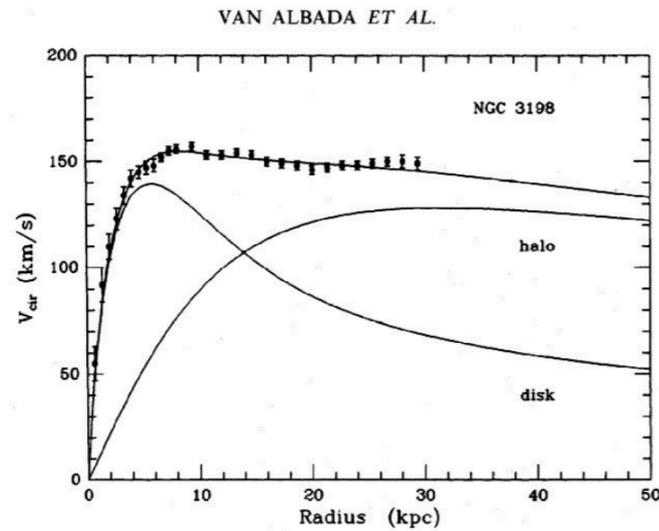
Galaxy rotation curves

ROTATION OF THE ANDROMEDA NEBULA FROM A SPECTROSCOPIC SURVEY OF EMISSION REGIONS*

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Department of Terrestrial Magnetism, Carnegie Institution of Washington and Lowell Observatory, and Kitt Peak National Observatory‡

Received 1969 July 7; revised 1969 August 21



Existence of thin disks



Velocity dispersion of Coma cluster -

Fritz Zwicky - 1933

Velocity dispersion was not consistent with viral theorem.

A NUMERICAL STUDY OF THE STABILITY OF FLATTENED GALAXIES: OR, CAN COLD GALAXIES SURVIVE?*

J. P. OSTRIKER

Princeton University Observatory

AND

P. J. E. PEEBLES

Joseph Henry Laboratories, Princeton University

Received 1973 May 29

$$Q \equiv \frac{\sigma \kappa}{3.36 G \Sigma_0} \geq 1.$$

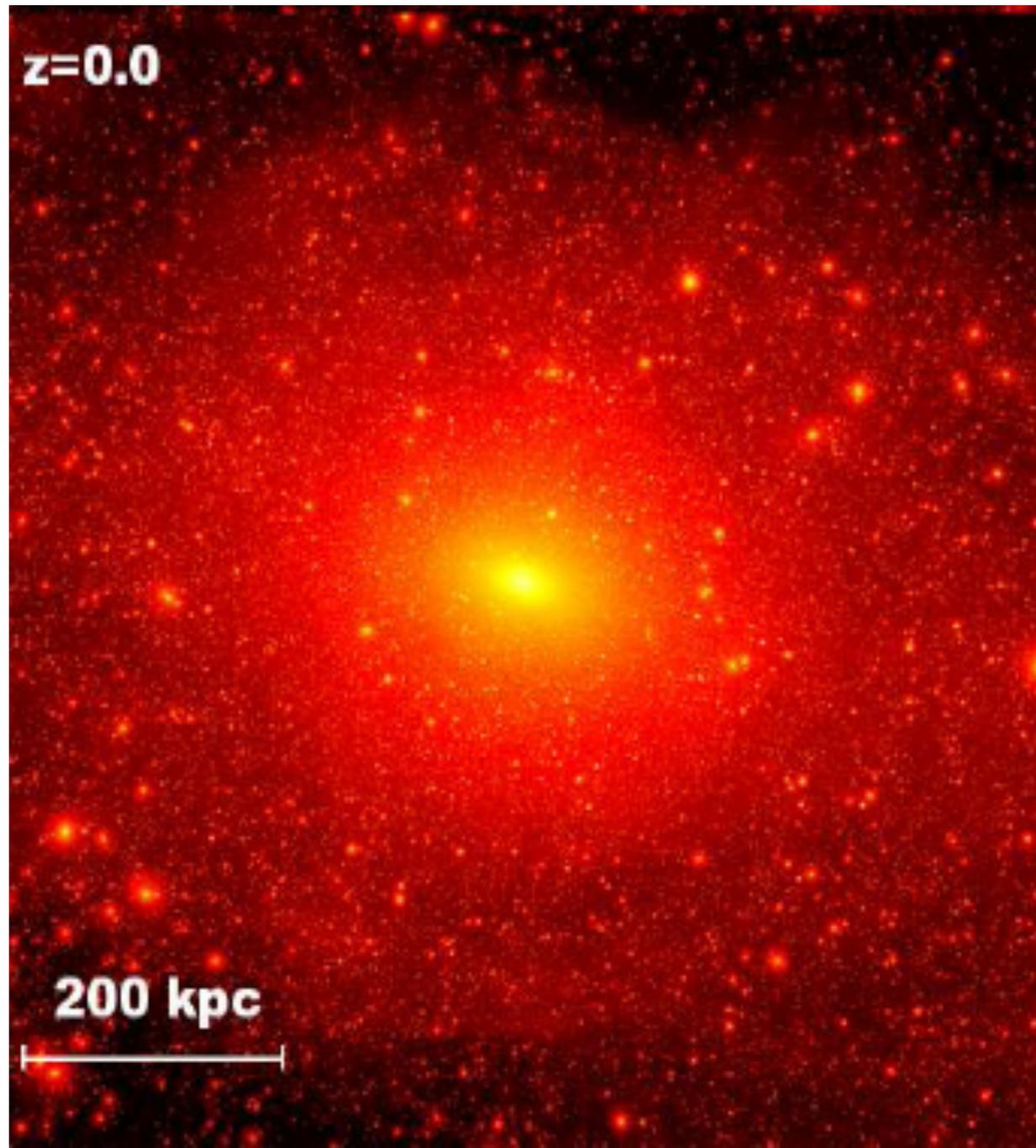
Direct evidence for the existence of dark matter



Image composite credit: X-ray: NASA / CXC / CfA / M.Markevitch et al.; Optical: NASA / STScI; Magellan / U.Arizona / D.Clowe et al.; Lensing Map: NASA / STScI; ESO WFI; Magellan / U.Arizona / D.Clowe et al.

Merging clusters -The bullet cluster system

What are Dark Matter Halos ?

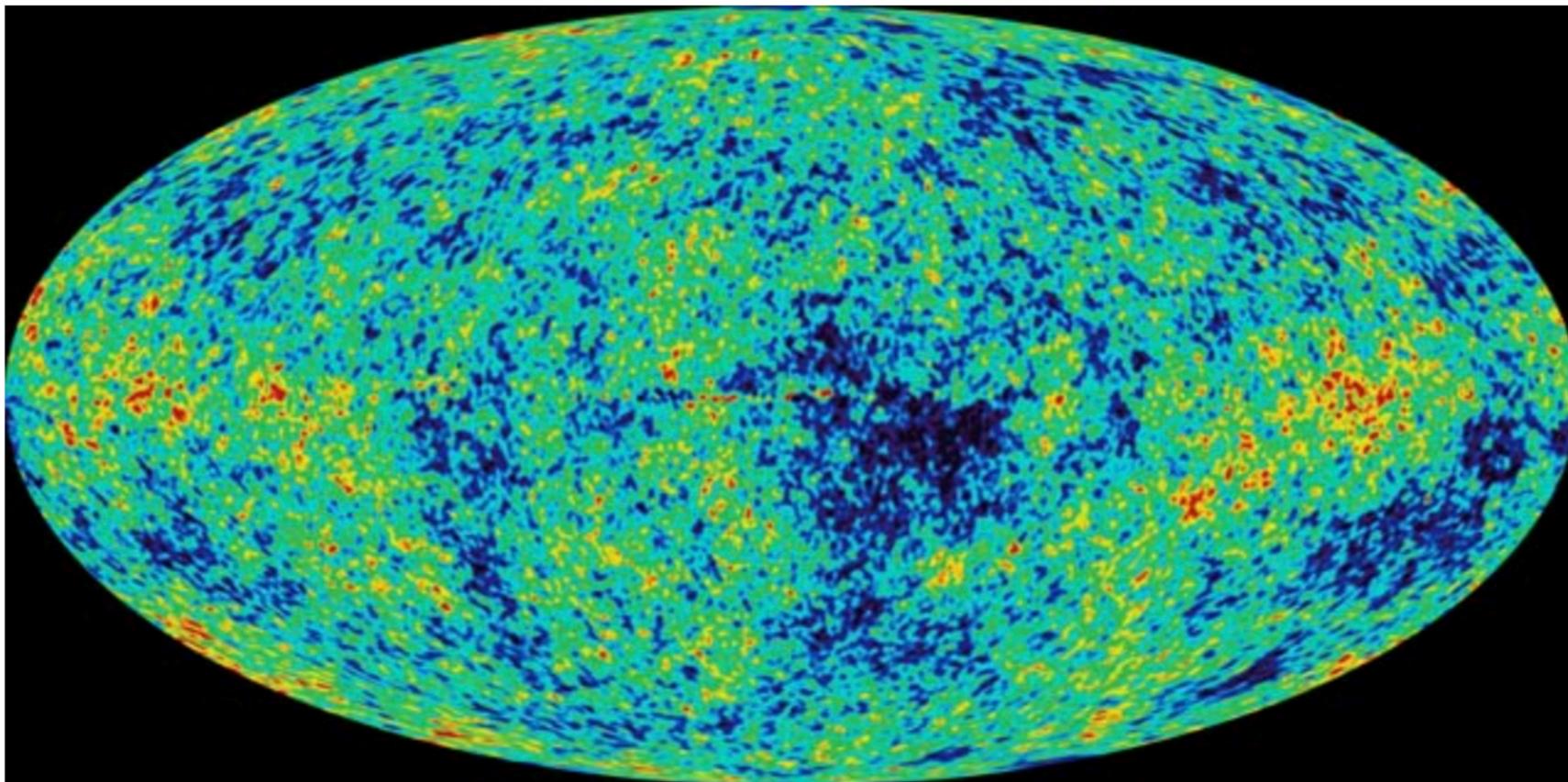


- Dark matter halos are endpoints of all cosmological structure formation
- Self-bound, virialized structures
- Harbor all stars, galaxies, quasars

Via Lactea simulation

Structure formation in the universe

Initial quantum fluctuations in the density of matter magnified by inflation



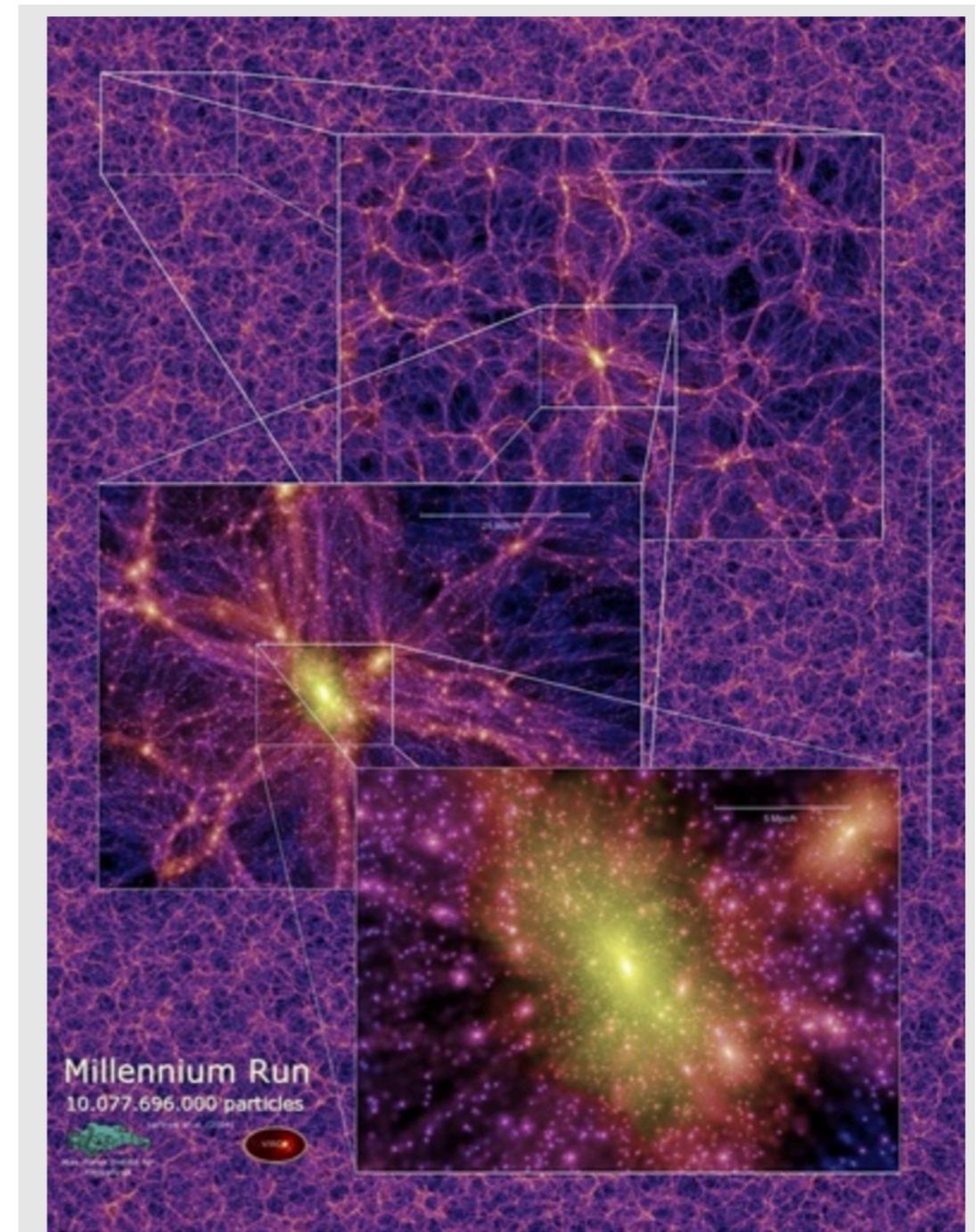
The cosmic microwave background

Density perturbations collapse gravitationally to form dark matter halos

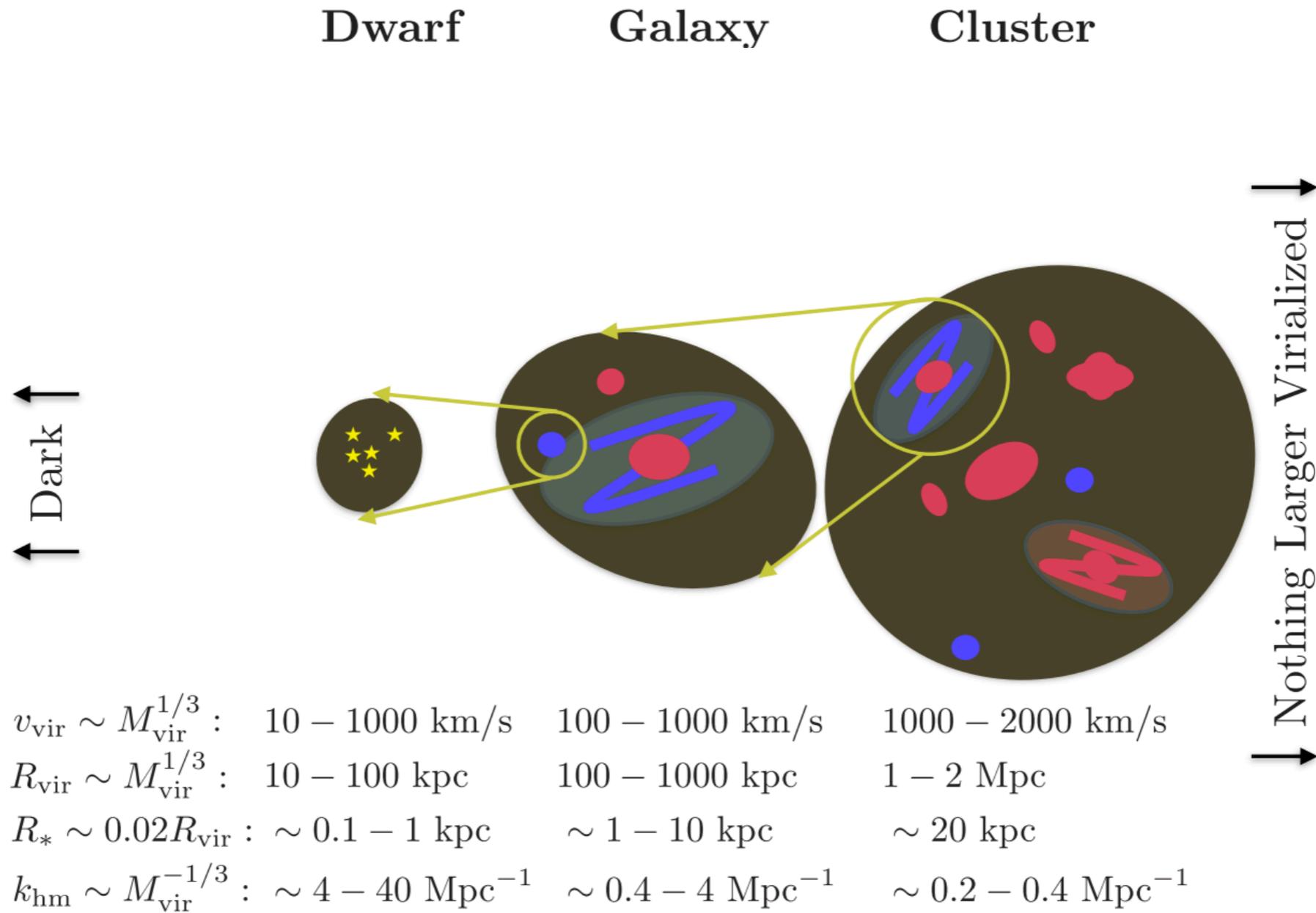
Structure formation in the universe



Density perturbations collapse gravitationally to form halos



Small halos form first and merge to form more massive halos

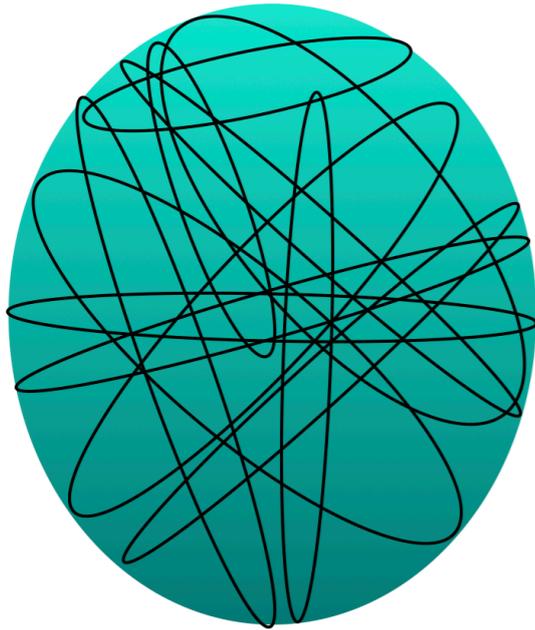


credit: Buckley and Peter 2017

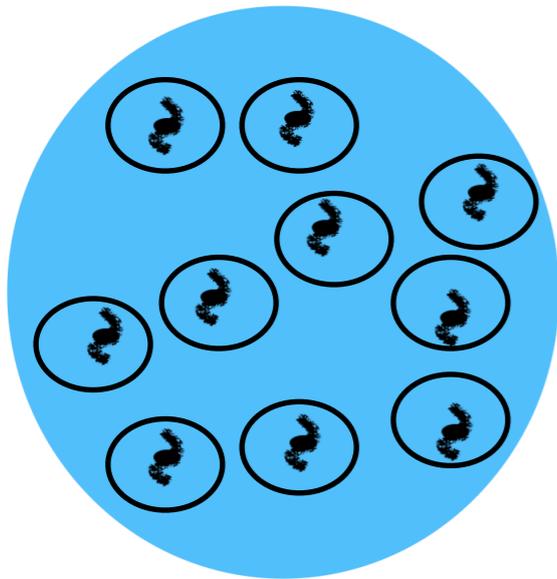
Hierarchical structure formation

Hierarchical structure formation

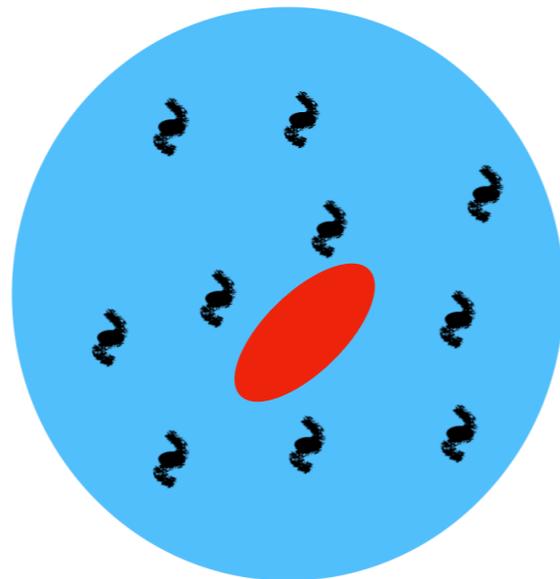
Main components of a halo



Dark matter **particles** that are orbiting in a central potential

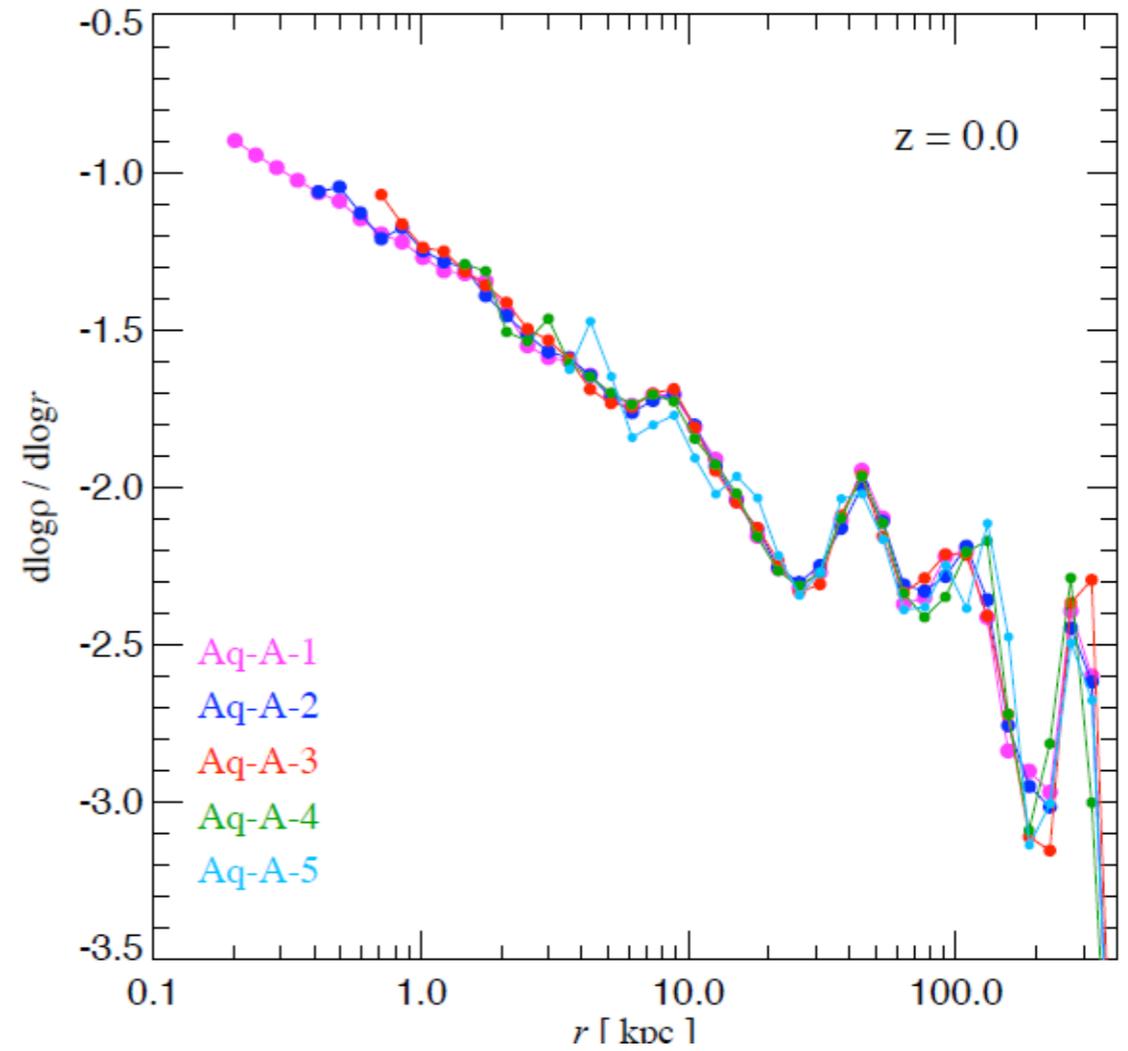
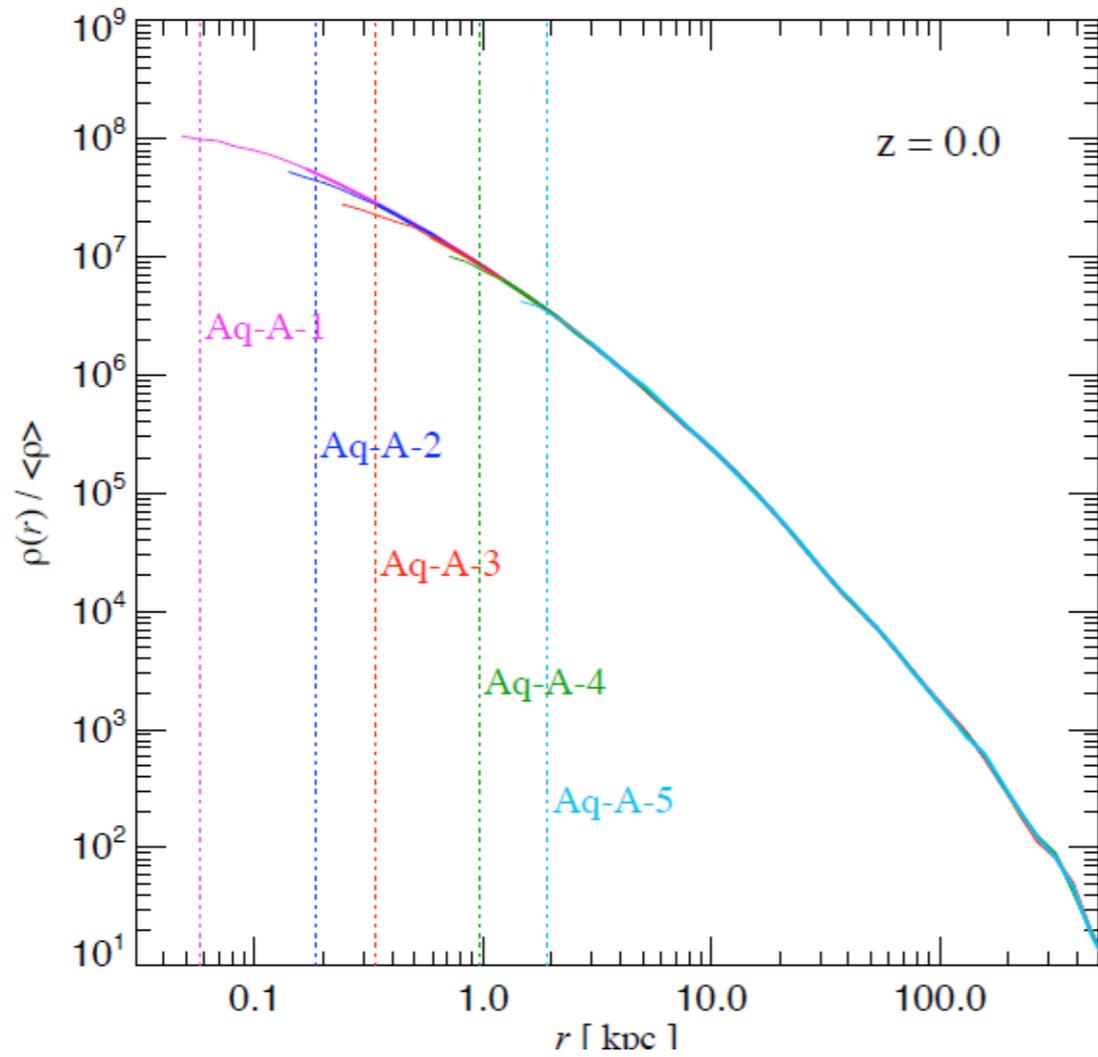


Halos grow hierarchically -
small objects form first and fall into massive halos
So halos contain **subhalos** that also harbor **galaxies**



Baryonic matter in the form of **diffuse stars,**
gas and galaxies

The density profiles of dark matter halos



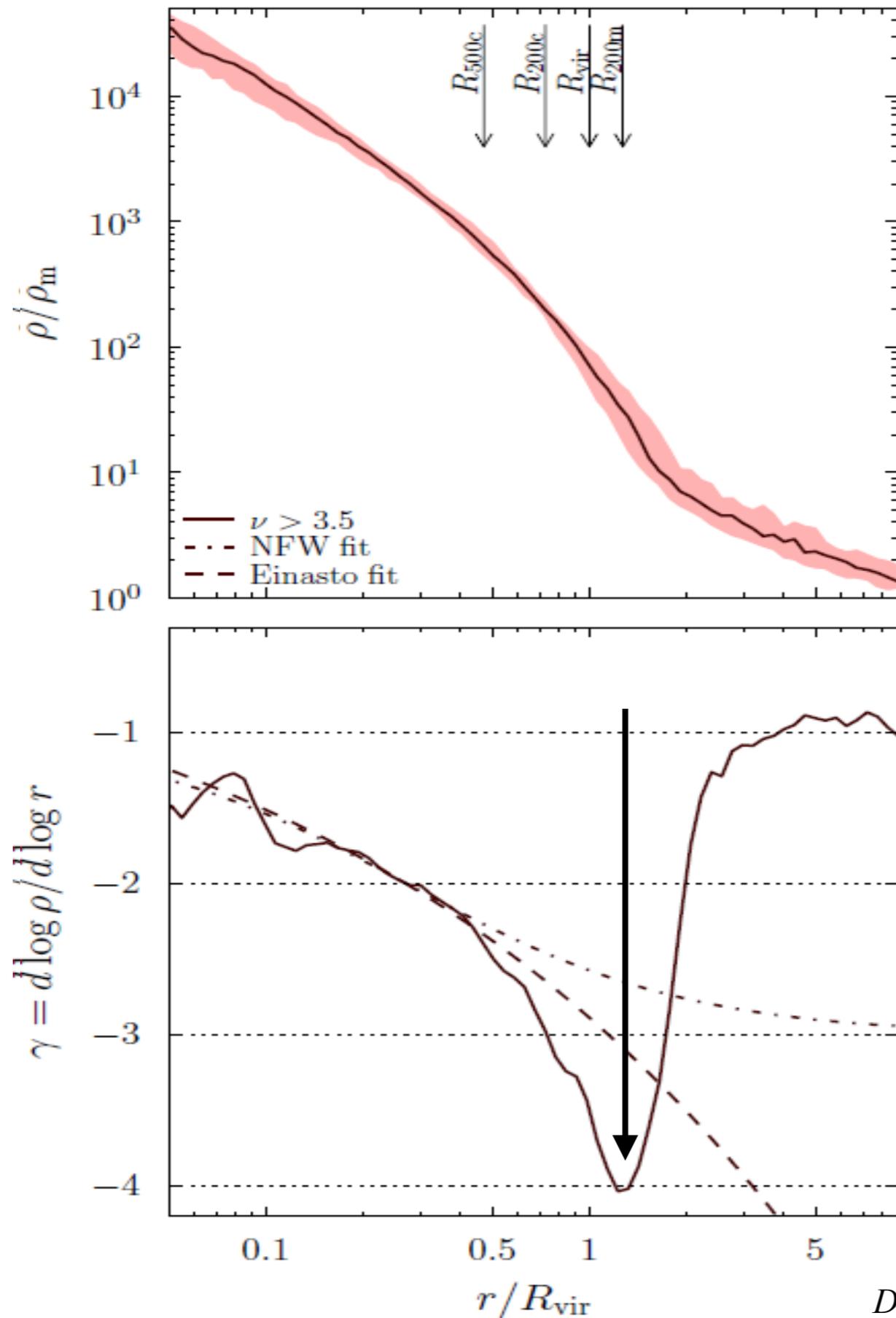
“Aquarius” Springel et. al 2008

- The density of halos is well described by the NFW profiles

$$\frac{\rho_0}{\frac{R}{R_s} \left(1 + \frac{R}{R_s} \right)^2}$$

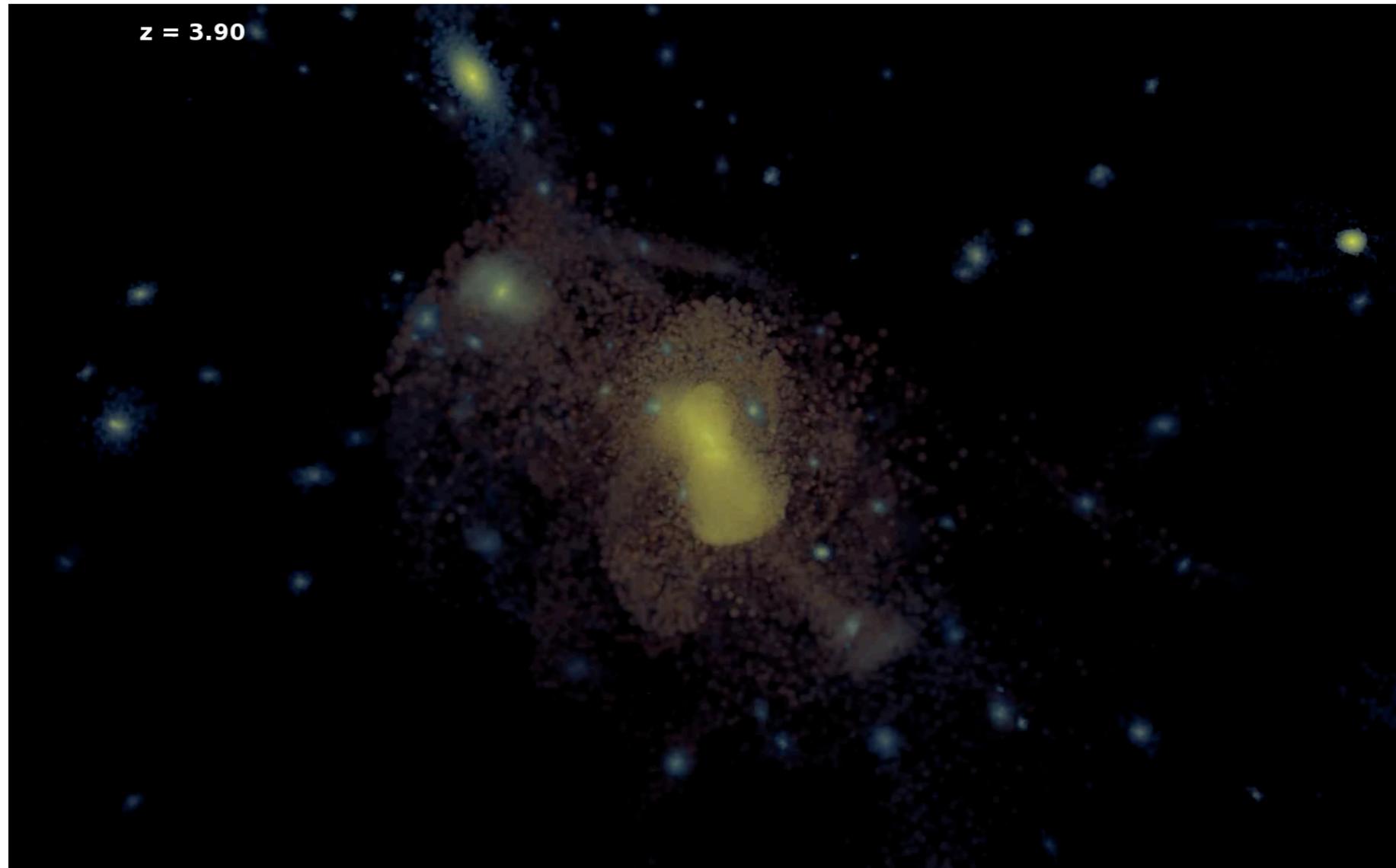
- Slope is -1 in the inner regions and rolls over to -3 in the outskirts of the halo.

Outer density profiles of Dark Matter Halos



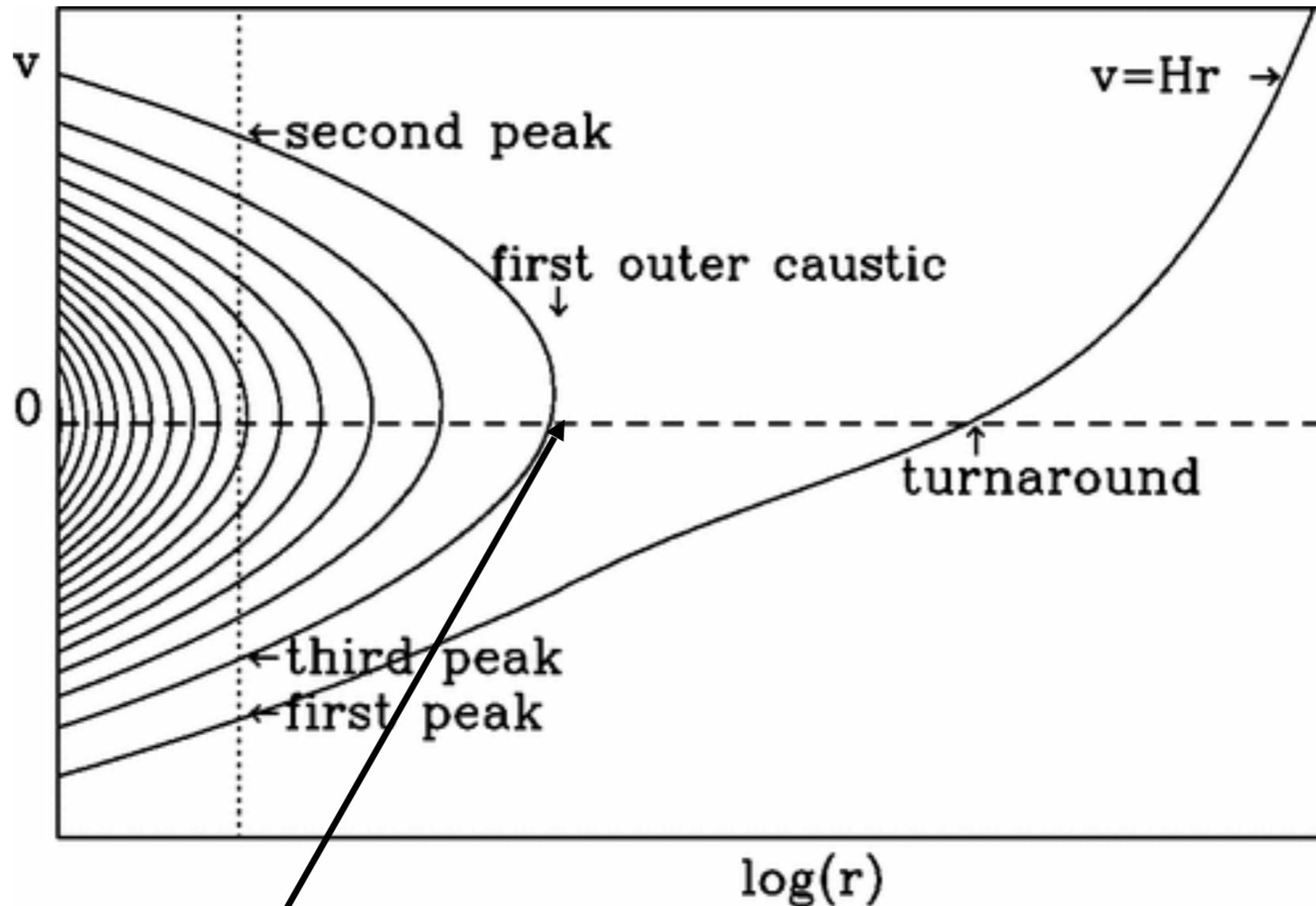
- Deviation from NFW and Einasto profile in the outer regions of the halo
- Slope of the local density deviates in a narrow confined region

The evolution of dark matter halos



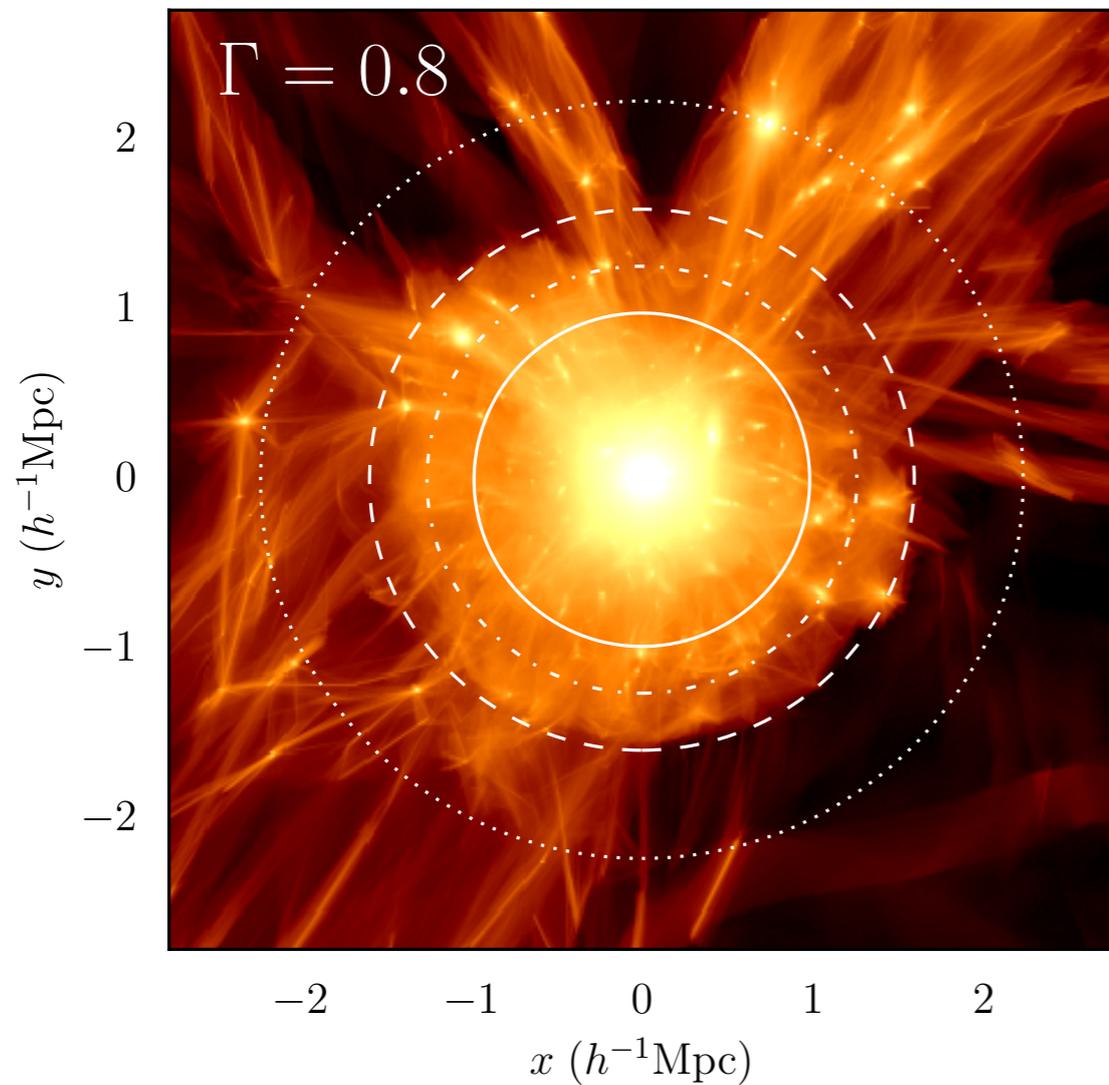
Phase space Diagram of Halo evolution

For spherical potential and smooth accretion

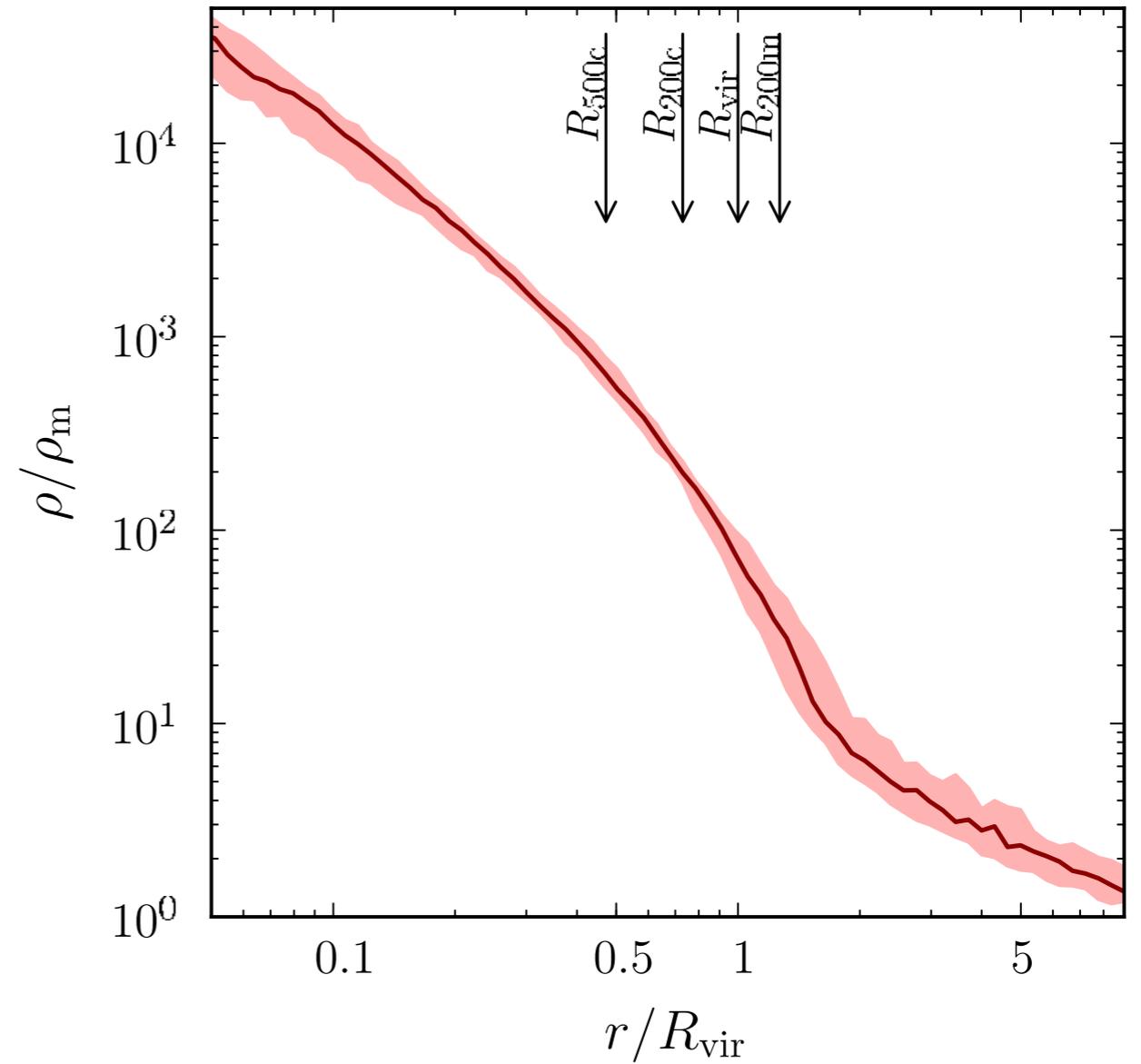


Splashback - corresponds to first apoapses passage after collapse

Where is the boundary of a halo?

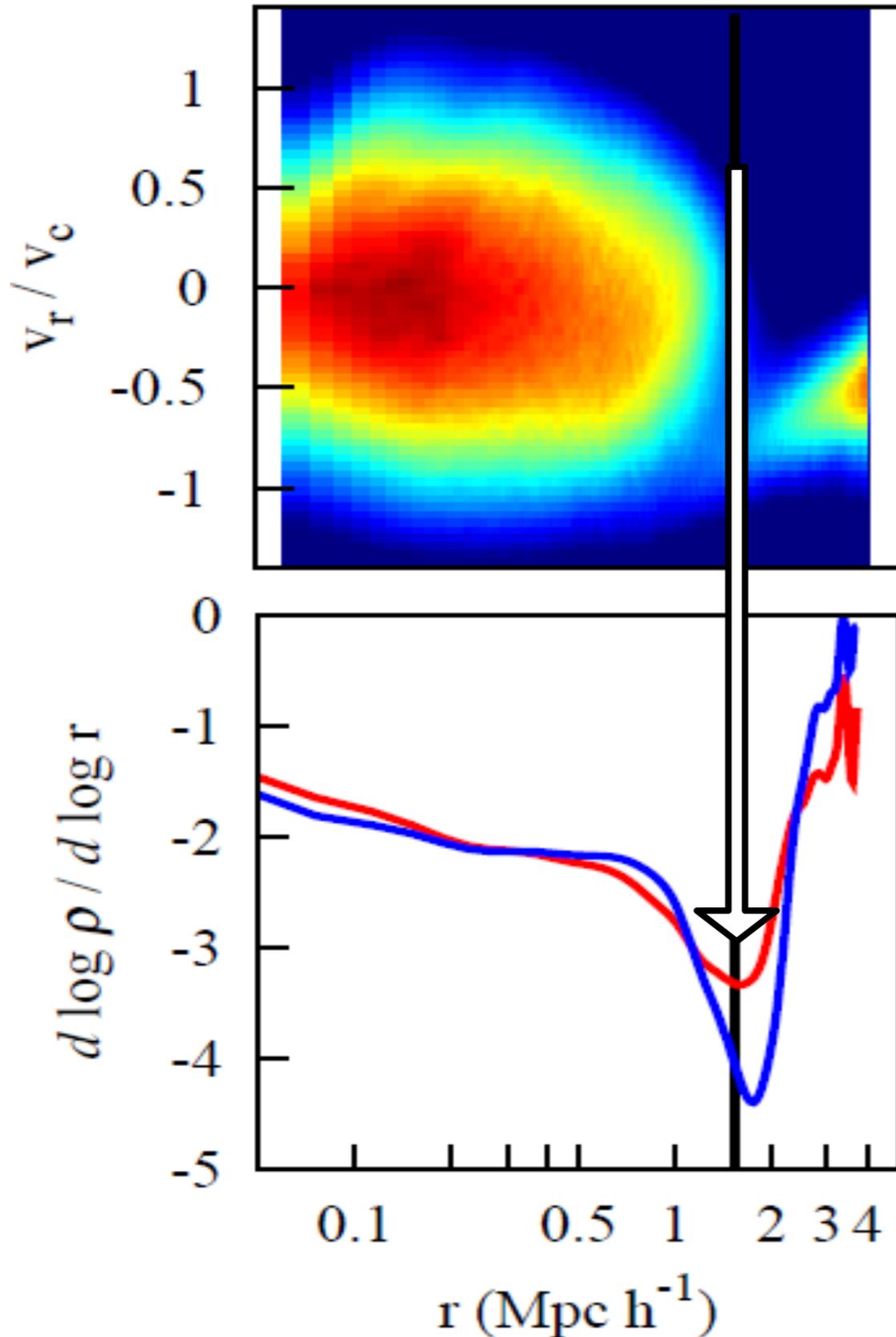


More et al. 2015



Diemer & Kravtsov 2014

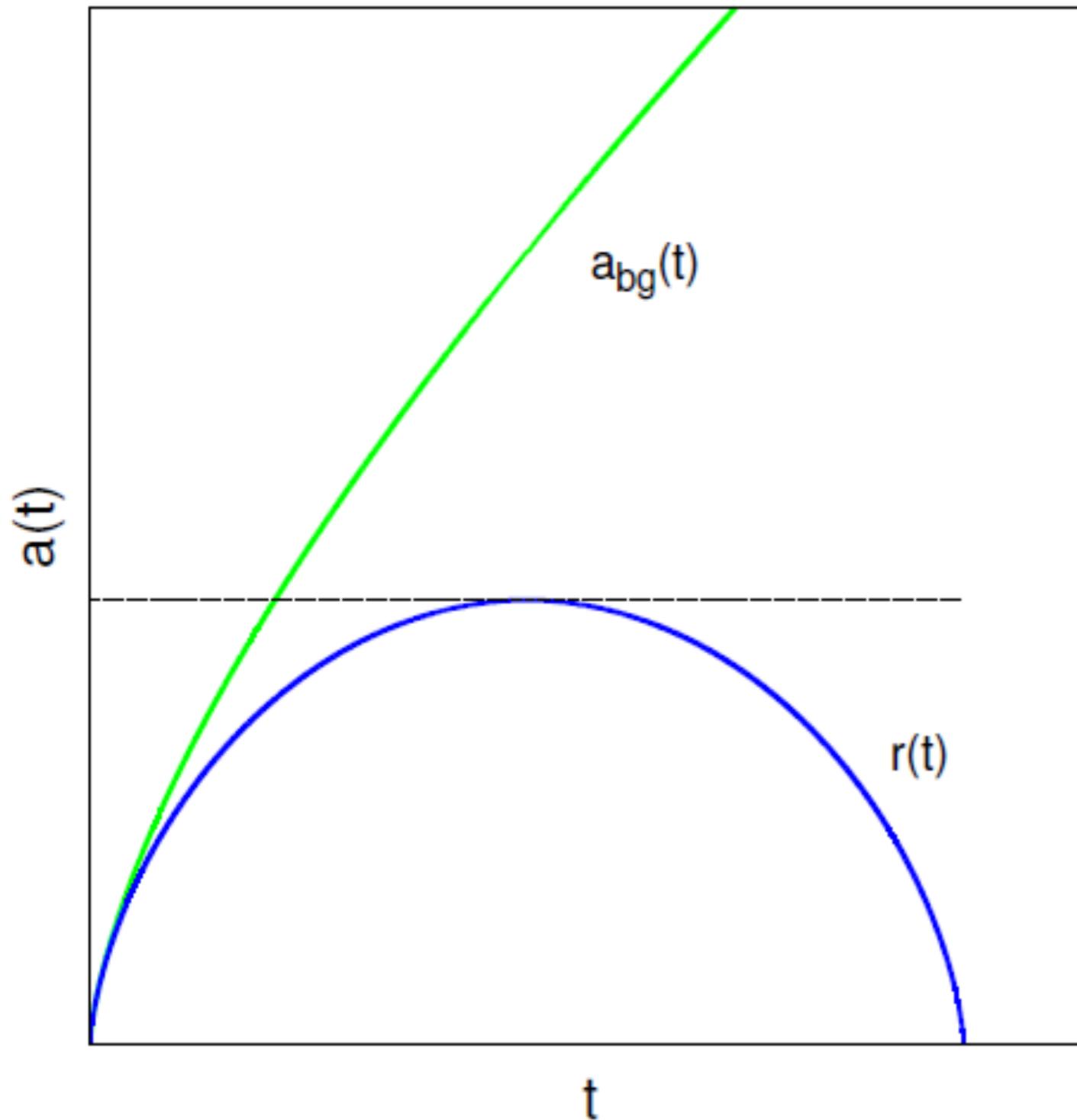
$$M_{\Delta} = \frac{4}{3}\pi R_{\Delta}^3 \Delta\rho_{\text{ref}}$$



- Phase space diagram of N-body halos from the Multidark simulation
- Halos stacked in the mass range of $1-4e14$ Msun
- Position of splashback coincides exactly with feature

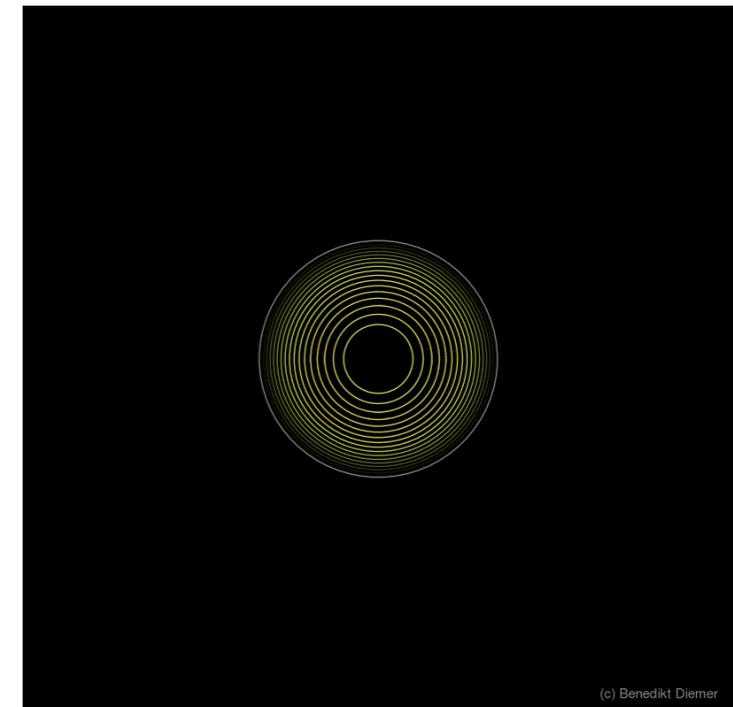
Phase space boundary at the location of turnaround of the most recently accreted material

Collapsing shells of matter around a dark matter over density

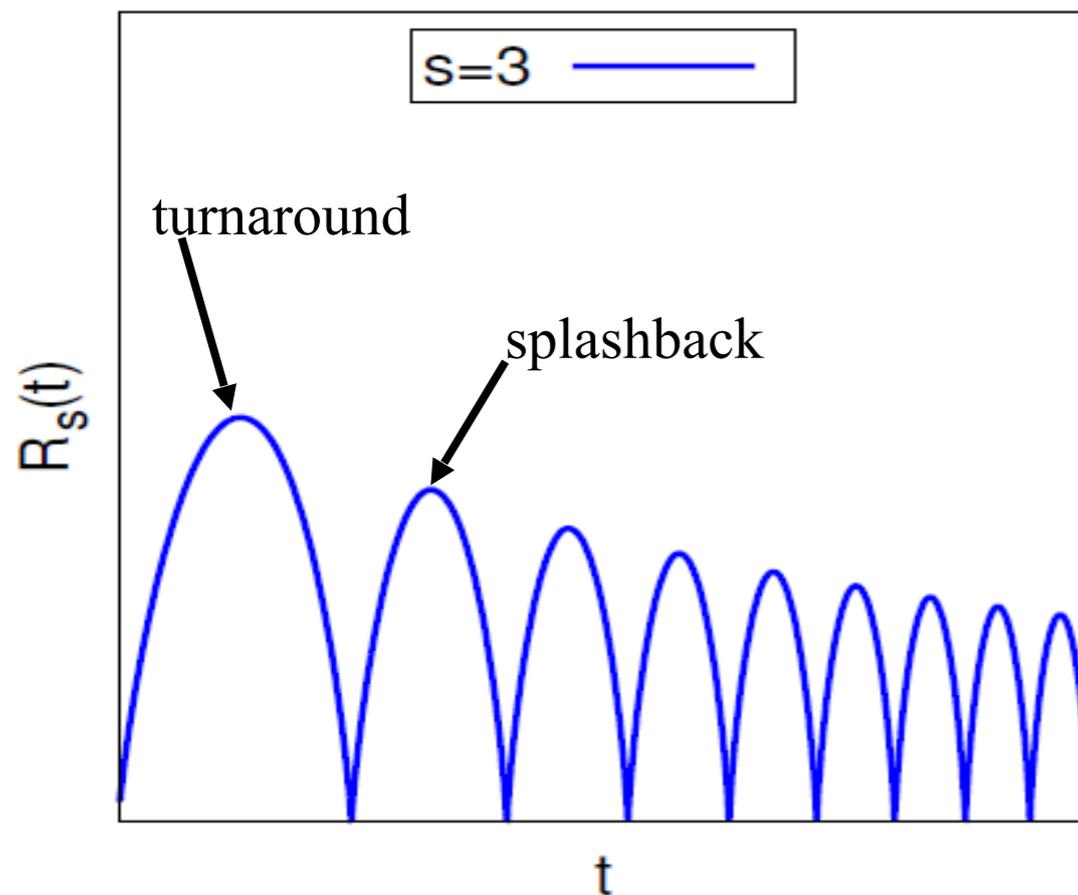
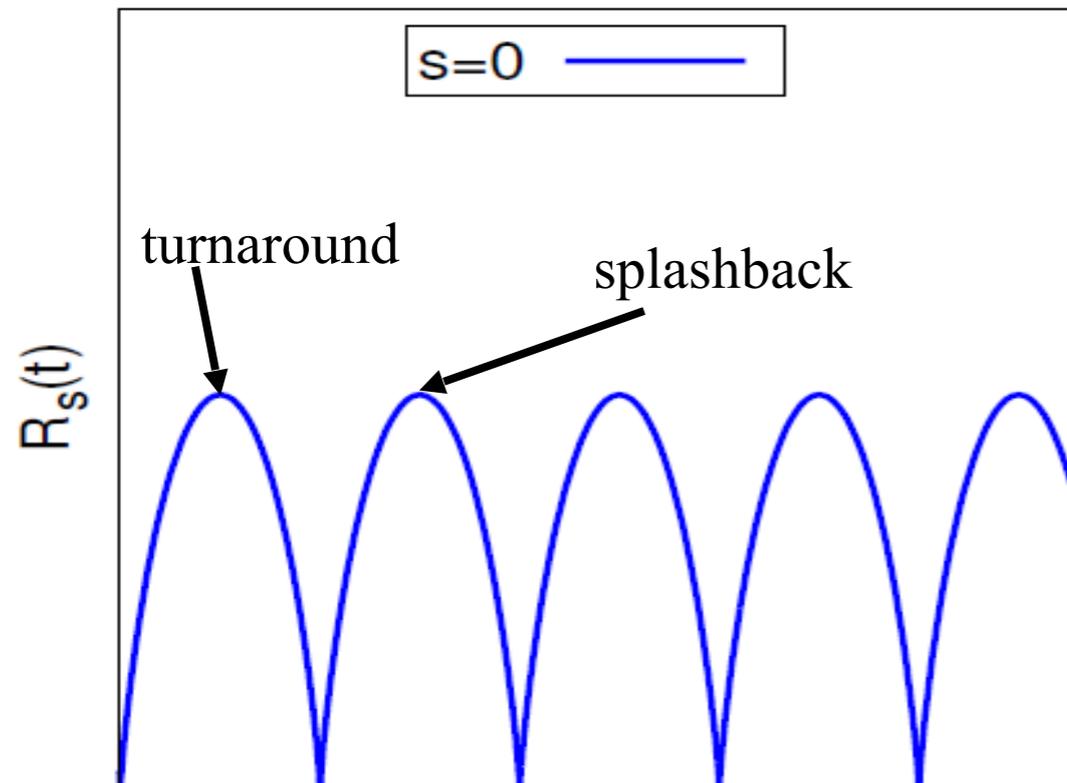


$$\frac{d^2 r}{dt^2} = -\frac{GM(< r)}{r^2}$$

$$\frac{d^2 r}{dt^2} = -\frac{GM(< r)}{r^2} + \frac{\Lambda c^2}{3} r$$



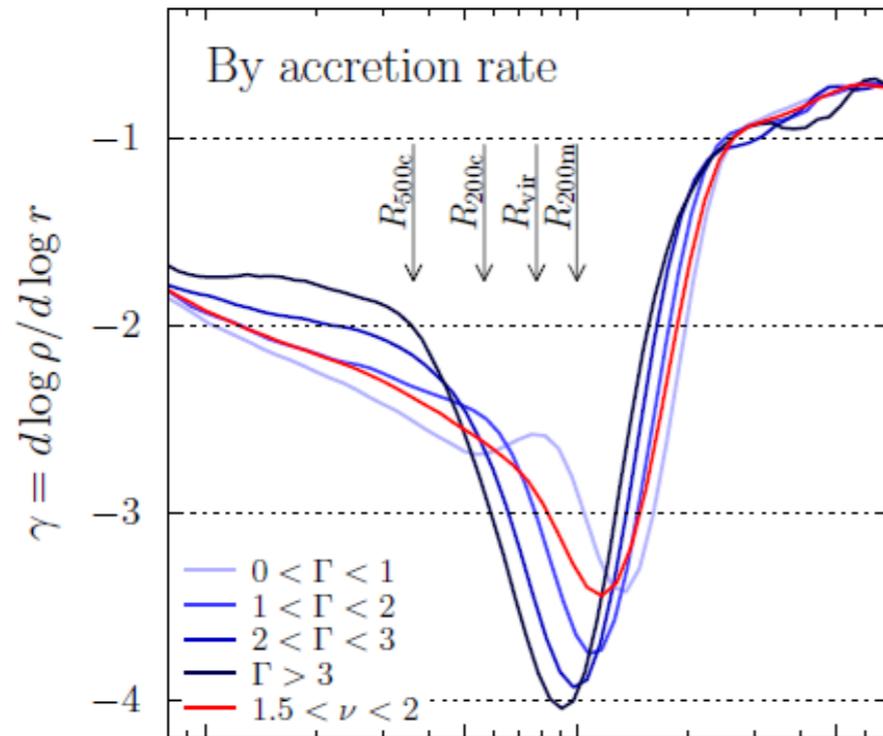
Particle Orbits



$$\frac{d^2 r}{dt^2} = -\frac{GM(< r, t)}{r^2} + \frac{\Lambda c^2}{3} r$$

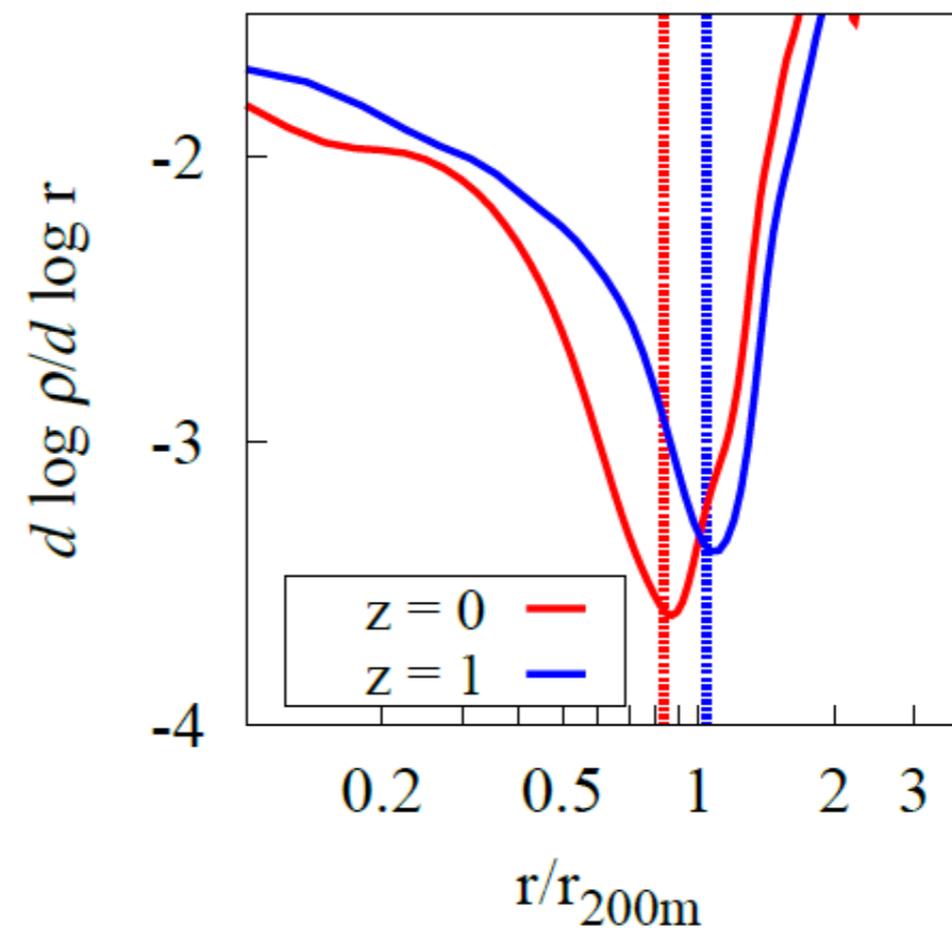
- For a constant potential the subsequent orbits are exactly the same
- Mass accretion - potential becomes deeper with time - Subsequent orbits shrink and become faster

Function of Accretion Rate and halo redshift

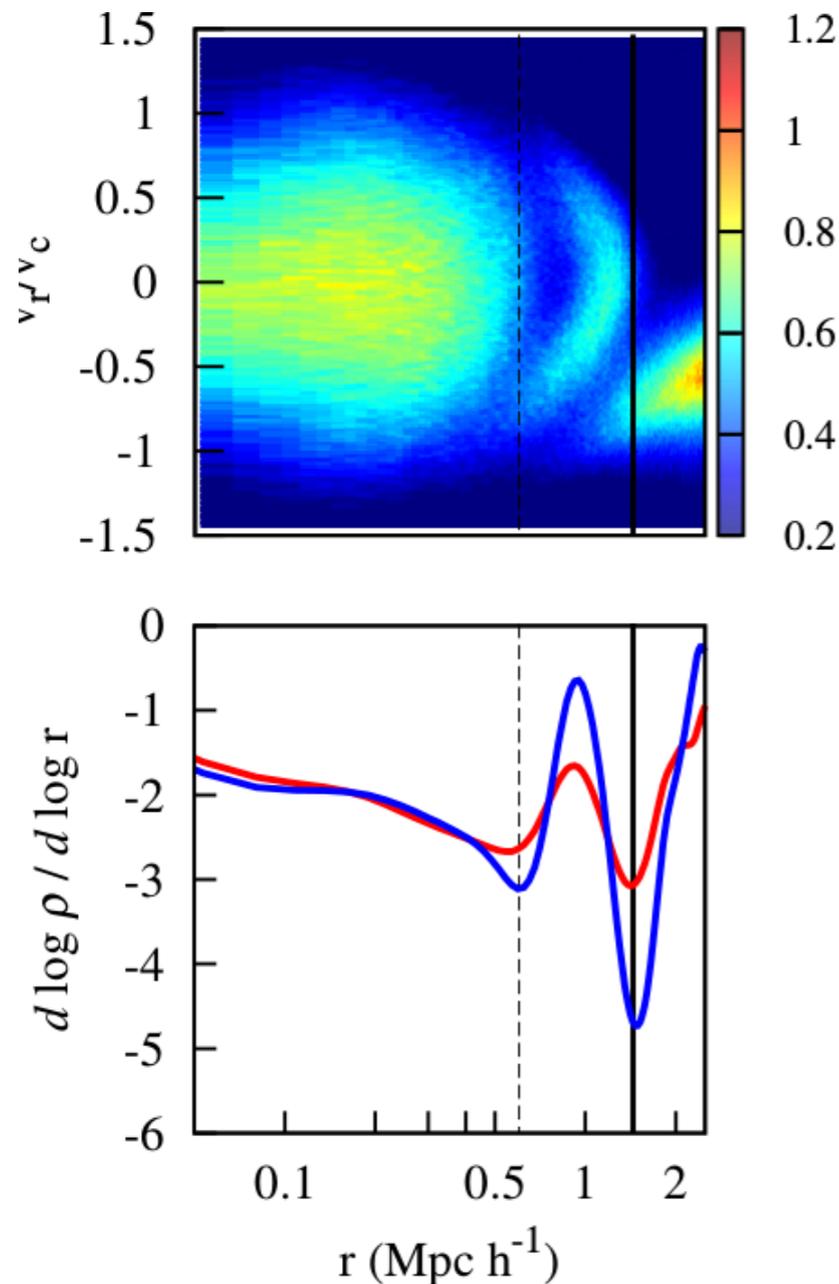


Faster a halo grows, the smaller is its splashback radius in units of R_{200} .

At a given accretion rate it is a function of redshift

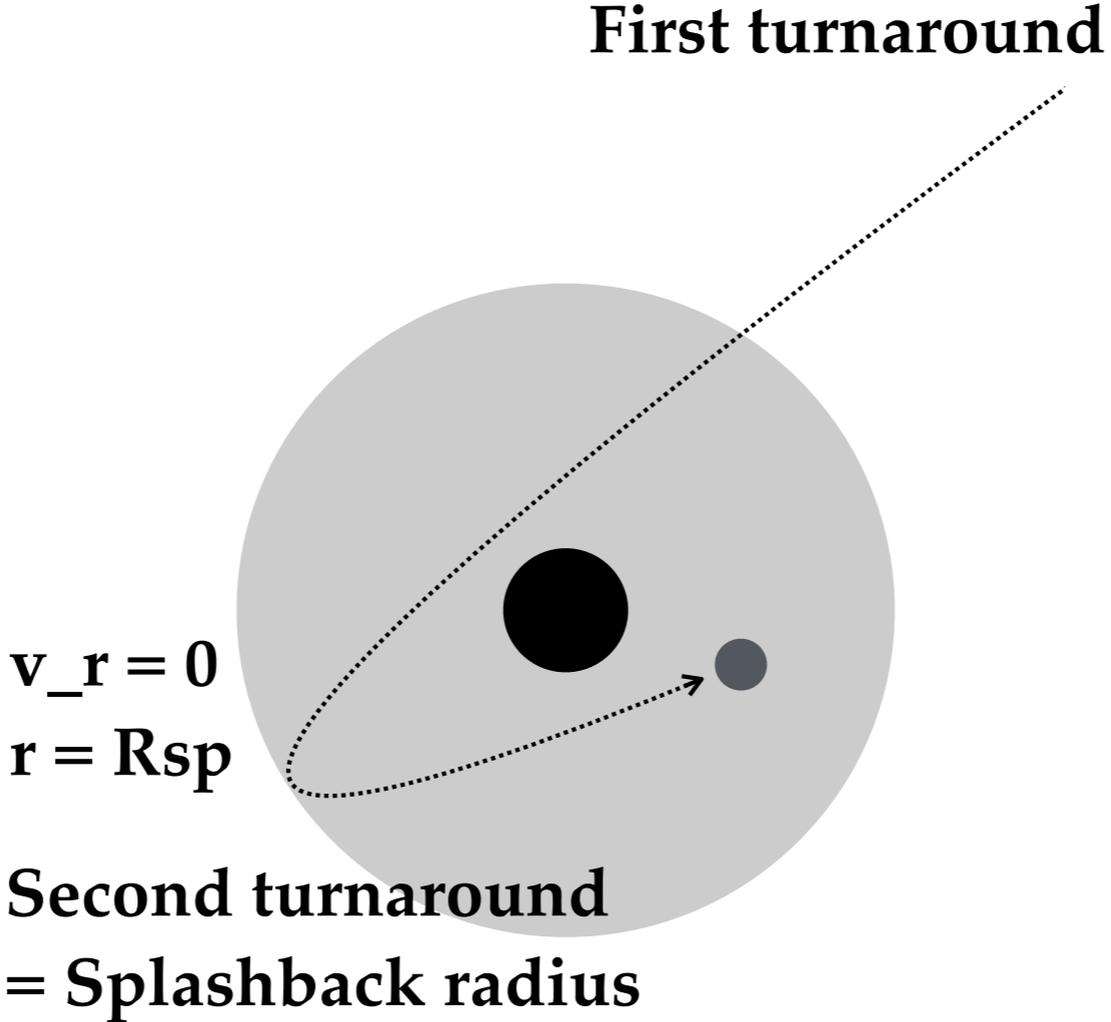


Why is this feature interesting?



- It forms the boundary of the halo
- Physical definition of halo mass
- The splashback radius probes growth history of the halo.
- It forms at the boundary that separates the virialized region of a halo from the infalling region.
- Fundamental length scale in the halo structure, should be present if there is a dark matter halo.
- Simple to understand formed **by the most recently accreted material** - that is not yet phase mixed.
- Inner regions of halos are often dominated by baryons

The location of the splashback radius is set by **simple physical principles** -
Gravitational collapse of cold dark matter in an expanding universe.



Credit : Chihway Chang

Gravitational collapse of collisionless dark matter in an expanding universe

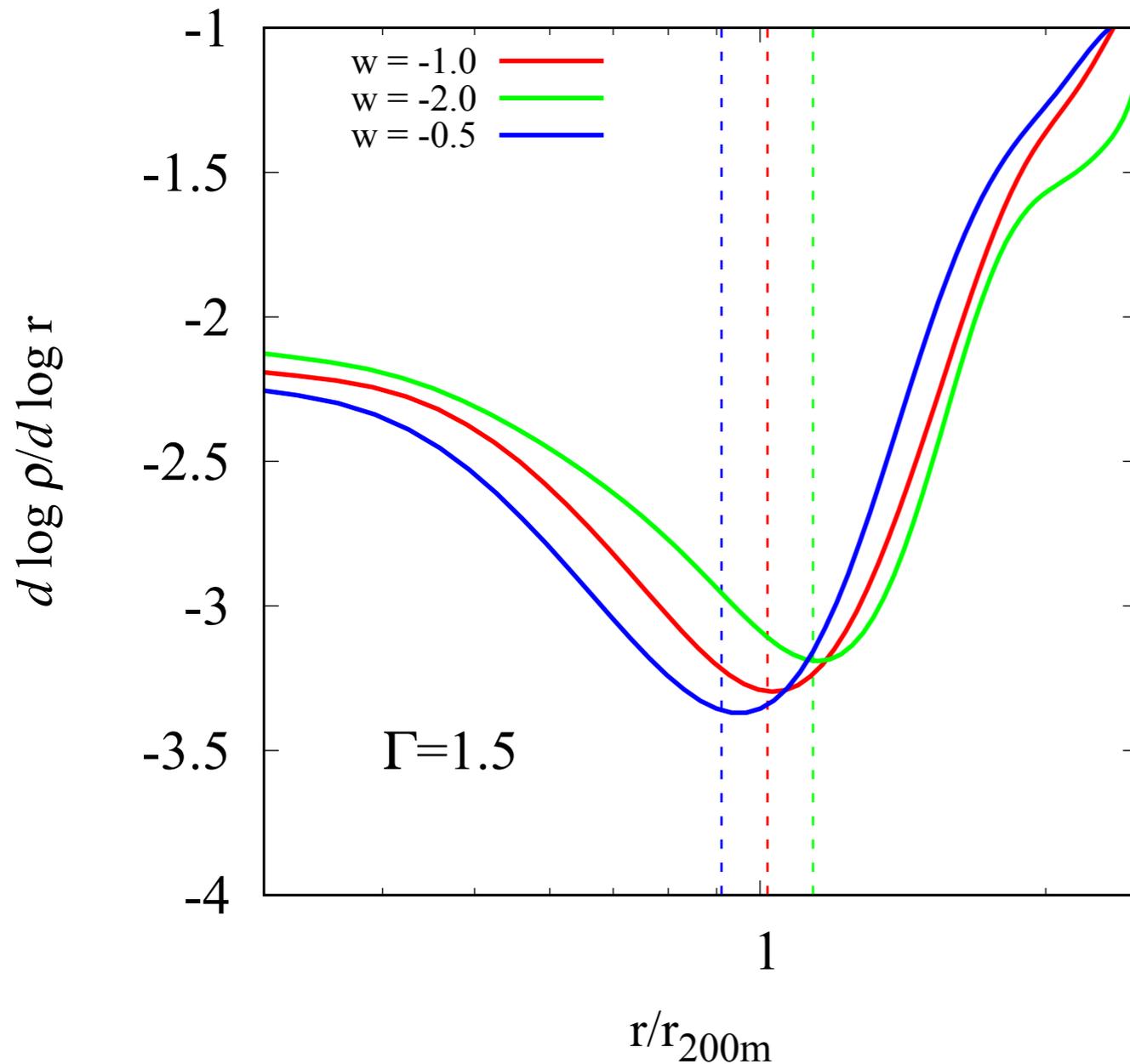
Gravitational collapse of collisionless dark matter in an expanding universe

What happens if we change gravity?

If dark matter self-interacts?

If universe is not Lambda CDM?

What happens to splashback if you change the equation of state parameter?



Adhikari et al. 2018

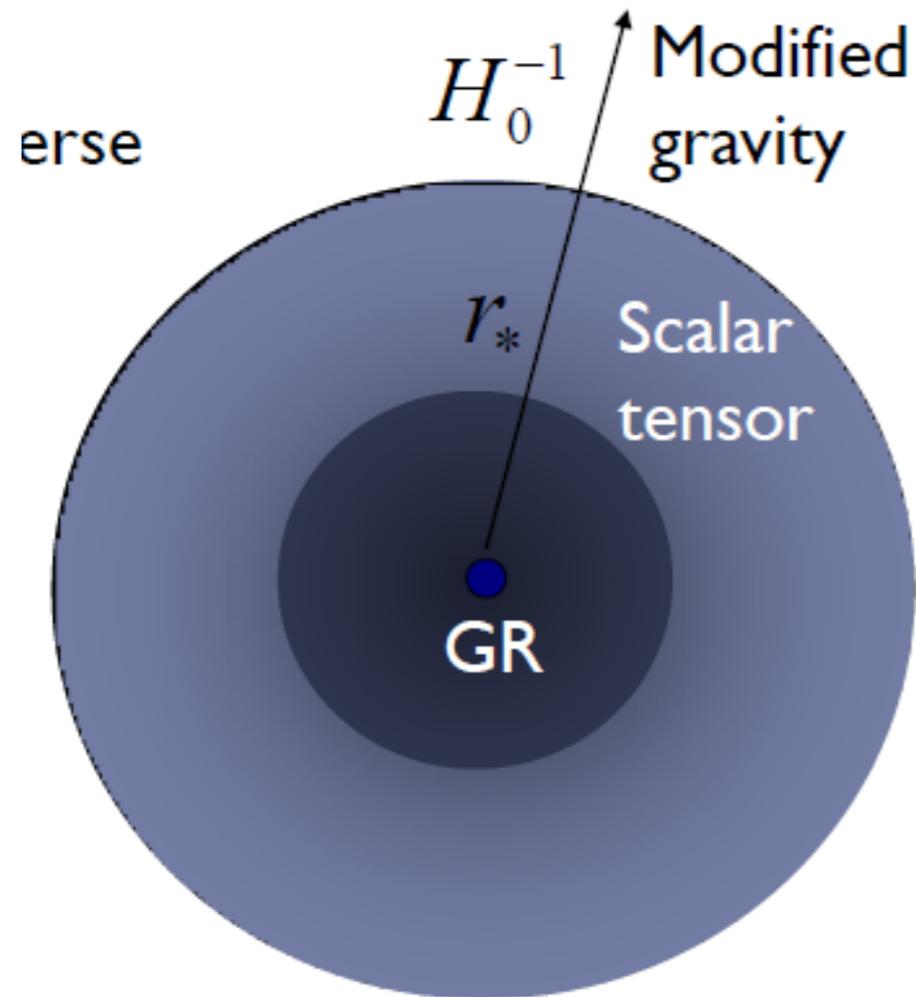
$$P_{DE} = w\rho \quad w \neq -1$$

$$\frac{\ddot{a}}{a} = H_0 \sqrt{\Omega_m a^{-3} + \Omega_{DE} a^{-3(1+w)}}$$

$$\ddot{r} = -\frac{GM}{r^2} - \frac{H_0^2}{2} \Omega_{DE} (1 + 3w) r^{-2-3w}$$

Splashback is a weak function of the w

What happens if we change gravity?



Large scales - Gravity is modified so that the universe accelerates

Intermediate scales - Gravity is still modified by a fifth force

Small scales - Solar system tests constrain gravity to normal GR

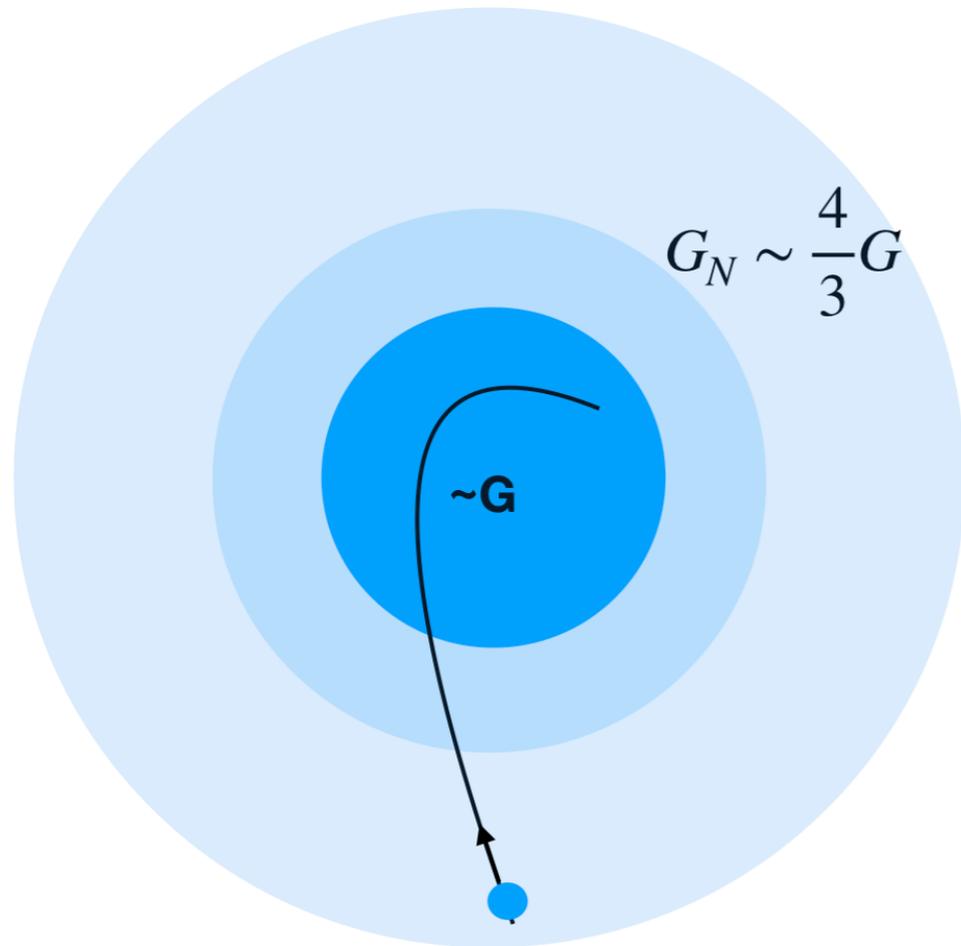
Screening mechanism :

Chameleon screening. - Mass of scalar mode becomes large in dense regions (f(R))

Vainshtein screening - non-linear derivative of fifth force becomes large in dense regions (DGP)

What happens if we change gravity?

Does the location of splashback radius change in modified gravity?



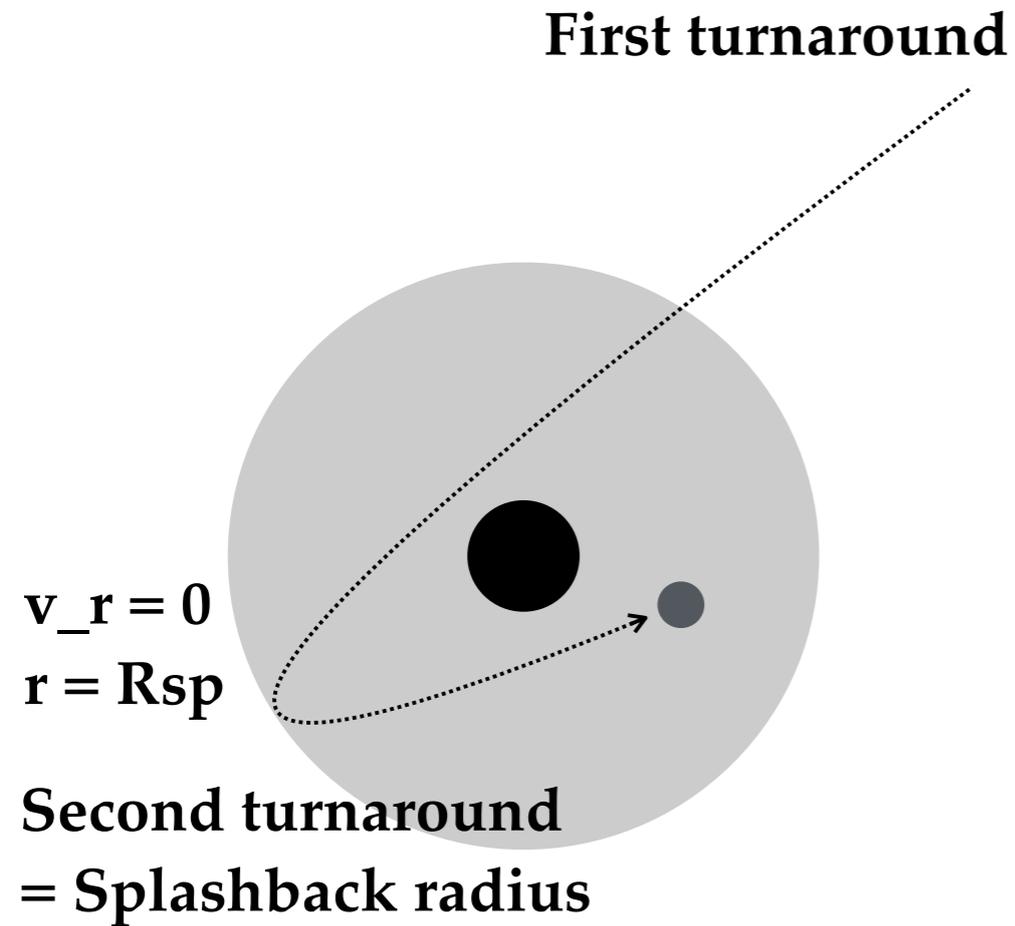
$$\ddot{r} = -\frac{GM}{r^2} + \frac{\Lambda r}{3} + F_{\pi}$$

$$F_{\pi} = -\alpha\pi' = -4\alpha^2 \frac{GM}{r^2} g\left(\frac{r}{r_*}\right)$$

$$g(\zeta) = \zeta^3 \left(\sqrt{1 + \zeta^{-3}} - 1 \right)$$

- i) Extra force mediated by the scalar field**
- ii) The enhanced gravity in the outskirts makes infall velocity higher.**

Splashback of Substructure in modified gravity

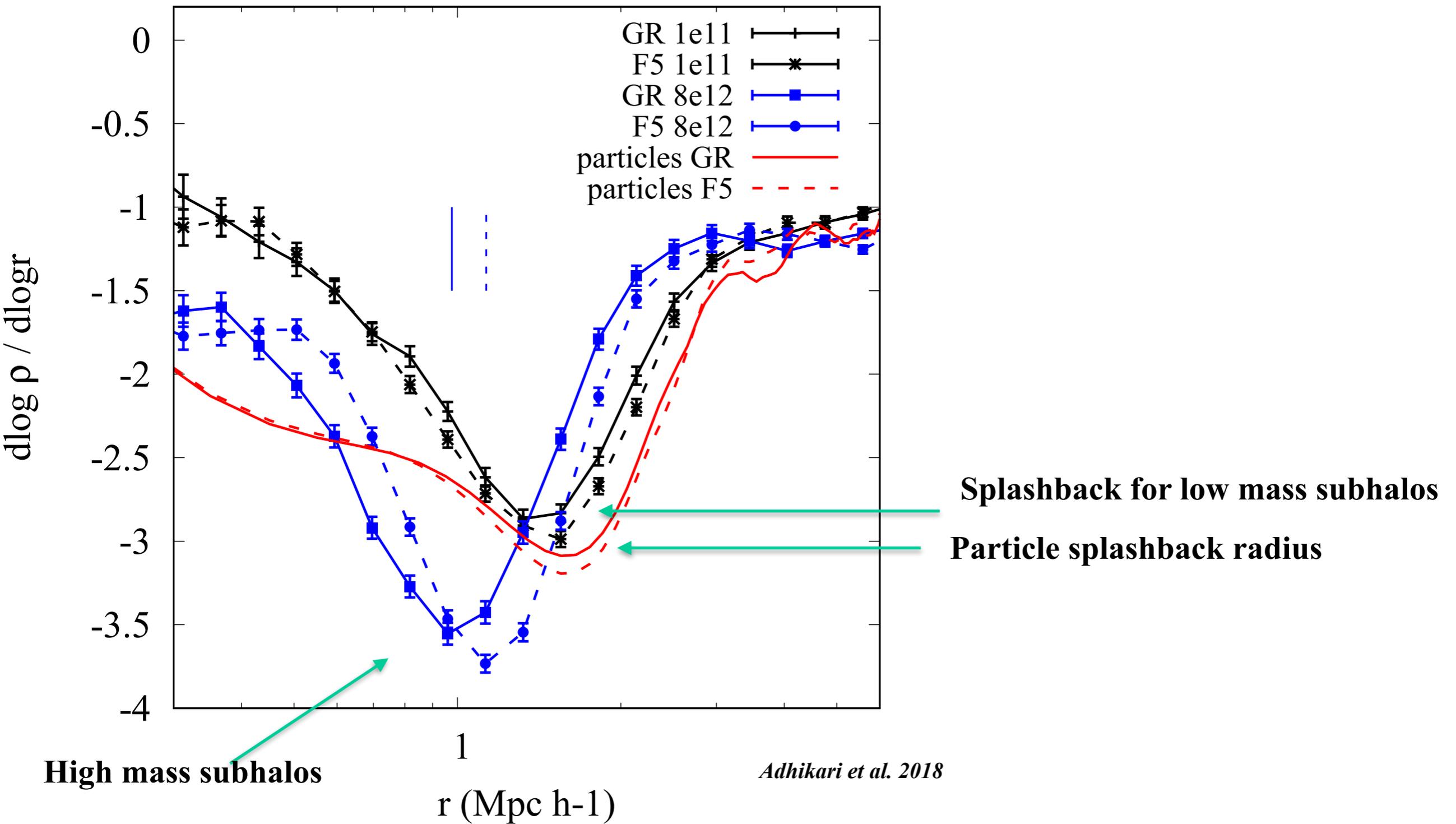


What happens to the subhalos?

Dynamical friction in subhalos

$$\frac{dv}{dt} \propto -\frac{G^2 M \rho}{v^3} v f(v, \sigma)$$

Faster a massive object moves, lower is the force of friction



High mass subhalos in feel lesser amount of dynamical friction in modified gravity - splashback at larger radius than their counterparts in GR

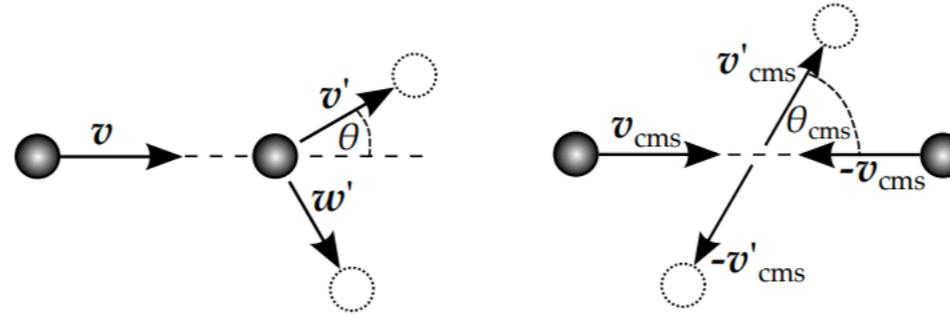
Gravitational collapse of collisionless dark matter in an expanding universe

What happens if we change gravity?

If dark matter self-interacts?

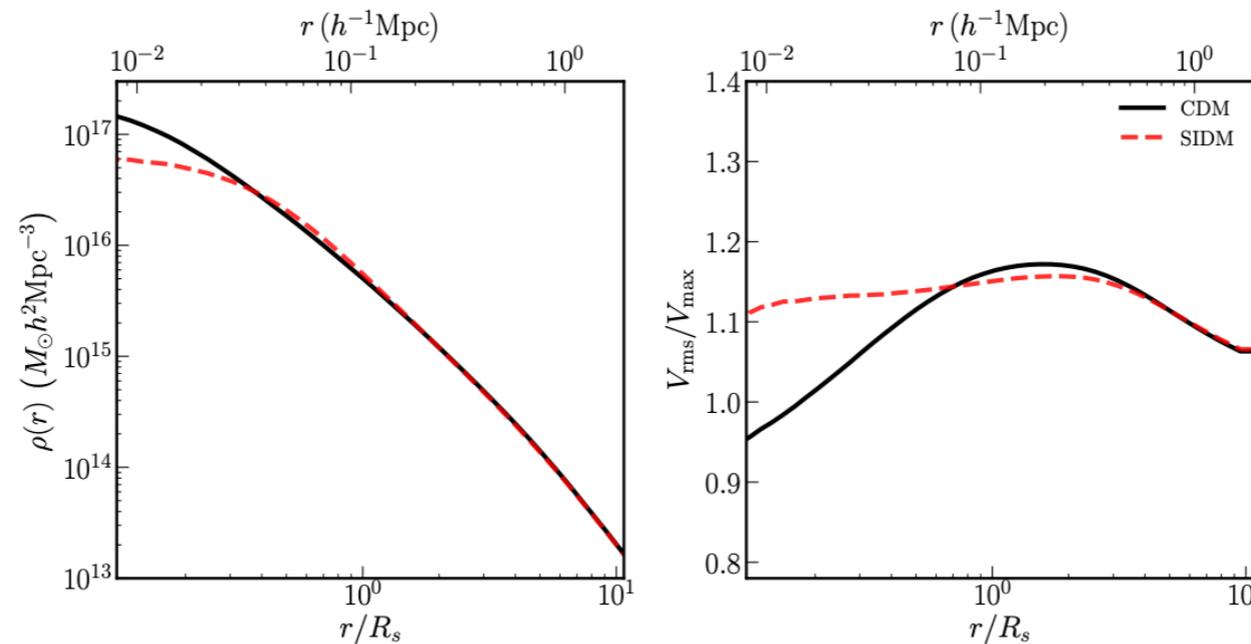
If universe is not Lambda CDM?

Self interacting dark matter and halo profiles

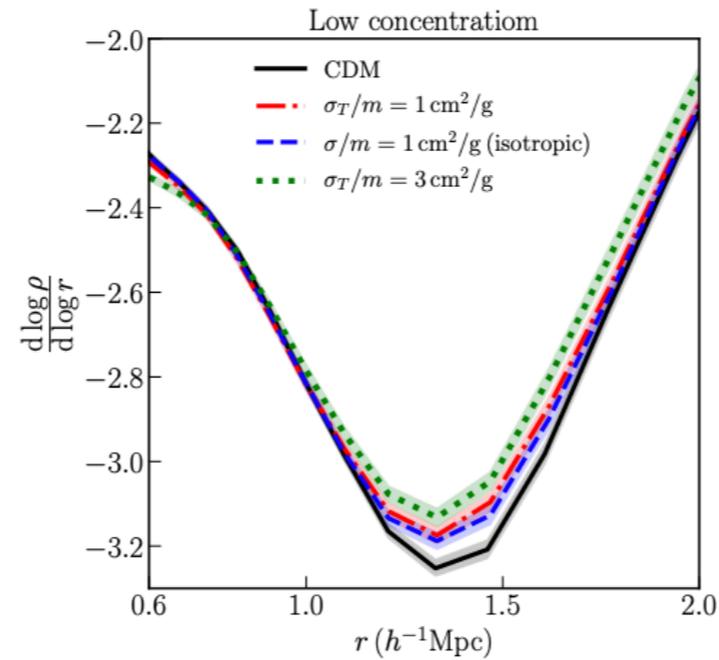
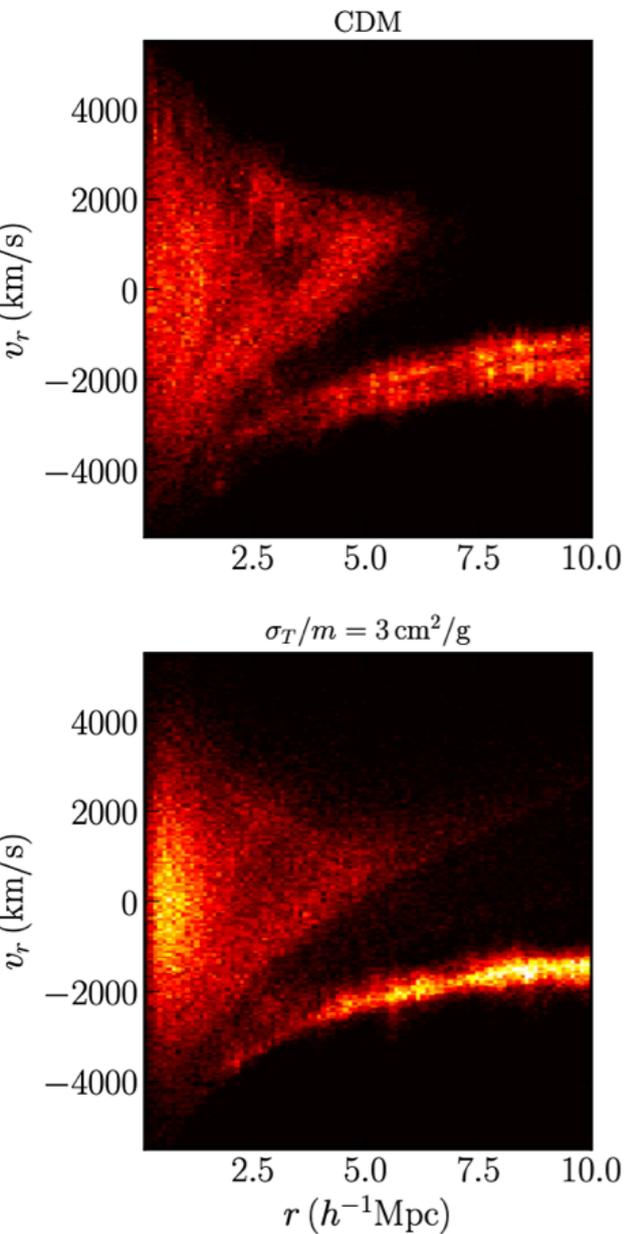


- Particles lose energy their orbits are altered
- Velocity dependent - subhalos and host are at different interaction cross-sections

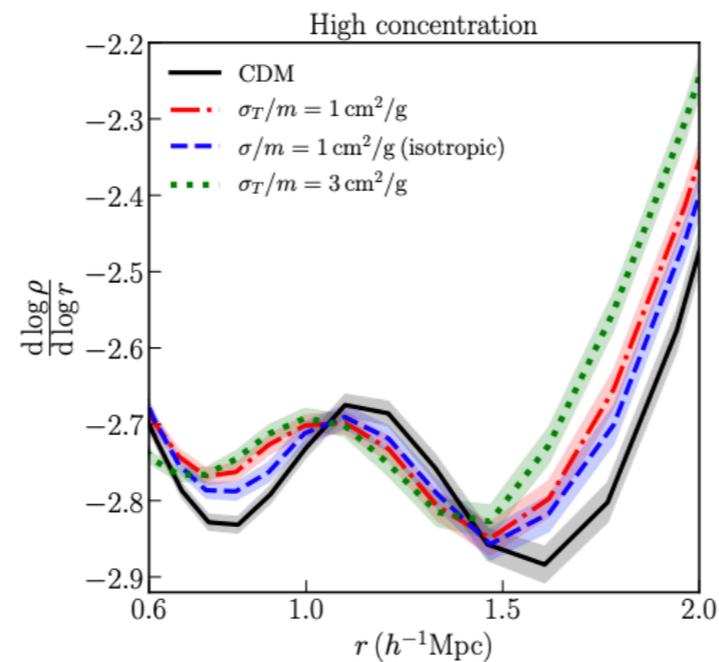
$$\frac{d\sigma}{d\Omega} = \frac{\sigma_0}{2 \left[1 + \frac{v^2}{w^2} \sin^2 \left(\frac{\theta}{2} \right) \right]^2}$$



In the case of self-interacting dark matter we see effects on splashback radius in older halos



Young halos



Old halos

Banerjee, Adhikari et al. 2019

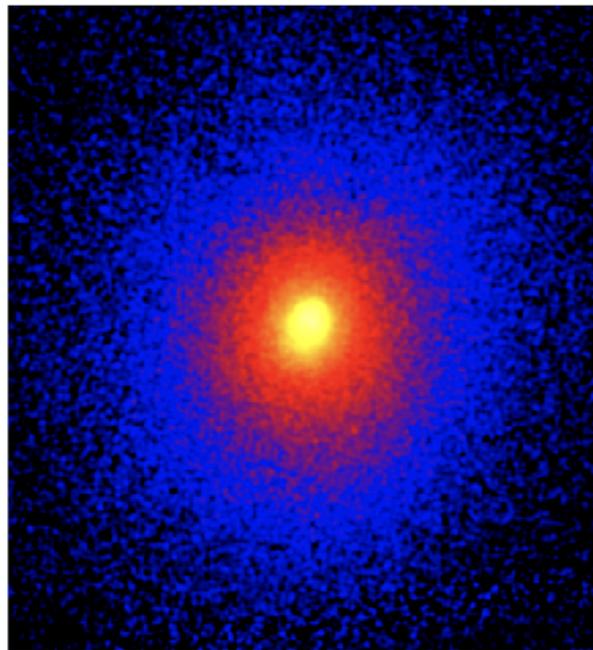
The movement in splashback becomes more prominent when halos are split on accretion history

Observations of the splashback radius

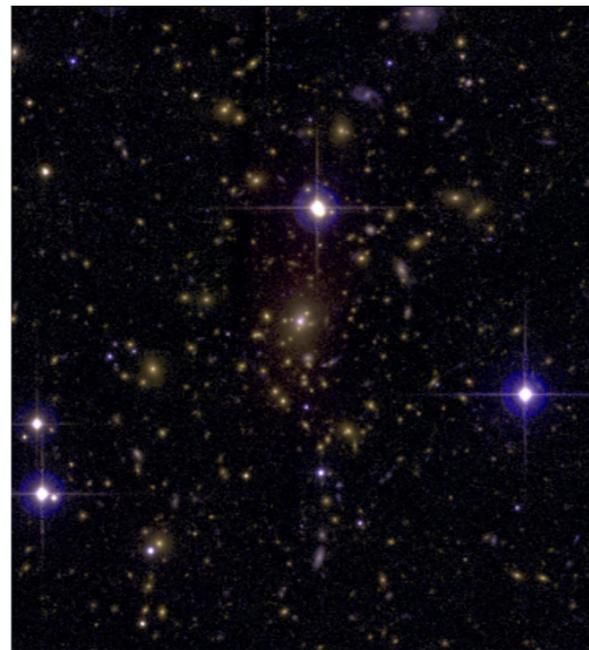
How do we observe dark matter halos?

We study the most massive bound structures in the universe
Cluster mass halos

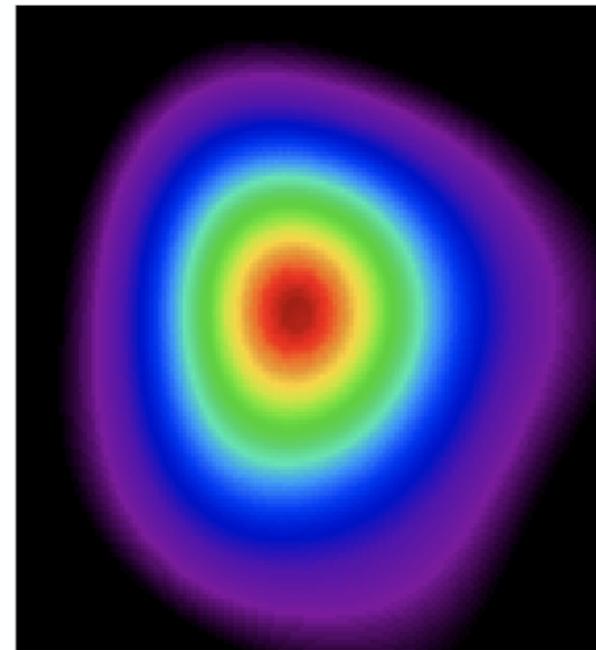
$$10^{14} - 10^{15} M_{sun}$$



X-ray



Optical



tSZ

They can be identified as “clusters” of galaxies in the sky

Galaxy clusters

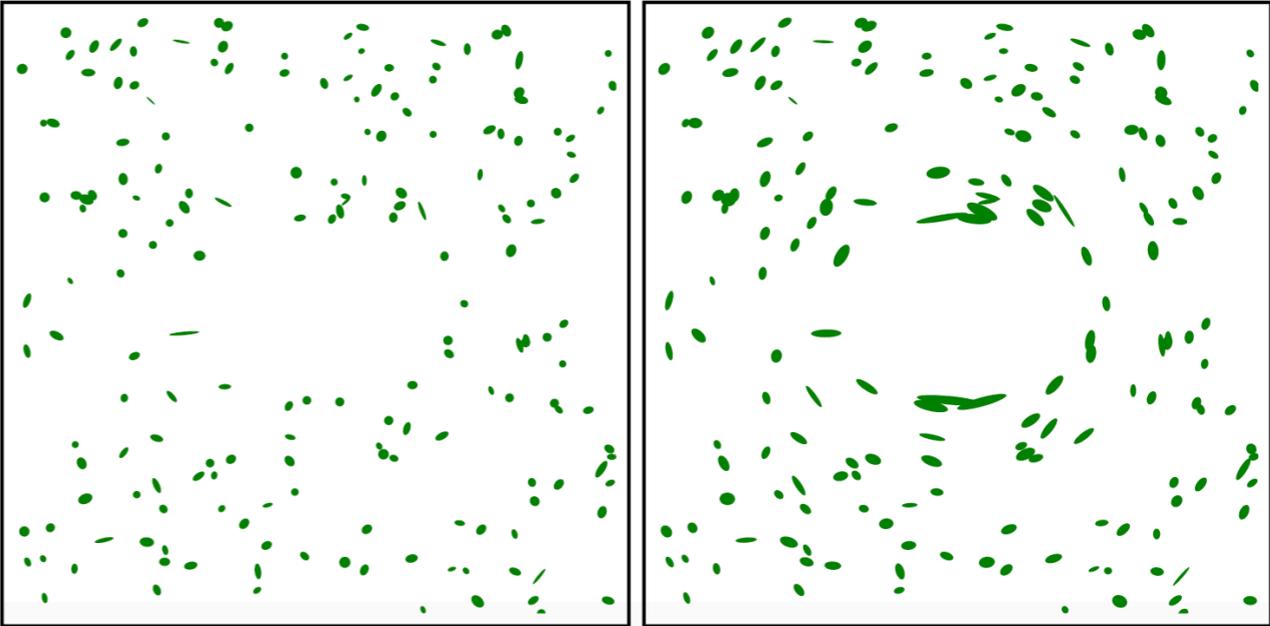
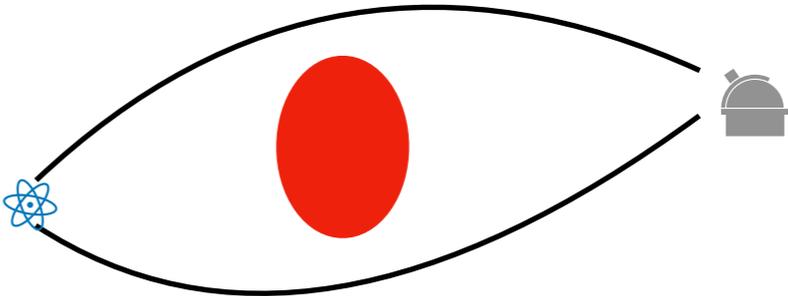
Distribution of Galaxies



Abell 2218

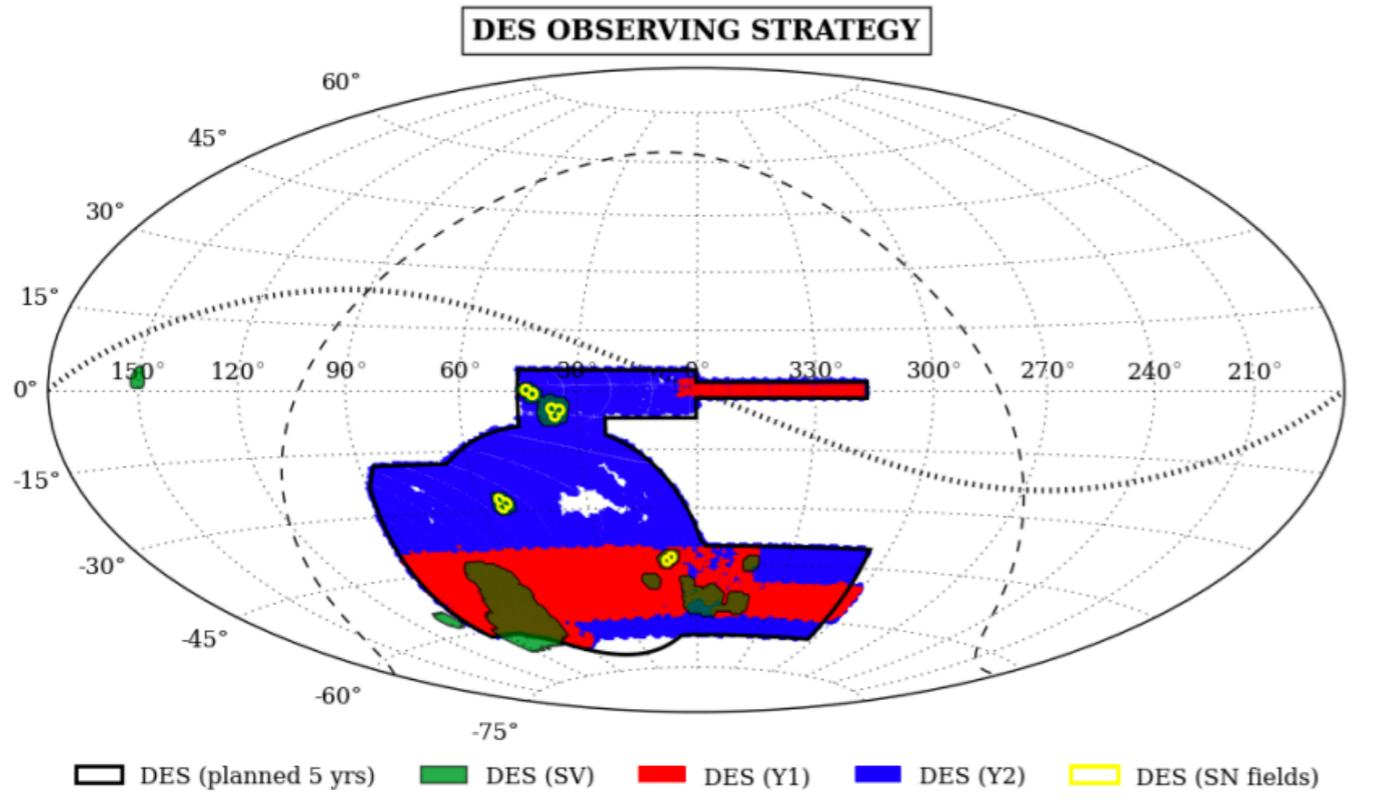
Study the distribution of galaxies that trace the potential of the parent dark matter halos

Lensing of background galaxies



Study the distortion of background galaxies due to massive halo in the line of sight

Dark Energy Survey (DES)



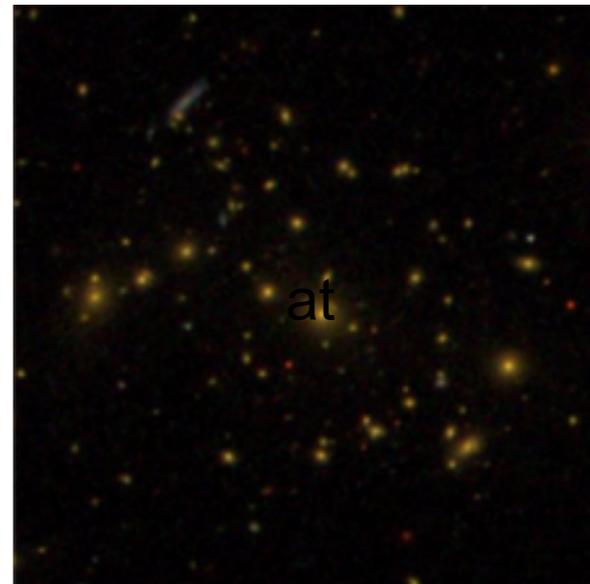
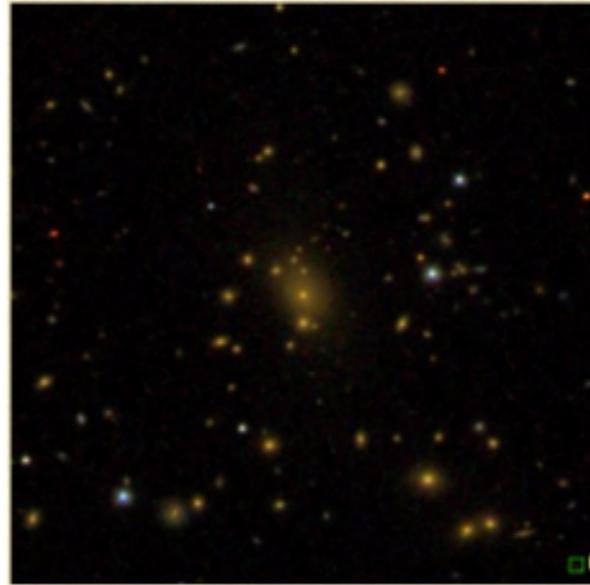
5000 sq. deg
Observes millions of galaxies



<https://www.darkenergysurvey.org/>

Blanco 4m telescope in Chile

Galaxy Clusters in SDSS data selected with the RedMaPPer algorithm



Clusters with richness $20 < \Lambda < 100$ corresponds to $M > 10^{14} M_{\odot} h^{-1}$

8648 RedMaPPer clusters

$0.1 < z < 0.33$

Observations of Splashback radius

DETECTION OF THE SPLASHBACK RADIUS AND HALO ASSEMBLY BIAS OF MASSIVE GALAXY CLUSTERS

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MASAMUNE OGURI^{1,9,13}, AND DAVID N. SPERGEL^{1,3}

The Halo Boundary of Galaxy Clusters in the SDSS

Eric Baxter^{1*}, Chihway Chang², Bhuvnesh Jain¹, Susmita Adhikari³,
Neal Dalal^{3,4}, Andrey Kravtsov^{2,5,6}, Surhud More⁷, Eduardo Rozo⁸,
Eli Rykoff^{9,10}, Ravi K. Sheth^{1,11}

THE SPLASHBACK FEATURE AROUND DES GALAXY CLUSTERS: GALAXY DENSITY AND WEAK LENSING PROFILES

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S. DESAI,³⁴ H. T. DIEHL,²³ J. P. DIETRICH,^{35,36} A. DRLICA-WAGNER,²³ T. F. EIFLER,^{14,37} B. FLAUGHER,²³ P. FOSALBA,³⁸
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G. GUTIERREZ,²³ K. HONSCHIED,^{16,17} D. J. JAMES,⁴² T. JELTEMA,⁴³ E. KRAUSE,⁵ K. KUEHN,⁴⁴ O. LAHAV,¹³ M. LIMA,^{30,45}
M. MARCH,² J. L. MARSHALL,⁴⁶ P. MARTINI,^{16,47} P. MELCHIOR,⁴⁸ F. MENANTEAU,^{32,33} R. MIQUEL,^{3,49} J. J. MOHR,^{7,35,36}
B. NORD,²³ R. L. C. OGANDO,^{30,31} A. A. PLAZAS,¹⁴ E. SANCHEZ,⁵⁰ V. SCARPINE,²³ R. SCHINDLER,⁹ M. SCHUBNEL,⁴¹
I. SEVILLA-NOARBE,⁵⁰ M. SMITH,⁵¹ R. C. SMITH,²⁵ M. SOARES-SANTOS,²³ F. SOBREIRA,^{30,52} E. SUCHYTA,⁵³
M. E. C. SWANSON,³³ G. TARLE,⁴¹ AND J. WELLER^{6,7,35}

(DES COLLABORATION)





Cluster - galaxy cross correlation

Measurement - Number density of galaxy in projection as a function of radius

$$\Sigma(R) = \int_{-h_{\max}}^{h_{\max}} dh \rho(\sqrt{R^2 + h^2})$$

$$\rho(r) = \rho^{\text{coll}}(r) + \rho^{\text{infall}}(r),$$

$$\rho^{\text{coll}}(r) = \rho^{\text{Ein}}(r) f_{\text{trans}}(r)$$

$$\rho^{\text{Ein}}(r) = \rho_s \exp\left(-\frac{2}{\alpha} \left[\left(\frac{r}{r_s}\right)^\alpha - 1\right]\right)$$

$$f_{\text{trans}}(r) = \left[1 + \left(\frac{r}{r_t}\right)^\beta\right]^{-\gamma/\beta},$$

$$\rho^{\text{infall}}(r) = \rho_0 \left(\frac{r}{r_0}\right)^{-s_e},$$

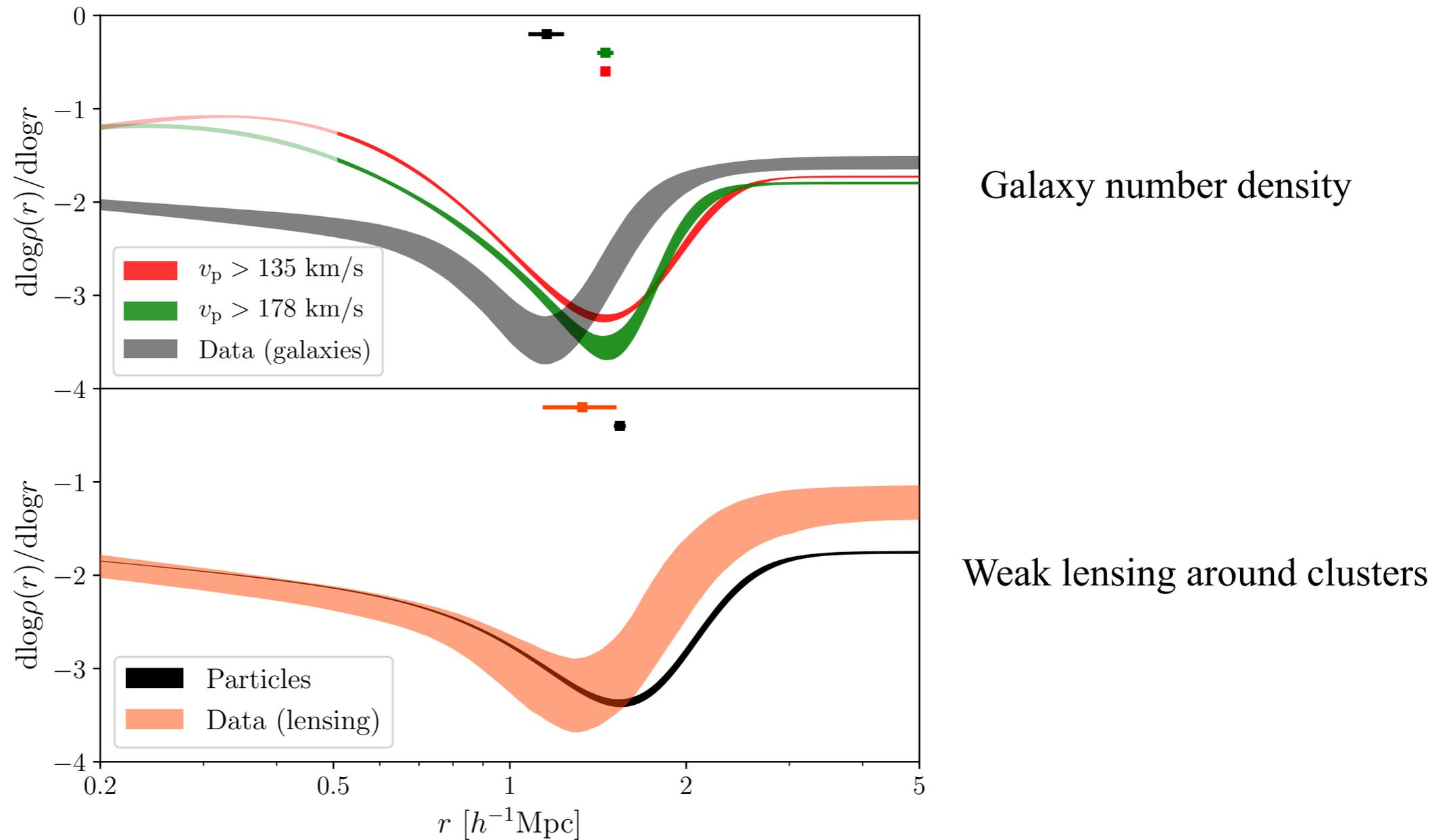
Stack clusters based on richness

richness > 20

$M > 1e14 M_{\text{sun}} h^{-1}$

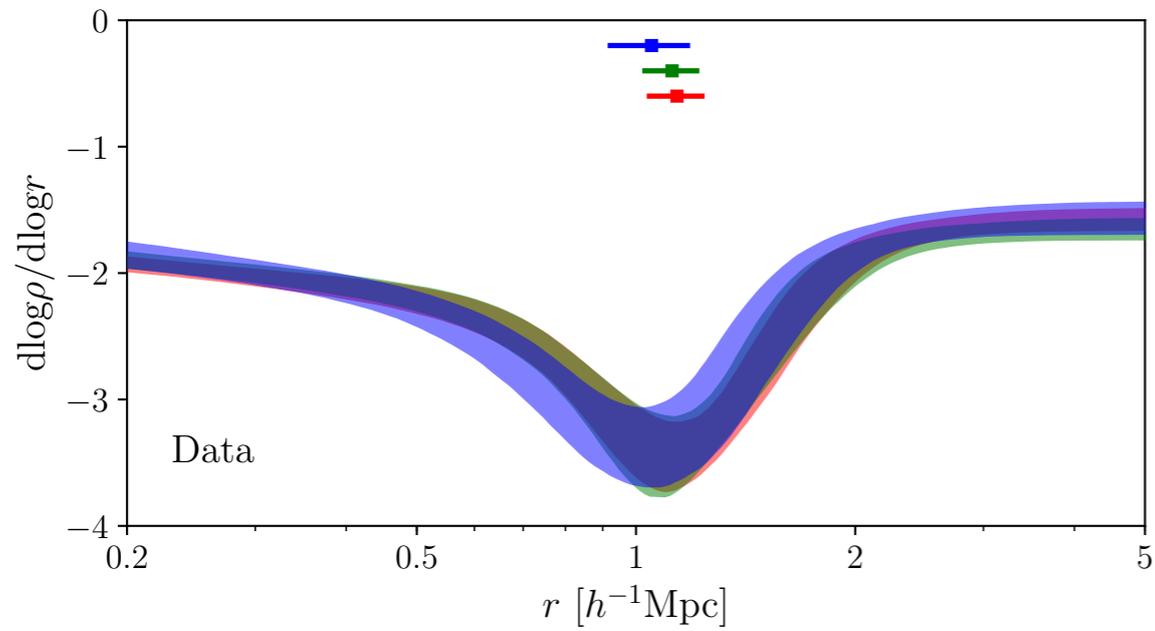
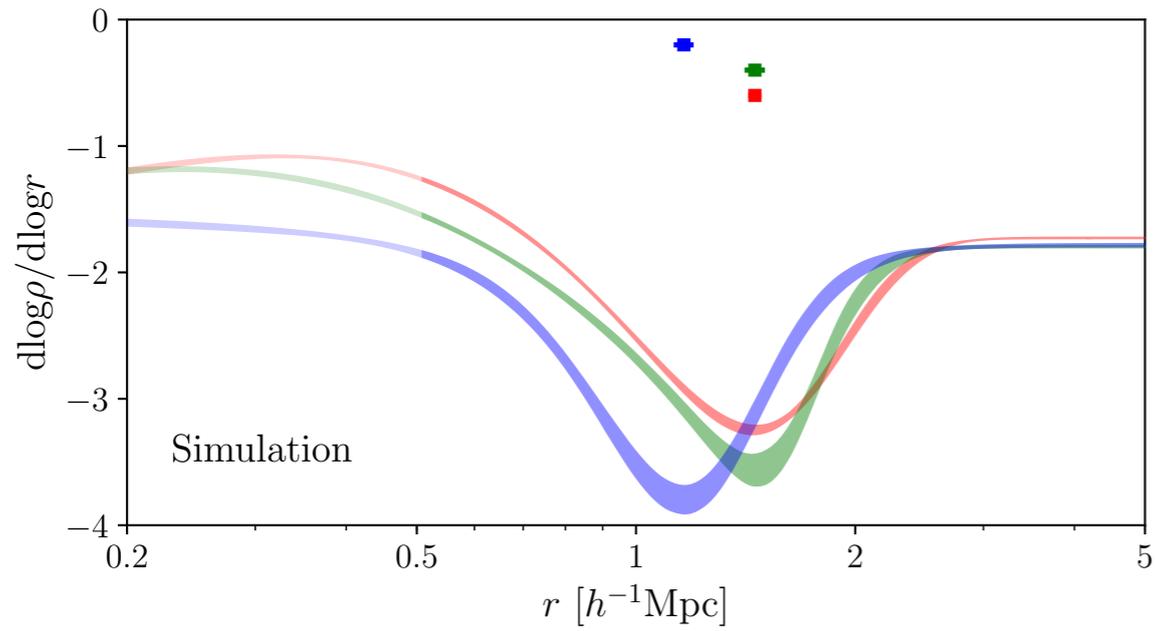
Splashback radius in DES Y1 results

First measurement in weak lensing around halos

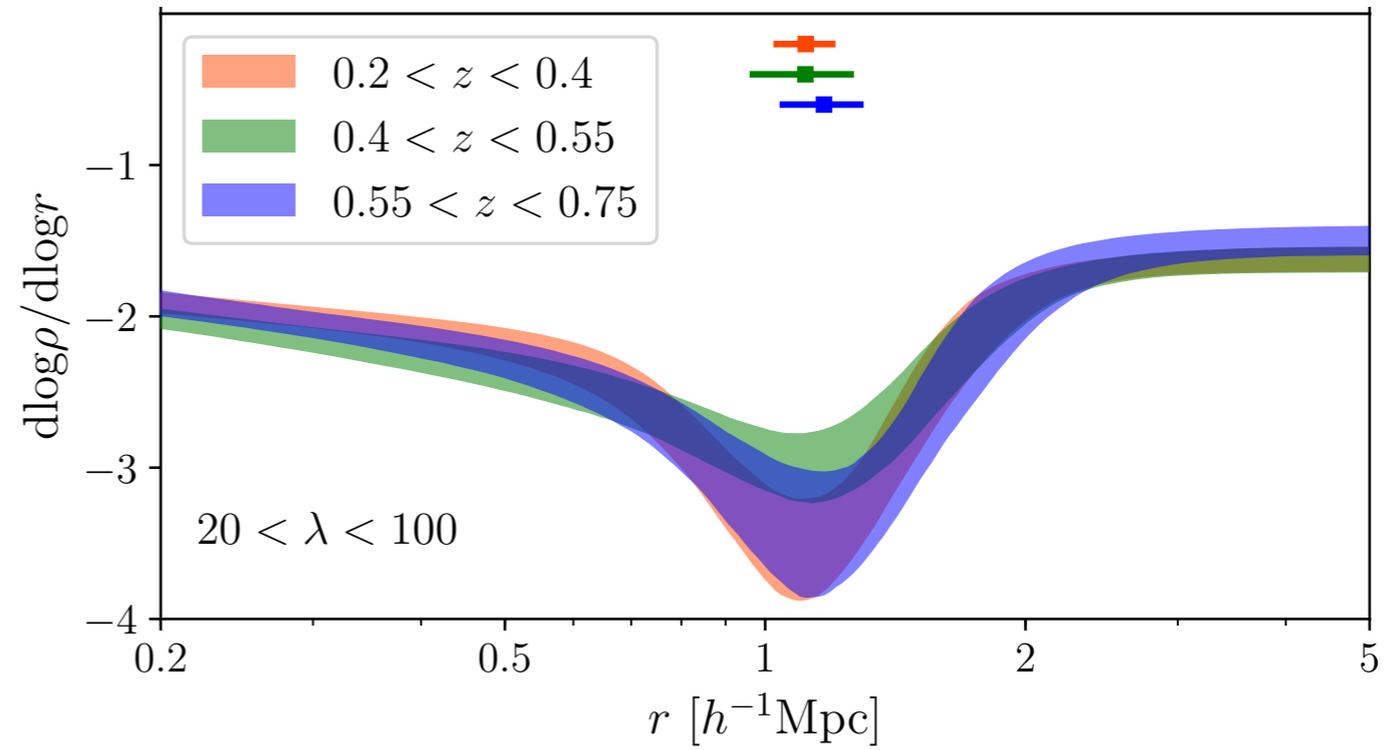


Chang, +, Adhikari et al. 2017

Discrepancy persists in the lensing splashback radius as well



Outstanding issues



Chang, Baxter, Jain, Sanchez, Adhikari et al. 2017

No movement with redshift of host cluster

No movement with galaxy magnitude

Why is splashback discrepant with simulations?

a) Dynamical Friction?

b) New Physics?

c) Observational bias?

Cluster selection?

Projection effects? (*Busch & White 2017*)

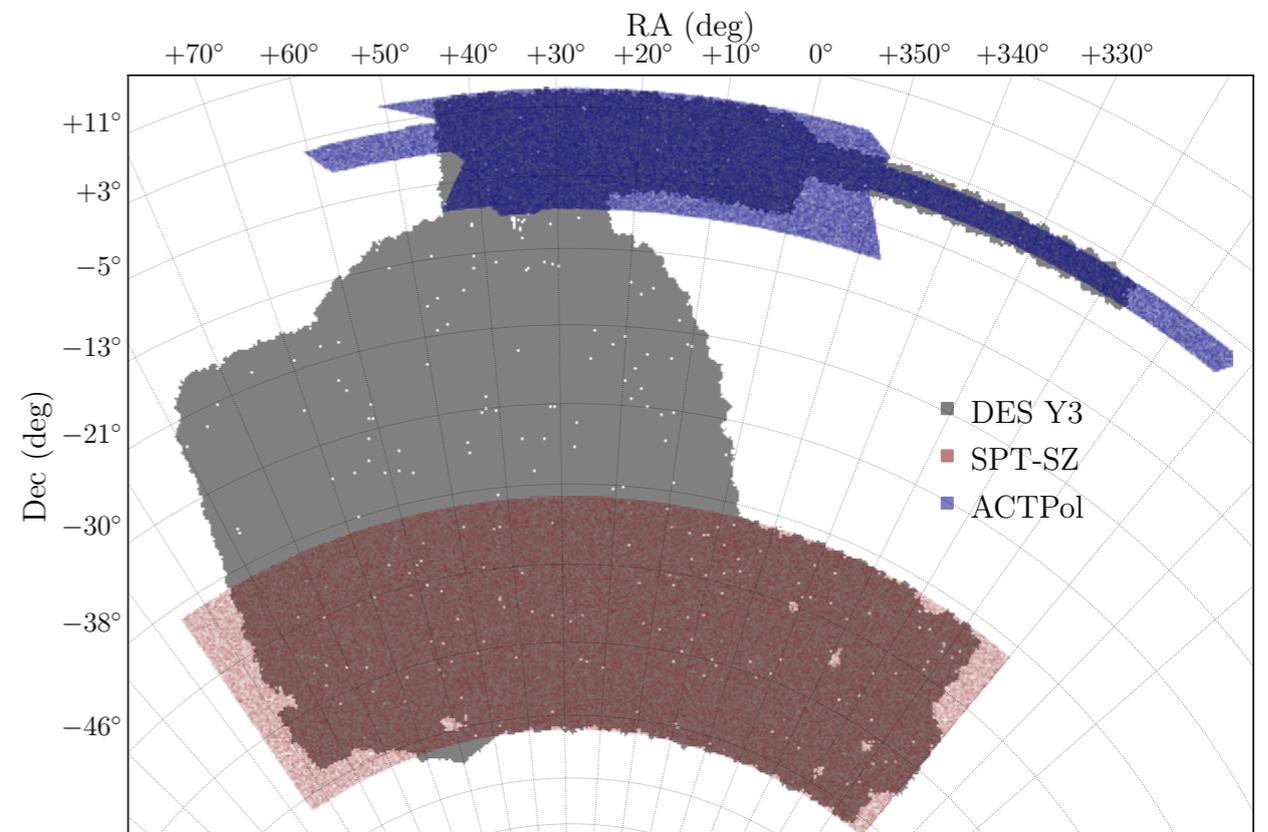
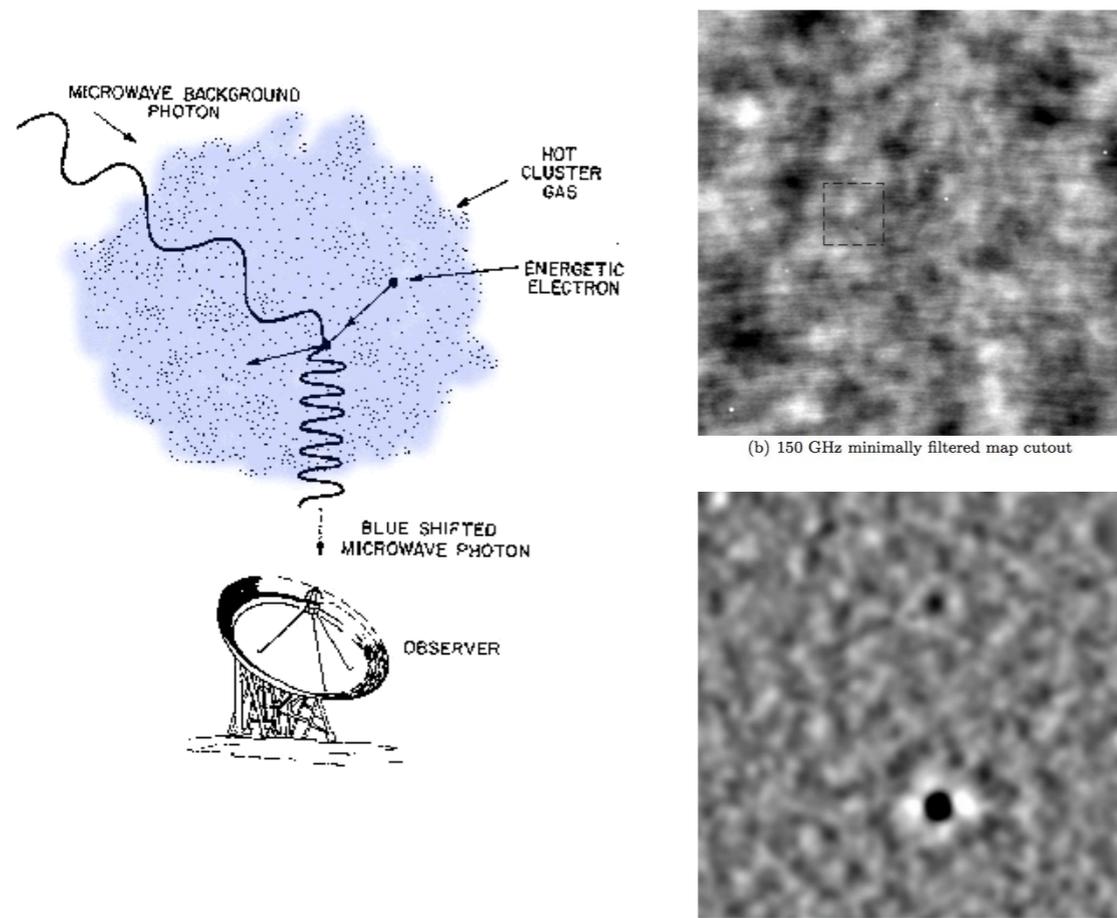
Orientation bias?

Aperture selection? (*Busch & White 2017*)

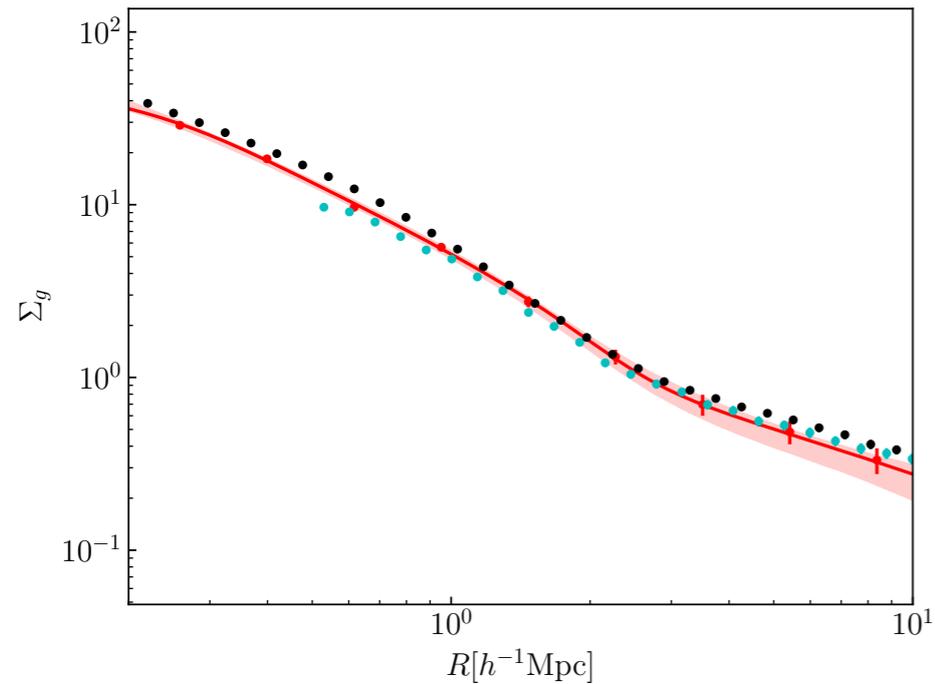
Different cluster selection method - SZ selected clusters

Splashback radius in SZ clusters from the South Pole telescope (SPT) and Atacama Cosmology telescope (ACT)

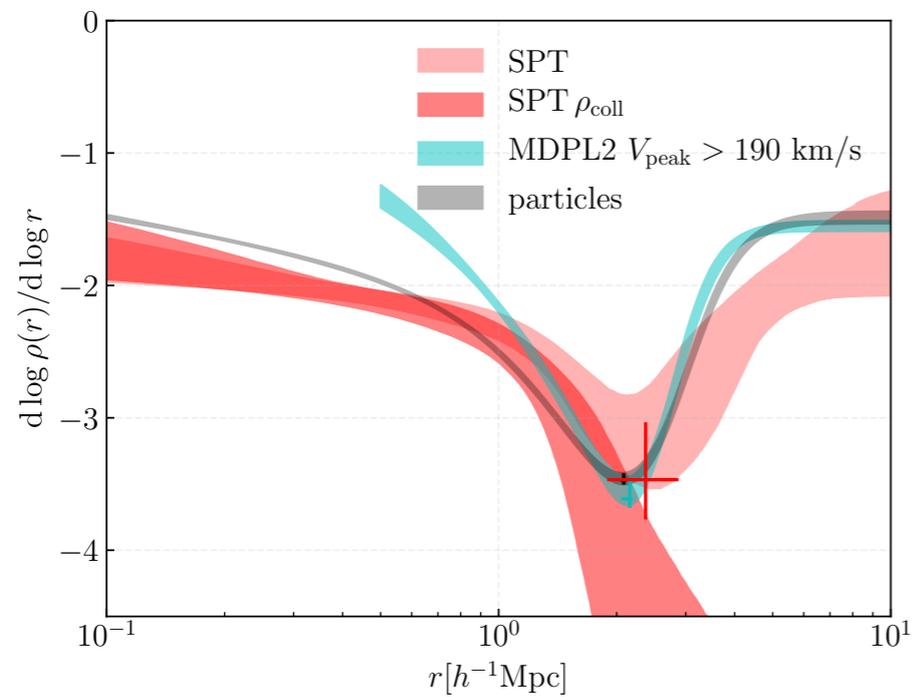
Clusters seen as a temperature decrement in CMB



Splashback radius in SPT SZ clusters, DES galaxies



Splashback radius SZ clusters are statistically consistent with simulations

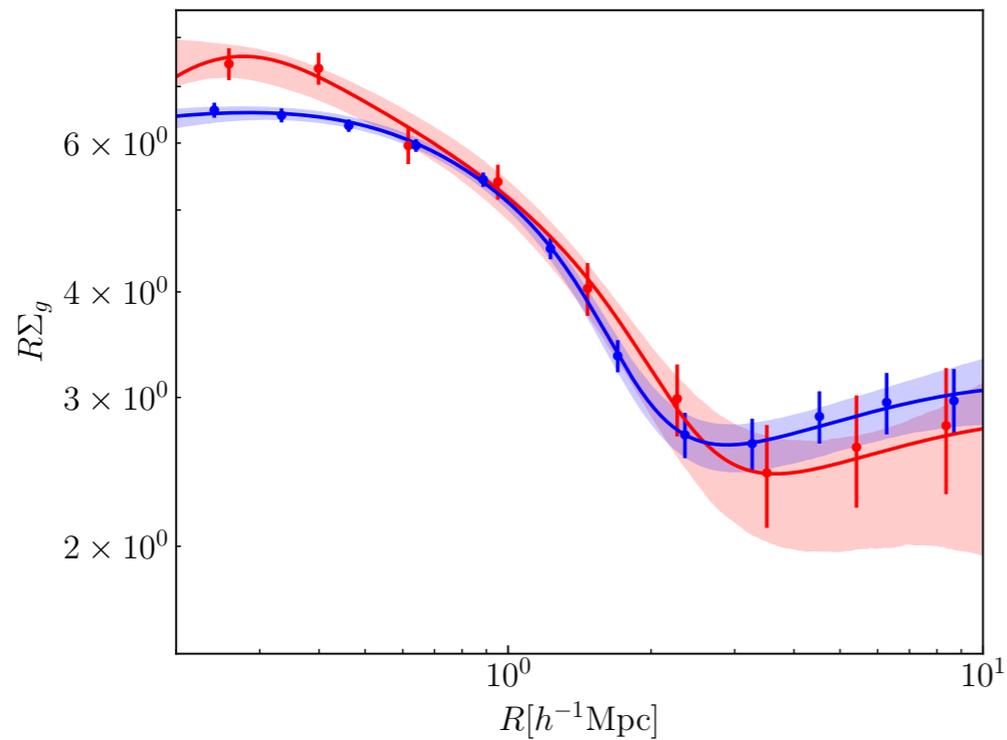


Pink - Slope of the fitted density profile
Black- Particles from MDPL2
Blue - Subhalos abundance matched

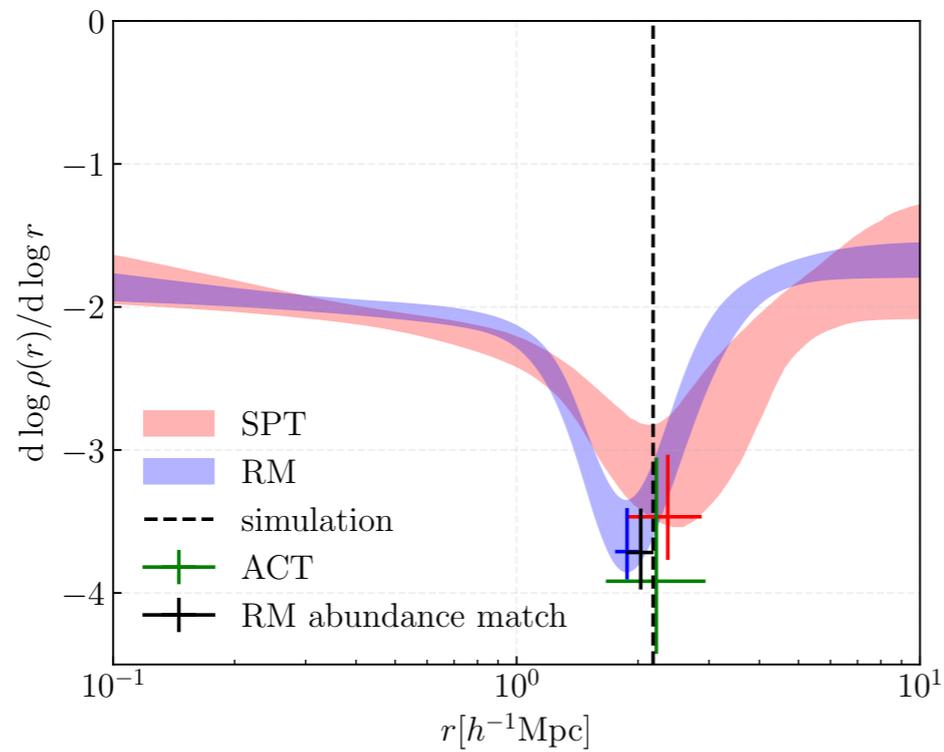
Hyeon-Shin, Adhikari et al. 2019

Consistent with Zuercher & More 2019 who did a similar analysis with Planck clusters

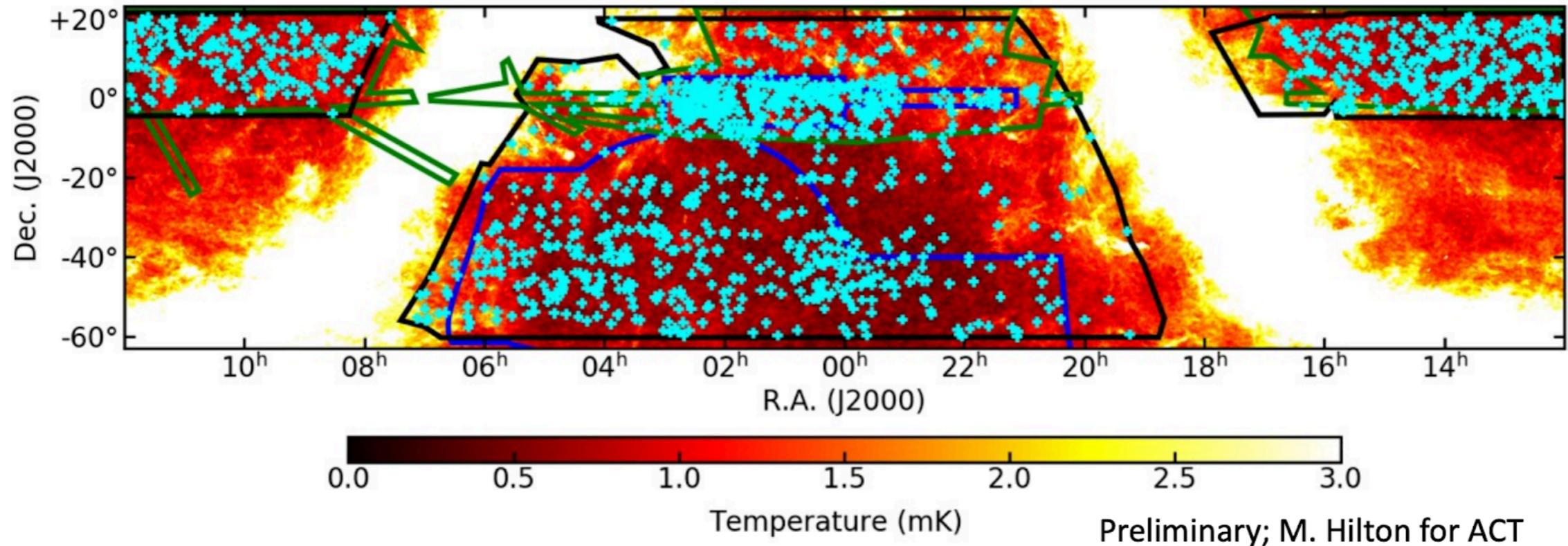
Comparison with RedMaPPer



**RM and SPT are consistent within 1 sigma,
but RM is inconsistent with sims.**



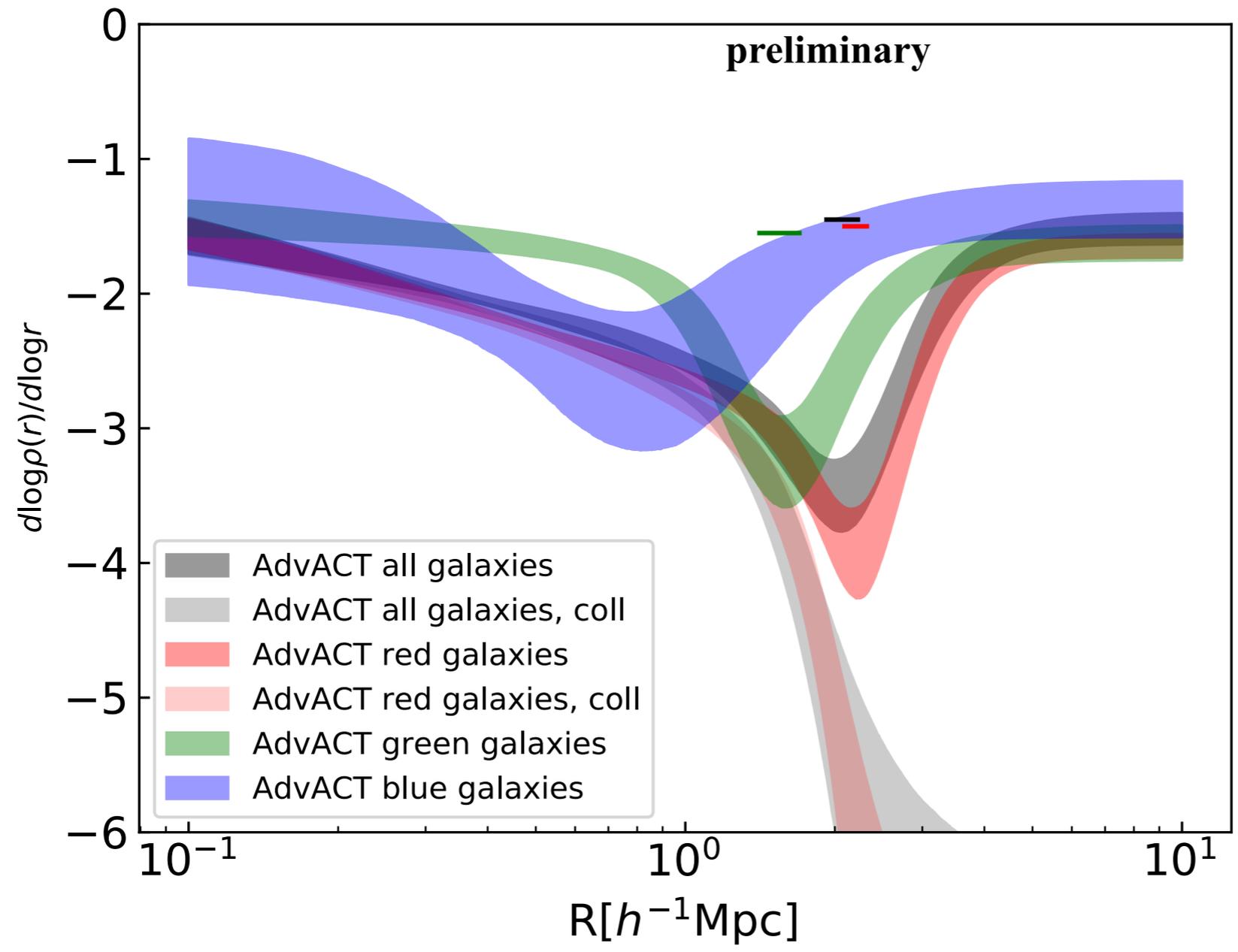
New AdvACT cluster sample



863 clusters (subject to change) in the DES footprint having $\text{SNR} > 4$, w/
 $0.15 < z < 0.7$
 $\langle M_{500c} \rangle = 3.0e14 \text{ Msun}/h$
 $\langle z \rangle = 0.44$

Galaxy quenching in Dark Matter halos

New AdvACT sample



(in prep)

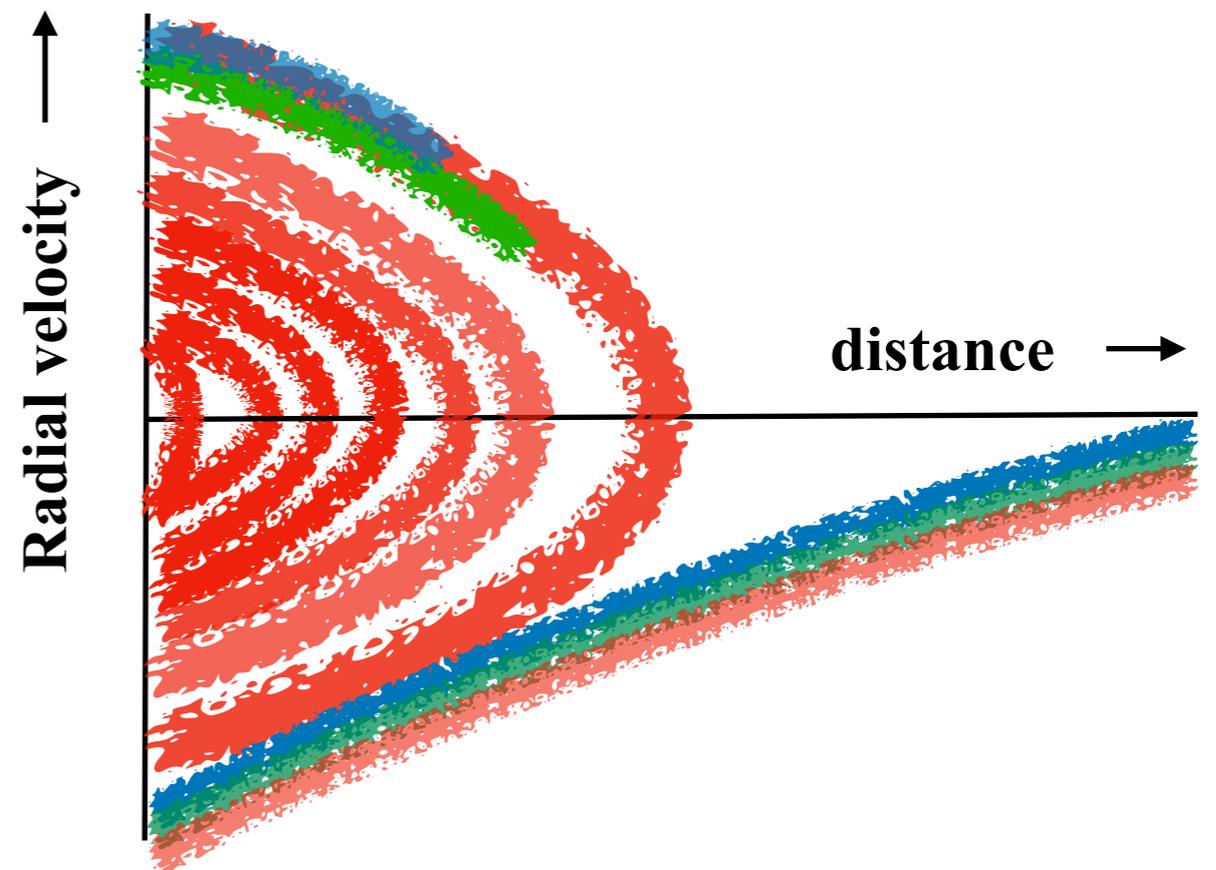
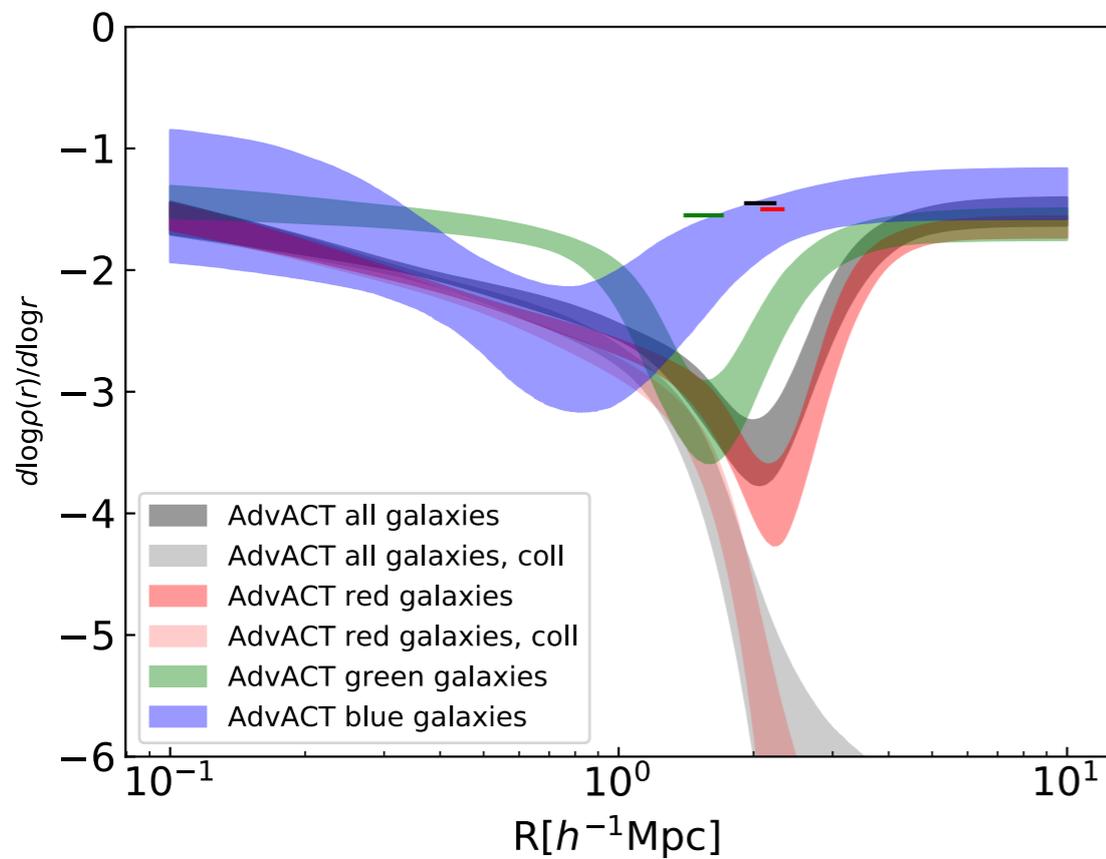
~700 cluste

$$\text{SFR}_{\text{sat}}(t) = \begin{cases} \text{SFR}_{\text{cen}}(t) & t < t_{Q, \text{start}} \\ \text{SFR}_{\text{cen}}(t_{Q, \text{start}}) e^{\left\{ -\frac{(t-t_{Q, \text{start}})}{\tau_{Q, \text{fade}}} \right\}} & t > t_{Q, \text{start}} \end{cases}$$

The splashback radius as a clock in the halo

Galaxies stop forming stars with time as they fall into a halo

Blue star-forming galaxies turn into red and dead galaxies



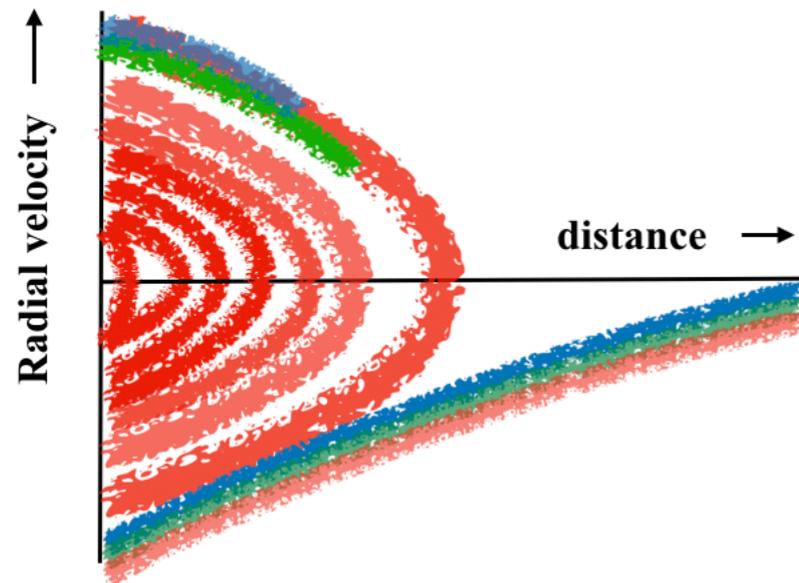
Longer delay , shorter quenching

Minimum traces the time spent in the cluster by a population of galaxies

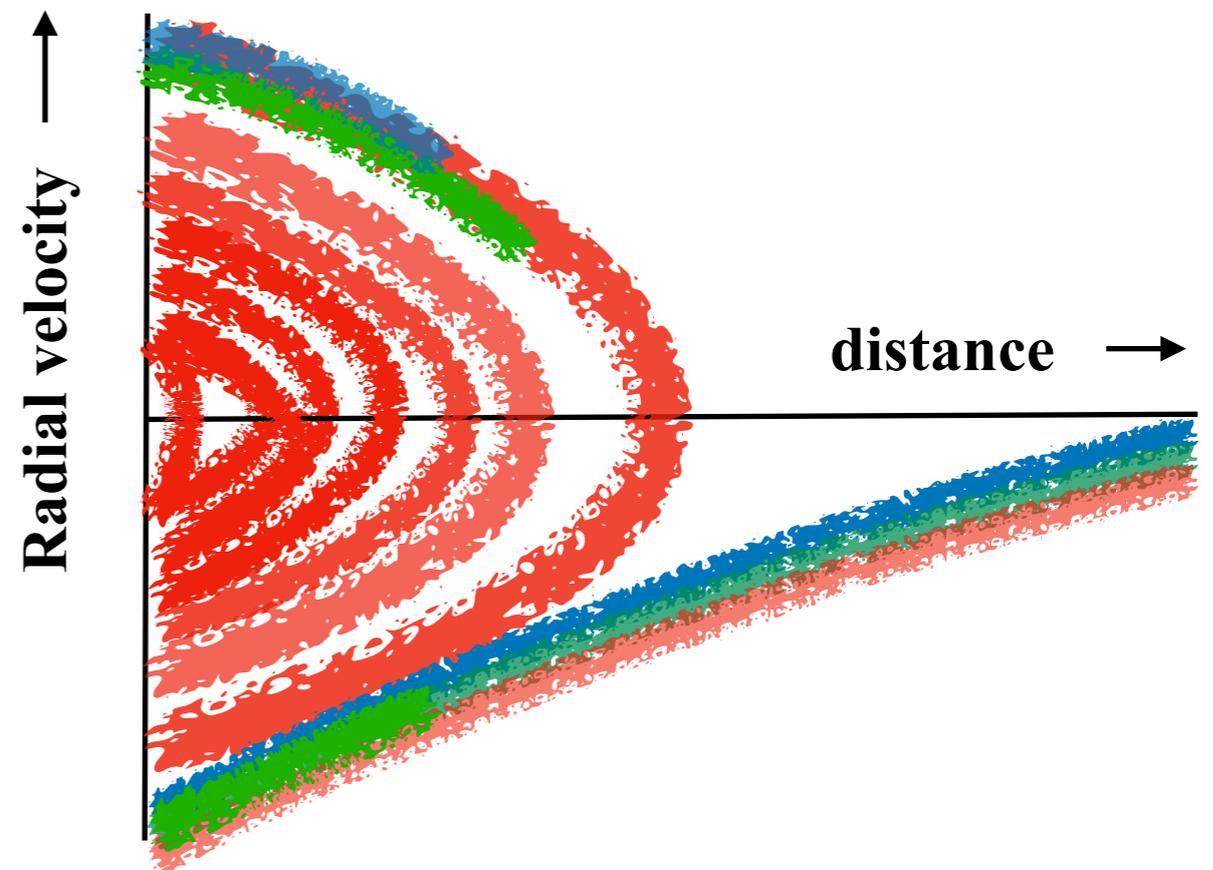
The splashback radius as a clock in the halo

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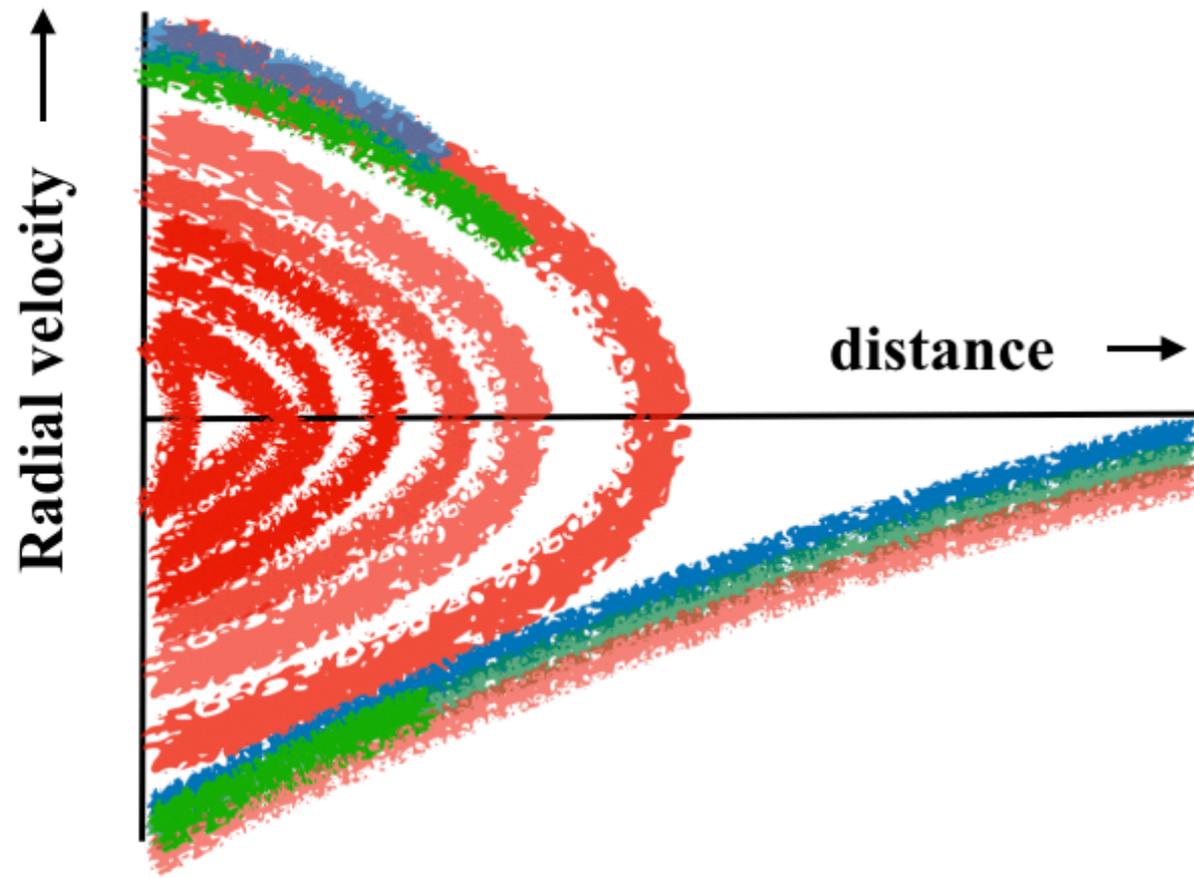


□ Longer delay , shorter quenching □

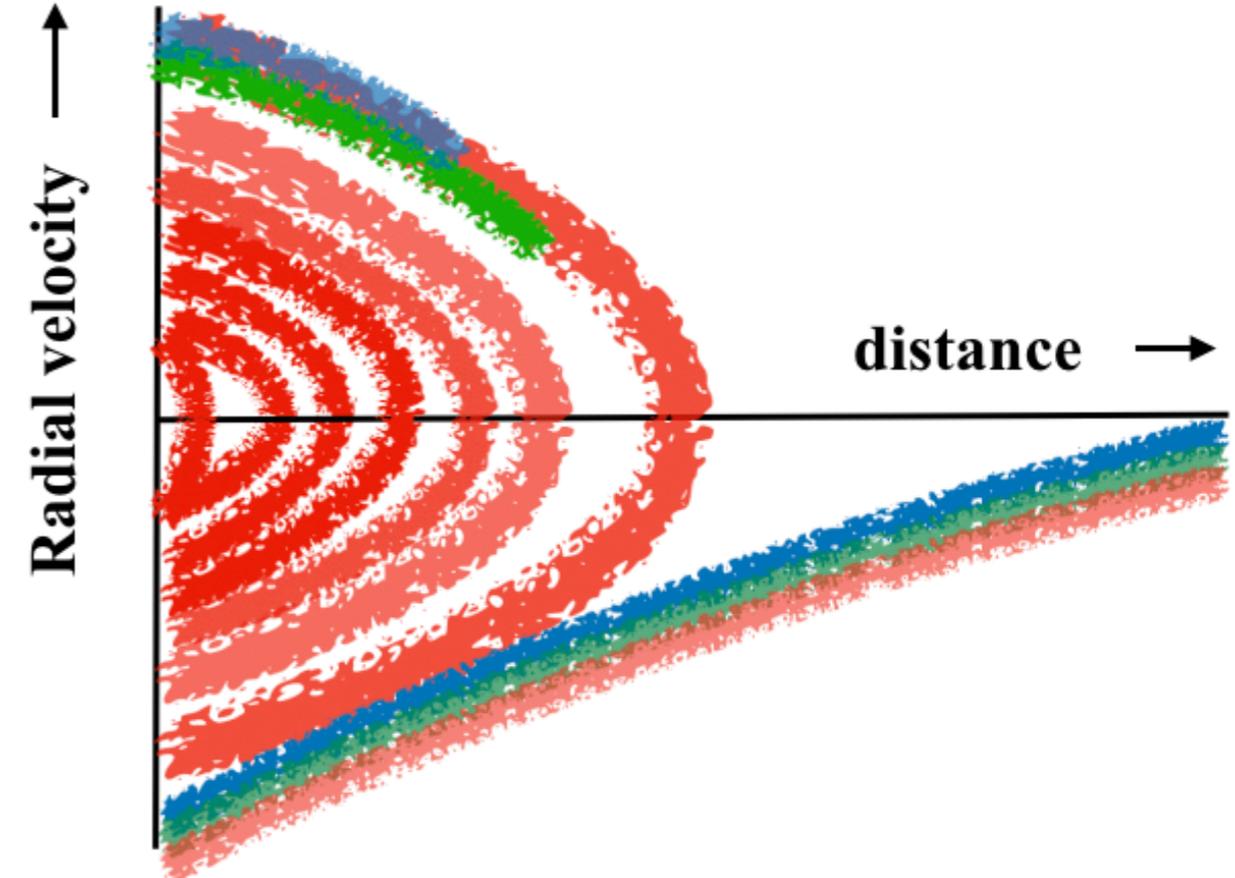


Short delay , long quenching

Minimum traces the time spent in the cluster by a population of galaxies



Short delay , long quenching



Longer delay , shorter quenching

Summary

- **The structure of dark matter halos contain information about the history of the universe**
- **The edges of halos can be understood through simple physical model**
- **The location of the edge is traced by the splashback radius that can be measured observationally**
- **Sensitive to modified gravity models**
- **Sensitive to models of self interacting dark matter, potentially any model that can change the energetics of dark matter particles**
- **A distinct scale in a halo that can tell us about galaxy evolution**